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Large Tunable Thermo-phase in SC-QD-SC Junction





The electric response due to thermal gradients



- An emerging player in energy conversion and production
- A powerful tool in **probing** strongly correlated effects

Thermoelectricity in Superconductors



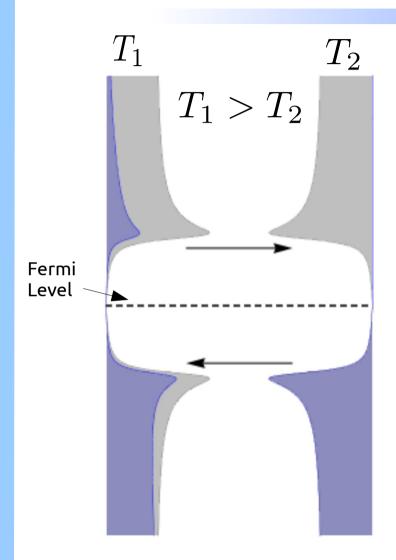
- Thermoelectricity in SC was long considered absent:
 - The SC induces thermo-phase instead of thermoelectricity under temperature gradients



- The SC in inherently particle-hole symmetric
- Theory ~40's Experiment ~ 70's

The SC-I-SC Junction



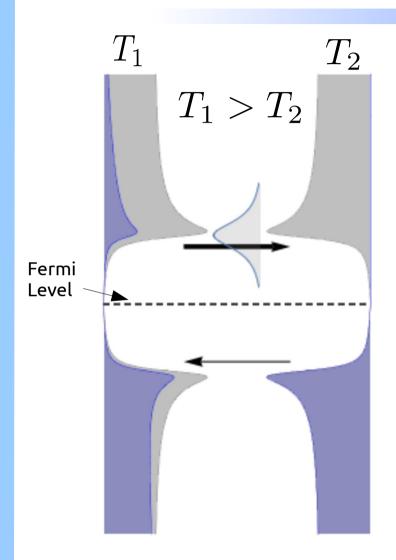


- p-h symmetric
- Any assymetry would come from impurities
- Rather well studied (The,Exp)

Small thermal response, mainly around the transition temperature

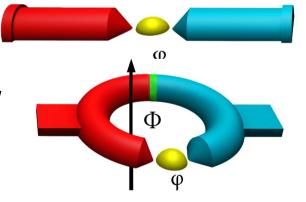
The SC-QD-SC Junction





- Offering tunable p-h symmetry breaking
- Offering substantial thermal response

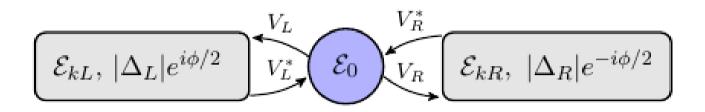
Experimentally relevant [1,2]



- [1] Giazotto et al, Nat Nano 11, 258–262 (2016)
- [2] De Franceschi et al, Nat Nano 5, 703–711 (2010)

Model and Methods

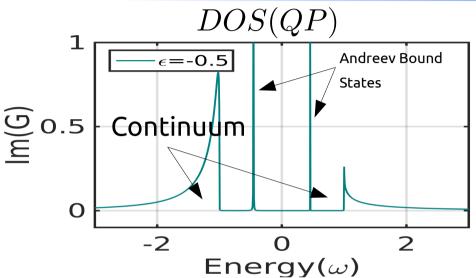




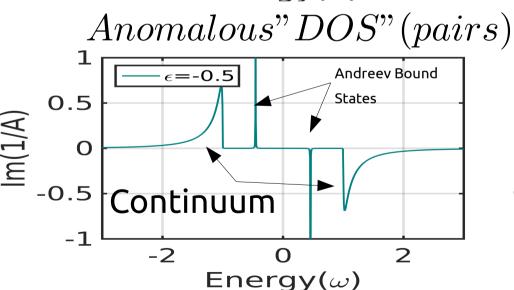
- An energy level between two SC reservoirs
- We apply a temperature gradient
- Using the Keldysh non-eq Greens Function we calculate the thermal response
- We get: $I(\phi, \Delta T, T_{av}, \epsilon)$ Temperature Level energy difference

Transmission





• Quasiparticle current travels through the ${\bf continuum}(\Delta T)$



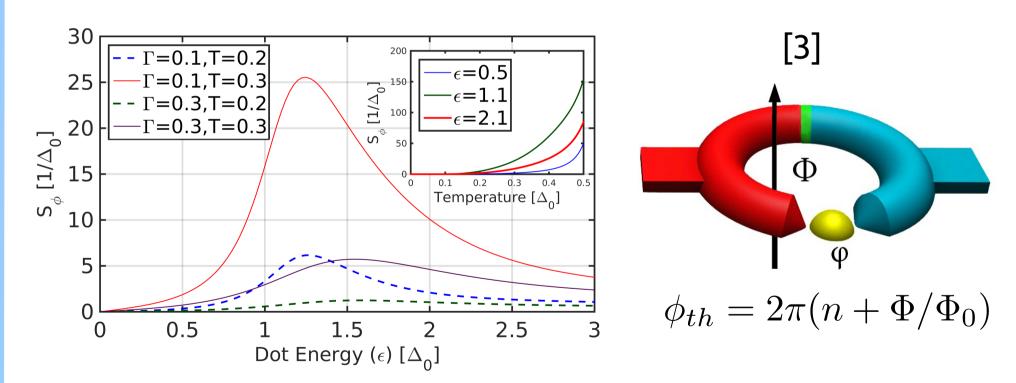
Supercurrent travels through the **ABS** and the **continuum**(ϕ)

 Pair->QP tranmission is p-h symmetric, no contribution

Thermo-Phase Seebeck Coefficient



$$S_{\phi} = -\left(\frac{dI/dT}{dI/d\phi}\right)_{I=0}, \quad \phi_{th} = S_{\phi}\Delta T$$

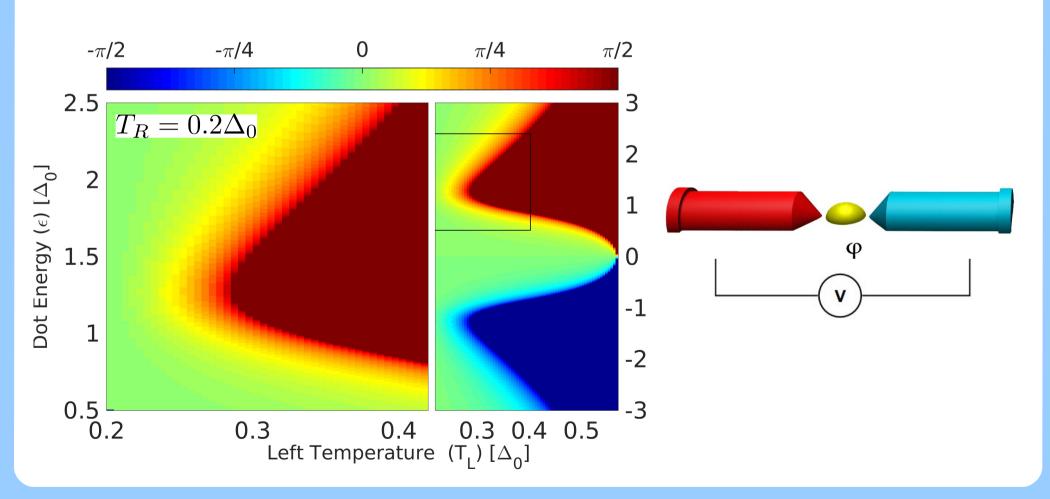


• [3] Garland et al, Physics letters A, 47 423 (1974)

Beyond Linear Response

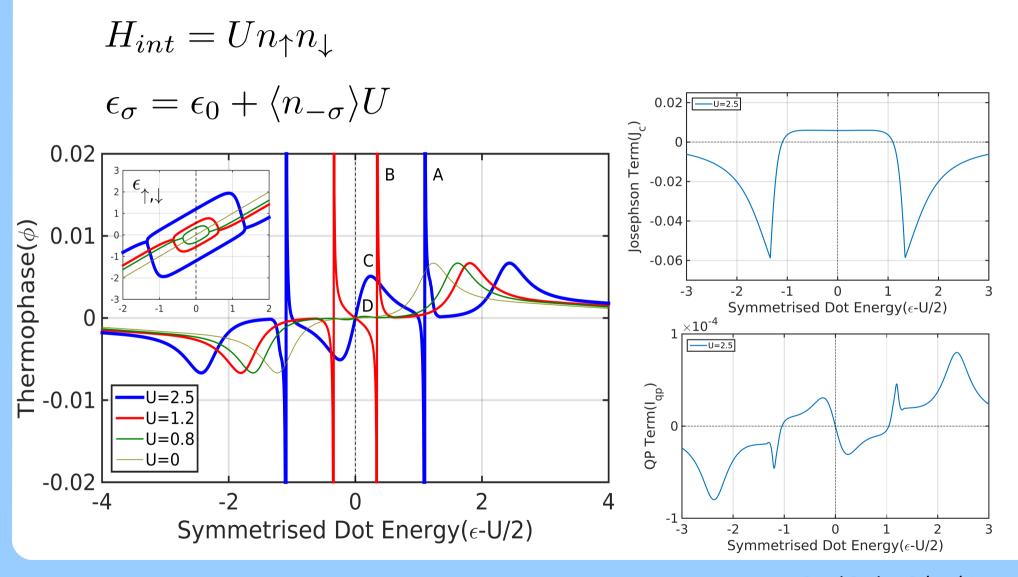


$$I(\phi_{th}, \Delta T, T, \epsilon) = 0$$



Thermo-phase with Coulomb Interaction



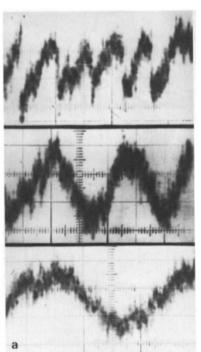


Capri Spring School, 18/04/16

Summary



We suggest and investigate a realizable setup for controlling and measuring substantial thermal response in SC



We lay the grounds and set the tools for a study of the AC thermal Josephson effect

[4] Panaitov et al, Phys lett A, 100, 301 (1984)



$$H_S = \epsilon_{ks} \sum_{s,k,\sigma} c_{sk\sigma}^{\dagger} c_{sk\sigma} + \sum_{s,k,\sigma} \Delta_s c_{sk\sigma}^{\dagger} c_{s-k-\sigma}^{\dagger} + H.c$$

$$H_{QD} = \sum_{\sigma} \epsilon_{\sigma} d_{\sigma}^{\dagger} d_{\sigma}$$

$$H_V = \sum_{s,k,\sigma} V_{ks} c_{sk\sigma}^{\dagger} d_{\sigma} + H.c$$

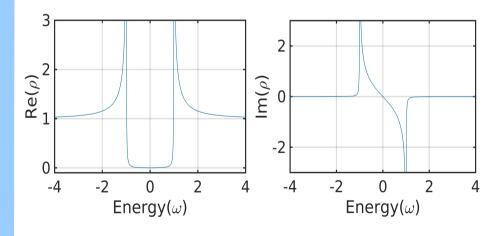
$$H_{int} = U n_{\uparrow} n_{\downarrow}$$



$$\hat{G}^r_{\sigma}(t) = -i\theta(t) \langle \{ \boldsymbol{\Psi}(\mathbf{t}), \boldsymbol{\Psi}^{\dagger}(0) \} \rangle,$$

$$\Psi^{\dagger} = \begin{pmatrix} d_{\sigma}^{\dagger} \\ d_{\overline{\sigma}} \end{pmatrix}$$

$$\hat{\Sigma}_s^r(\omega) = -\frac{i}{2} \Gamma_s \rho_s(\omega) \begin{pmatrix} 1 & -\frac{\Delta}{\omega} e^{i\phi_s} \\ -\frac{\Delta}{\omega} e^{-i\phi_s} & 1 \end{pmatrix}$$



$$\rho_s(\omega) = \begin{cases} \frac{|\omega|}{\sqrt{(\omega)^2 - \Delta_s^2}} & |\omega| > \Delta_s \\ \frac{\omega}{i\sqrt{\Delta_s^2 - (\omega)^2}} & |\omega| < \Delta_s \end{cases}$$



$$I_{qp} = \frac{e}{h} \sum_{\sigma} \int_{-\infty}^{\infty} d\omega \frac{f(\omega)}{dT} \Gamma \operatorname{Re}[\rho(\omega)] \operatorname{Im}[-G_{11}^{r}(\omega)] \Delta T$$

$$I_{sc} = \frac{e}{h} \sum_{\sigma} \int_{-\infty}^{\infty} d\omega f(\omega) \frac{\Delta^2 \Gamma^2}{\omega^2 - \Delta^2} \operatorname{Im}[1/A(\omega)] \sin(\phi)$$

