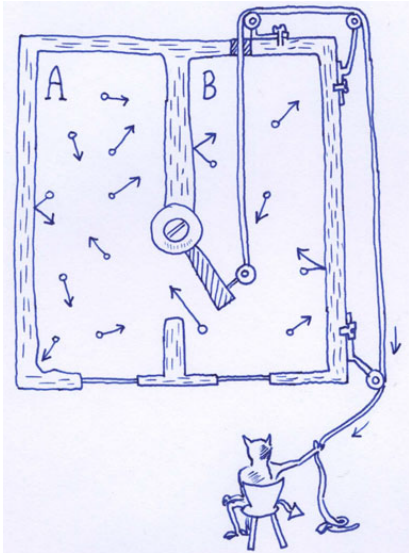


LECTURE 2, Jukka Pekola

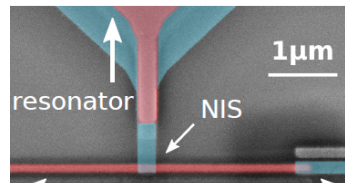
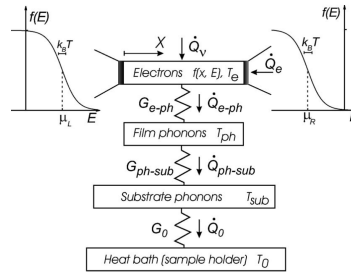
Outline

Maxwell's demon

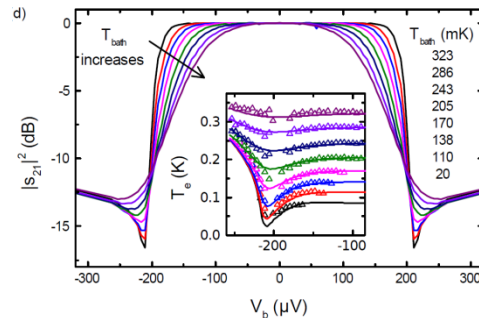
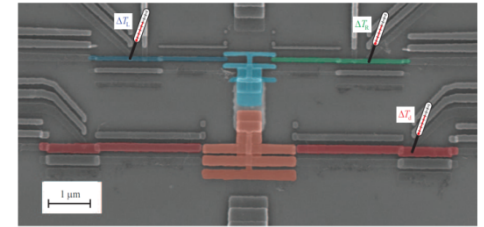


... role of information

Energy relaxation and thermometry

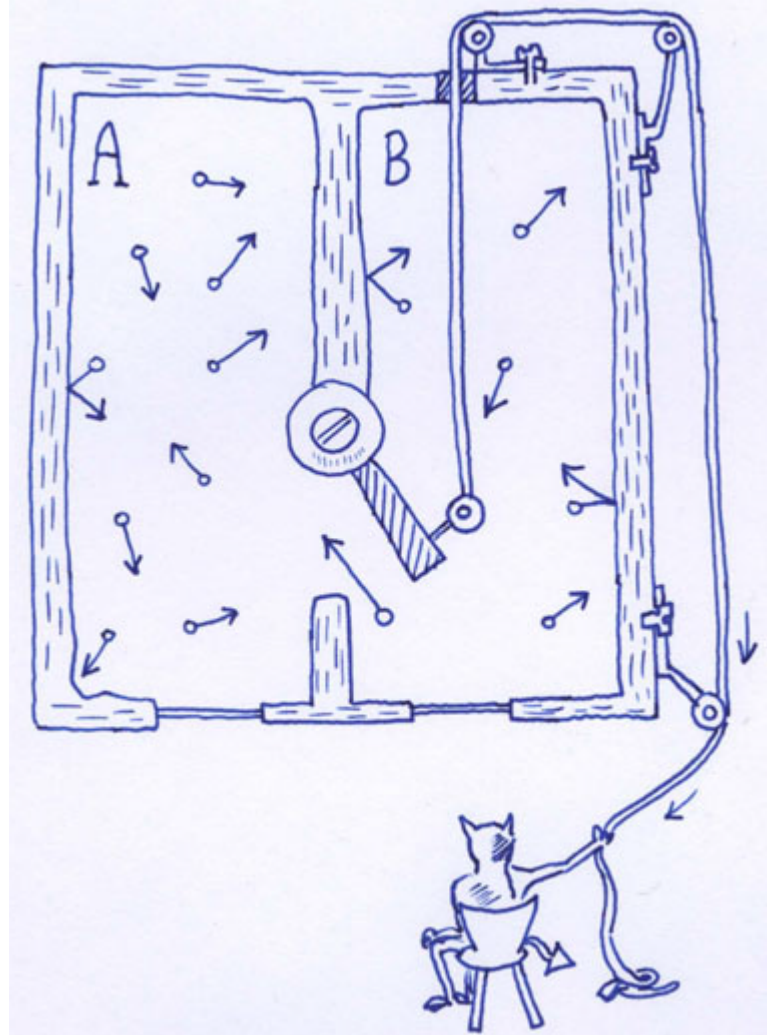


Autonomous Maxwell's Demon



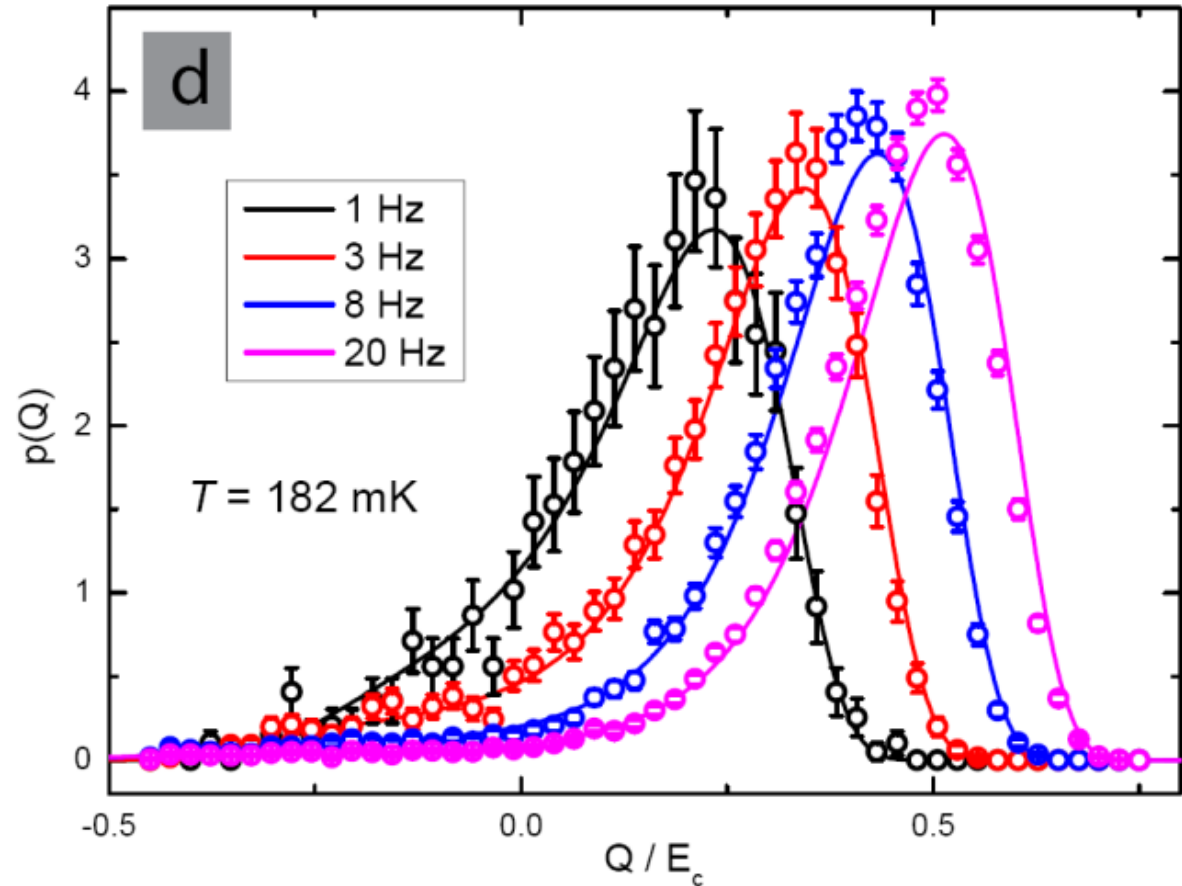
Maxwell's Demon

J. C. Maxwell 1867



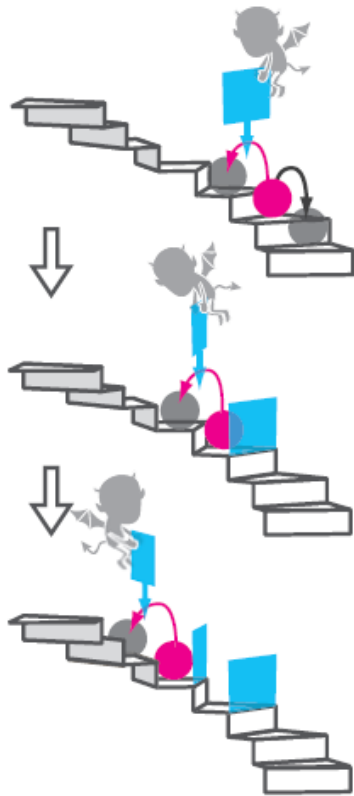
Negative heat

**Possible to
extract heat
from the
environment**

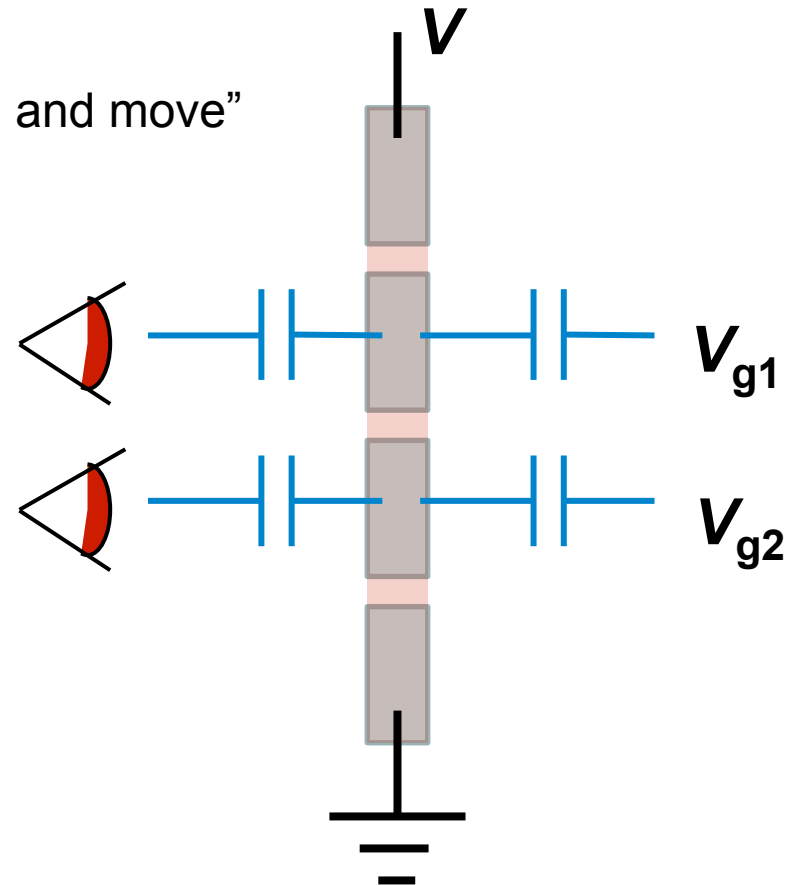


Provides means to realize Maxwell's demon using SETs

Electronic Maxwell's demon



"watch and move"



S. Toyabe et al., Nature Physics 2010

D. Averin et al., PRB 84, 245448 (2011).
G. Schaller et al., PRB 84, 085418 (2011).
P. Strassberg et al., PRL 110, 040601 (2013).
J. Bergli et al., Phys. Rev. E 88, 062139 (2013).

Szilard's engine (L. Szilard 1929)

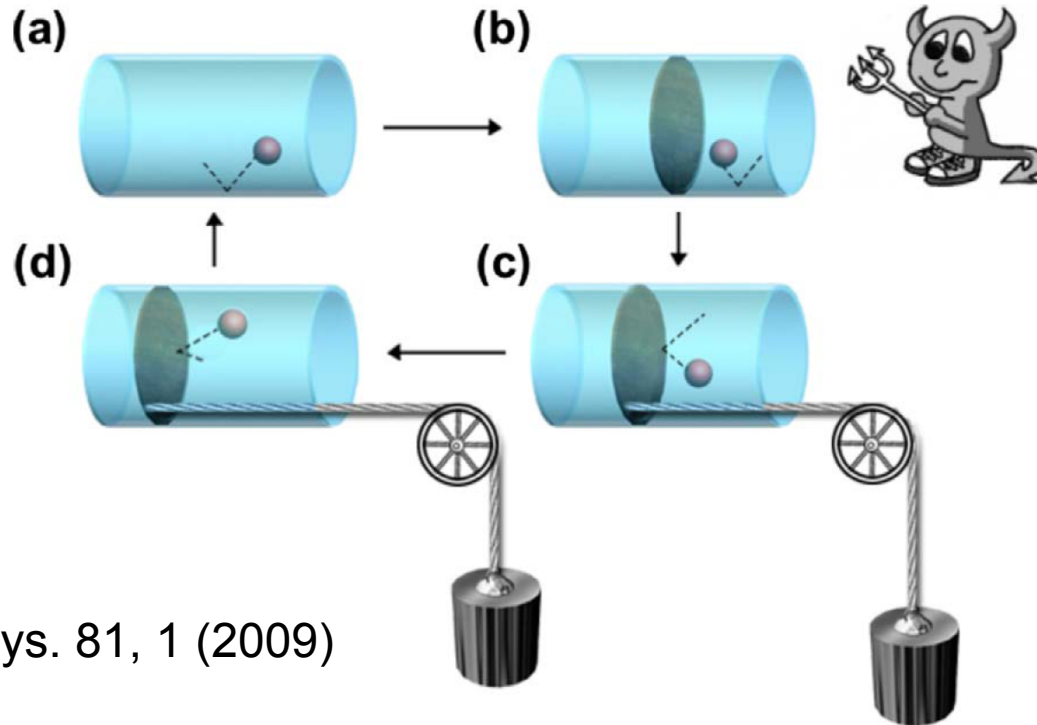


Figure from
Rev. Mod. Phys. 81, 1 (2009)

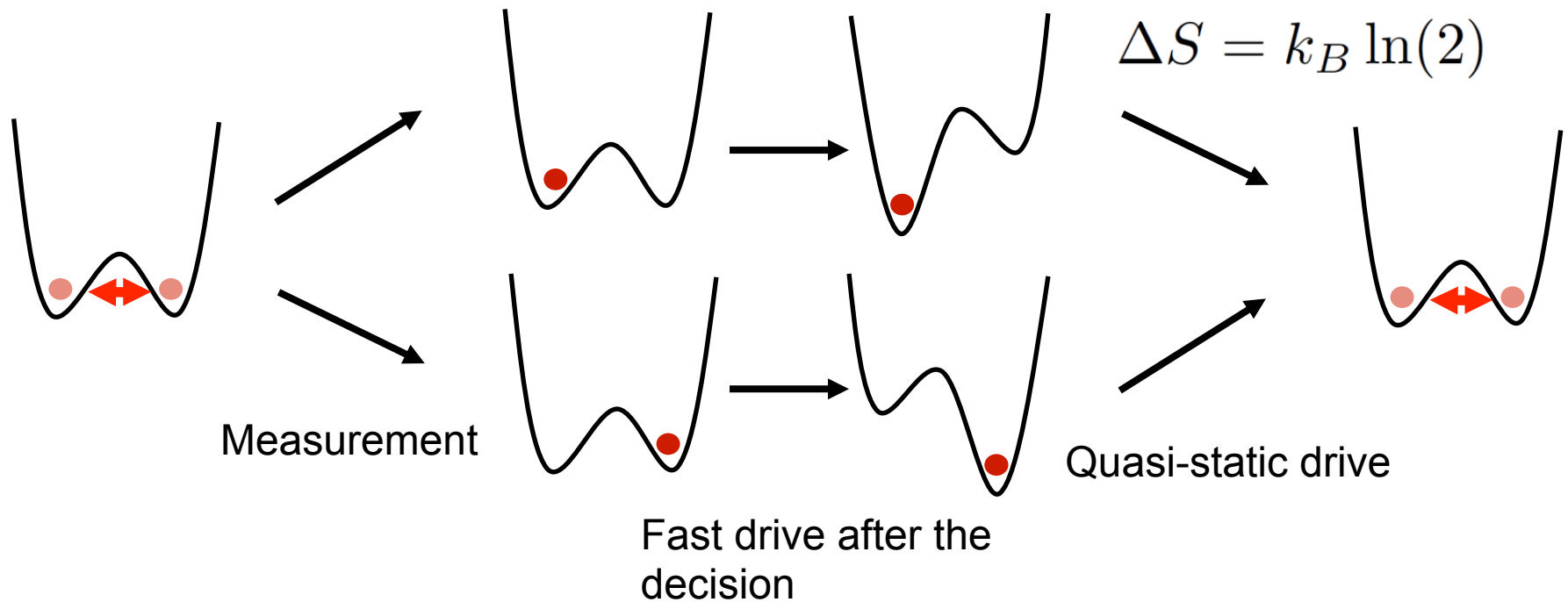
Isothermal expansion of the "single-molecule gas" does work against the load

$$W = Q = \int_{V/2}^V p dV = \int_{V/2}^V \frac{k_B T}{V} dV = k_B T \ln 2$$

Szilard's engine for single electrons

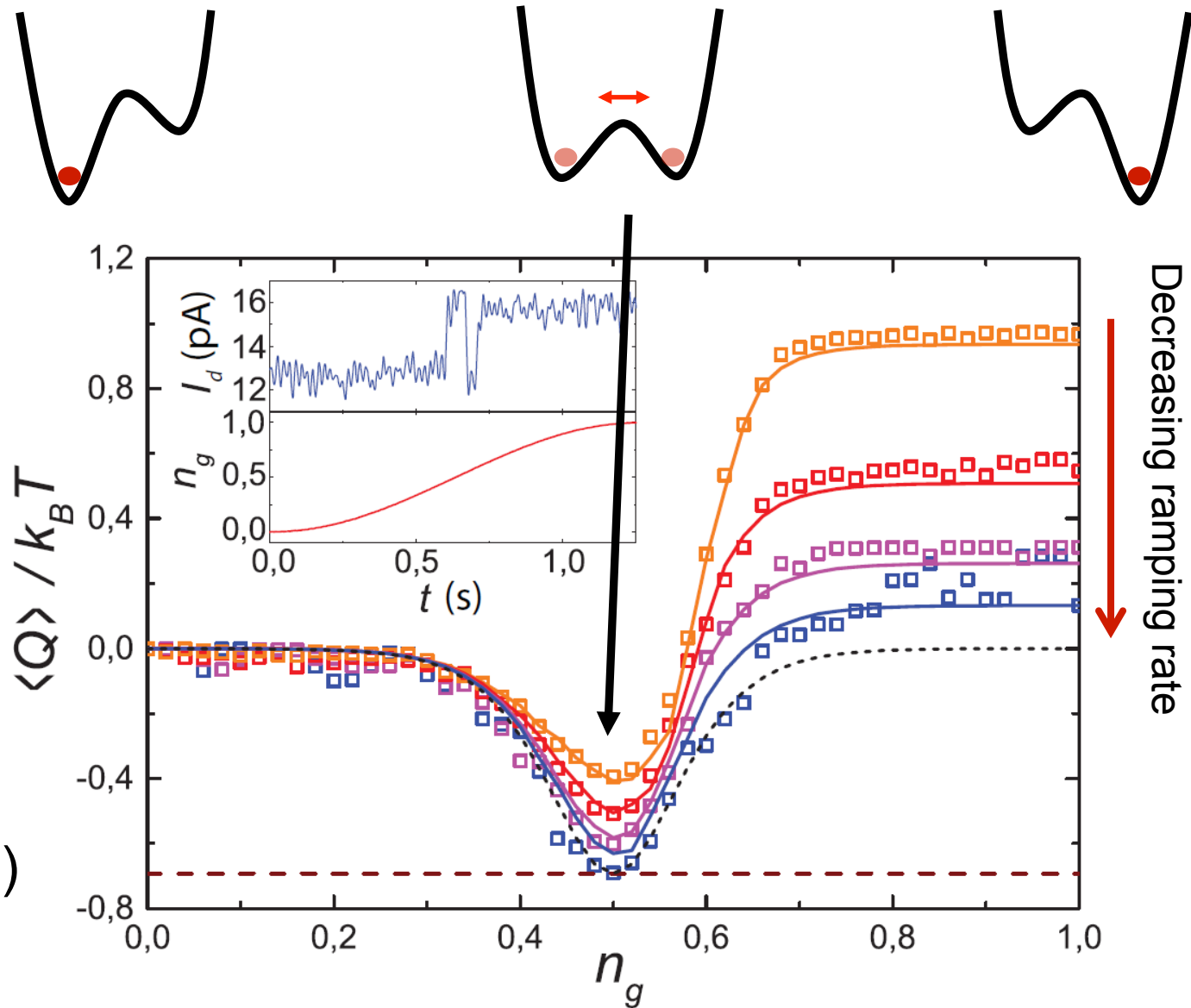
J. V. Koski et al., PNAS 111, 13786 (2014); PRL 113, 030601 (2014).

Entropy of the charge states: $S = -k_B \sum_{i=0,1} p(i) \ln[p(i)]$



In the full cycle (ideally): $Q = W = -k_B T \ln(2)$

Extracting heat from the bath

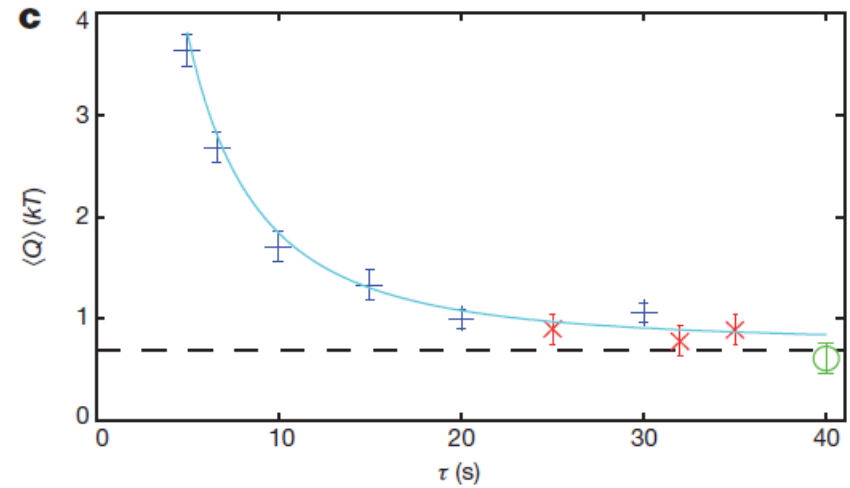
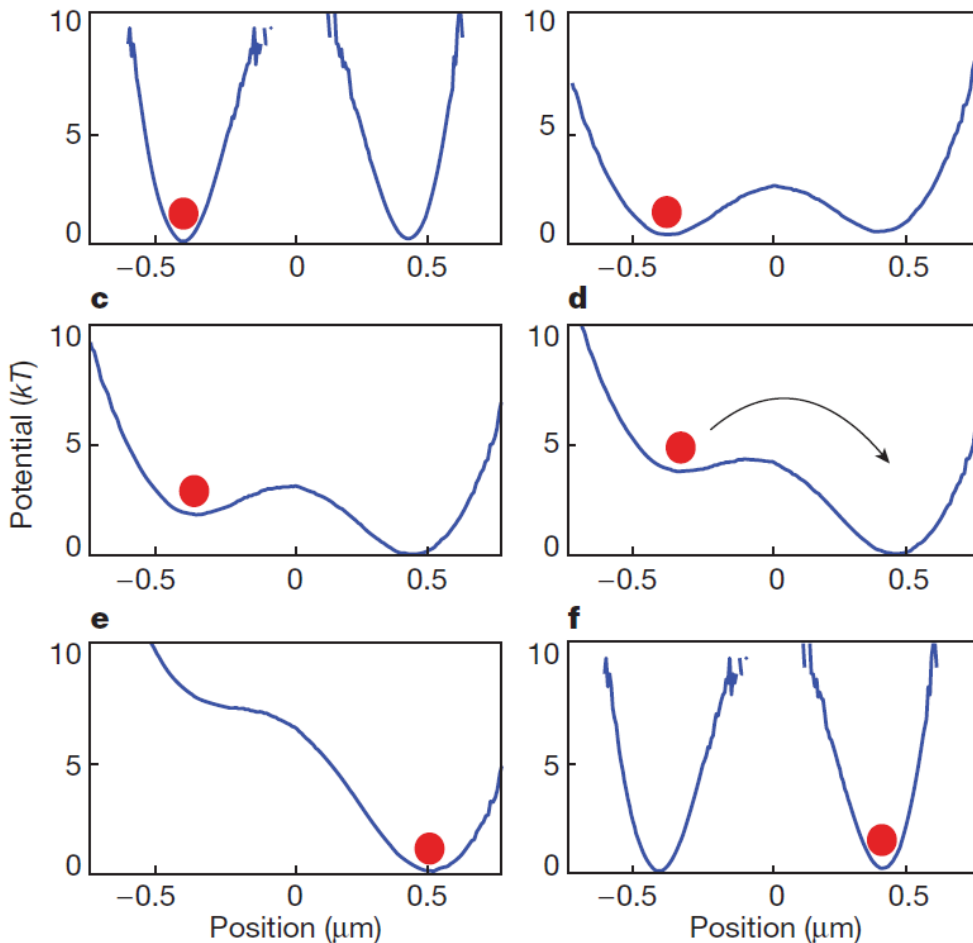


Erasure of information

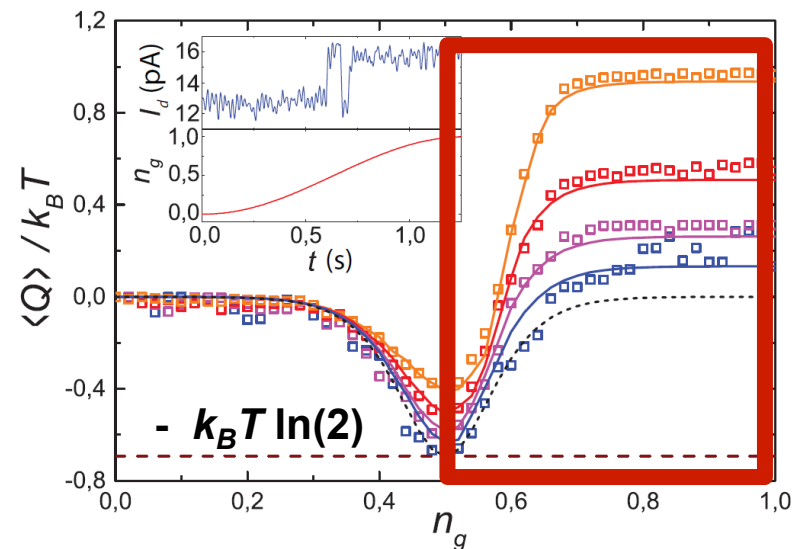
Landauer principle: erasure of a single bit costs energy of at least $k_B T \ln(2)$

Experiment on a colloidal particle:

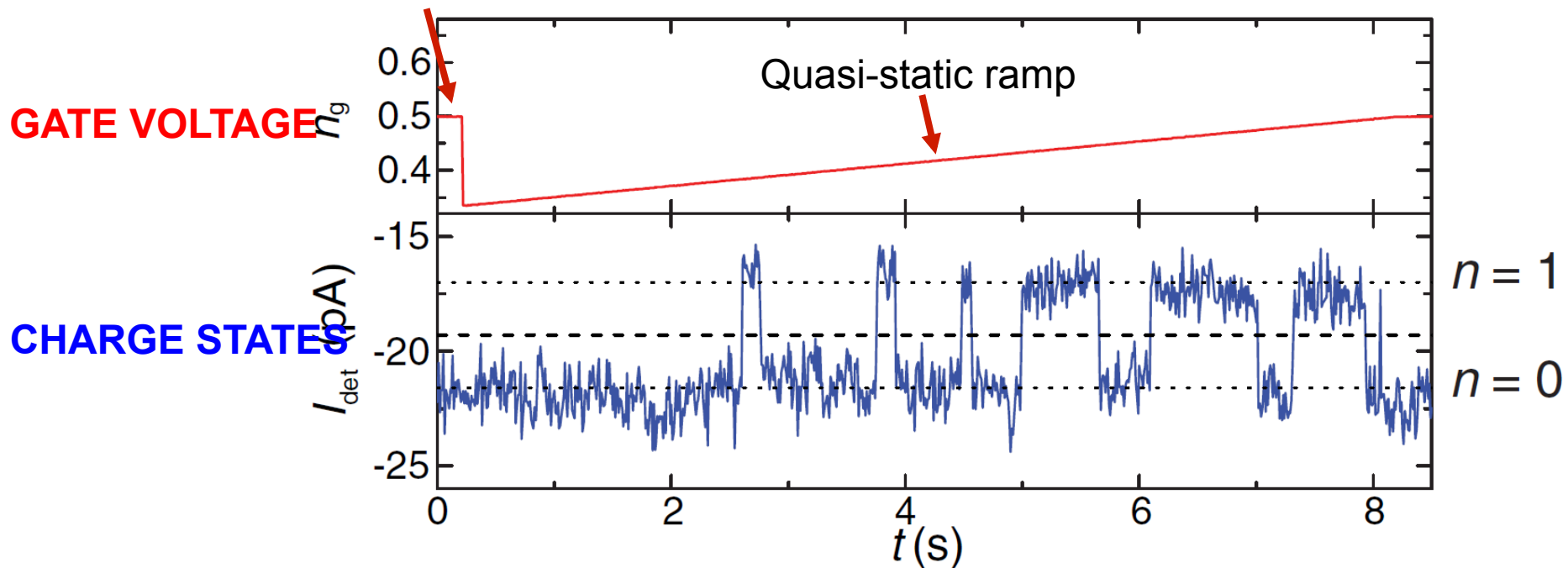
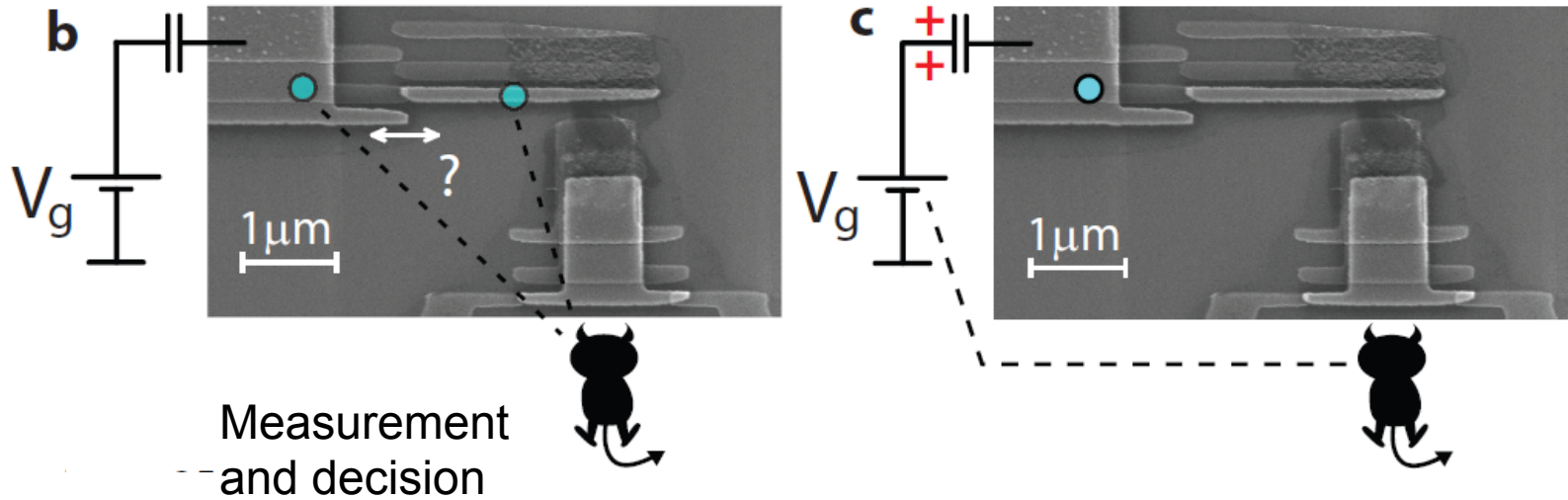
A. Berut, ... , S. Ciliberto et al., Nature 2012



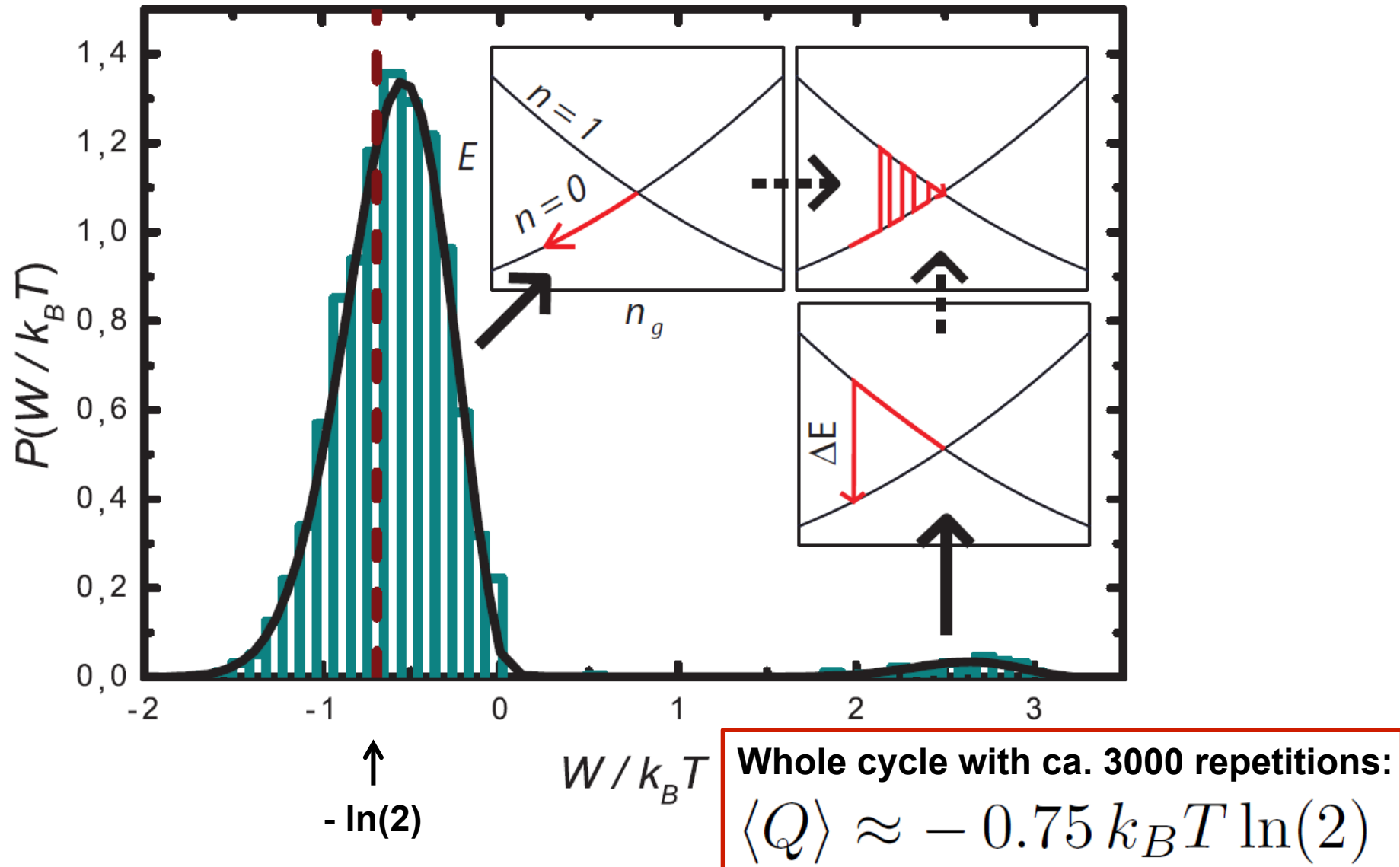
Corresponds to our experiment:



Realization of the Maxwell's demon with an electron



Measured distributions in the Szilard's engine



Sagawa-Ueda relation

$$\langle e^{-(W - \Delta F)/k_B T - I} \rangle = 1$$

T. Sagawa and M. Ueda, PRL 104, 090602 (2010)

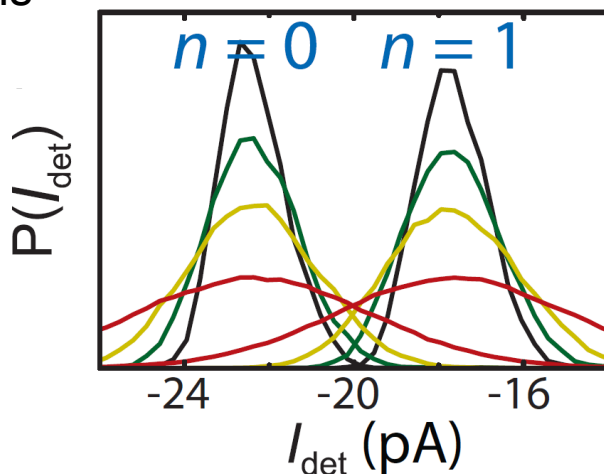
$$I(m, n) = \ln \left(\frac{P(n|m)}{P(n)} \right)$$

For a symmetric two-state system:

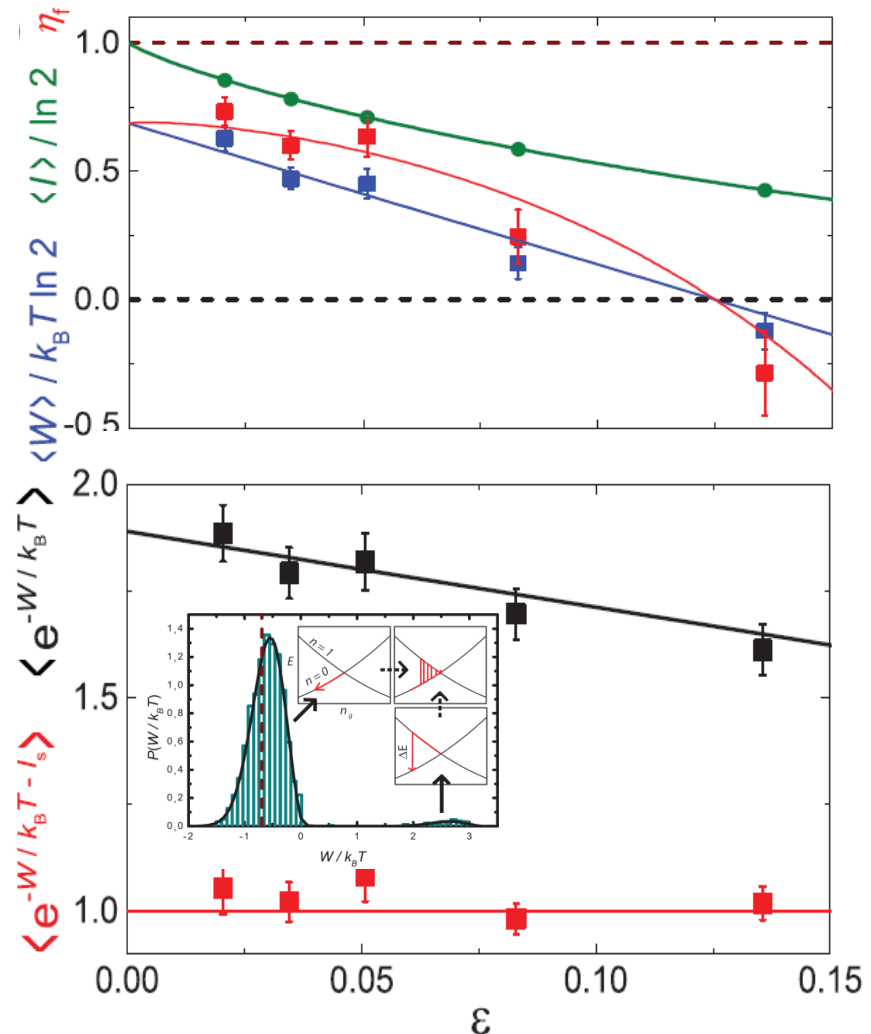
$$I(n = m) = \ln(2(1 - \epsilon))$$

$$I(n \neq m) = \ln(2\epsilon)$$

Measurements of n at different detector bandwidths

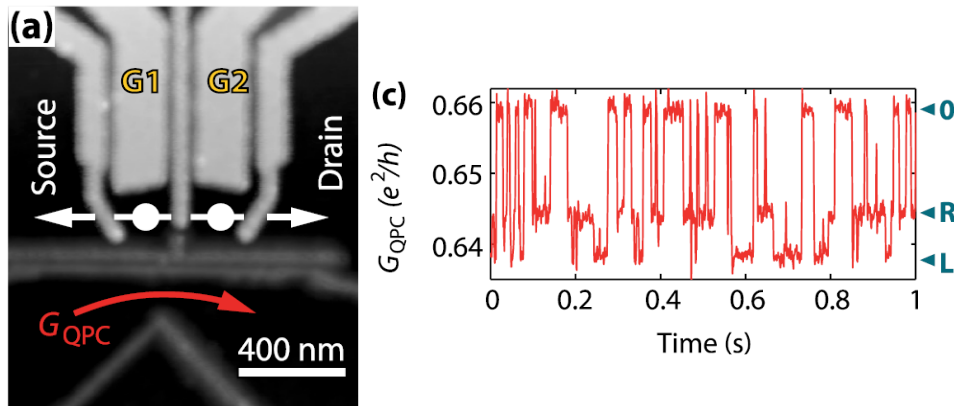


J. V. Koski et al., PRL 113, 030601 (2014)



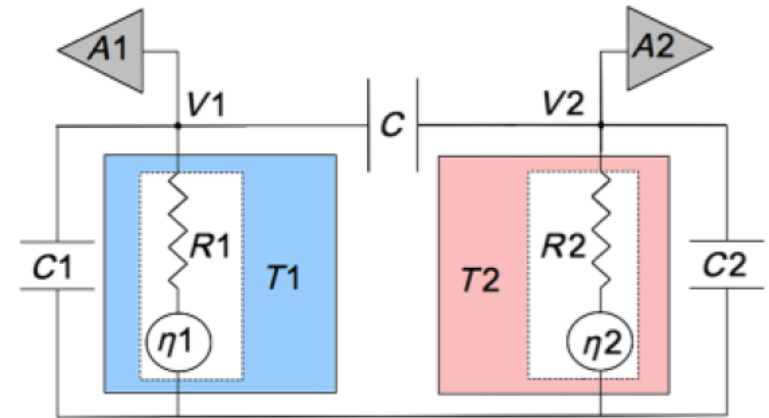
Work and heat in small systems: experimental situation

Typically an indirect measurement, hinging on understanding the system sufficiently well



$$Q = neV_{DS}$$

Y. Utsumi et al. PRB 81, 125331 (2010),
B. Kung et al. PRX 2, 011001 (2012)

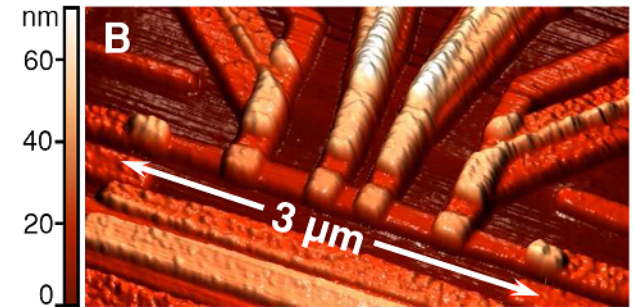
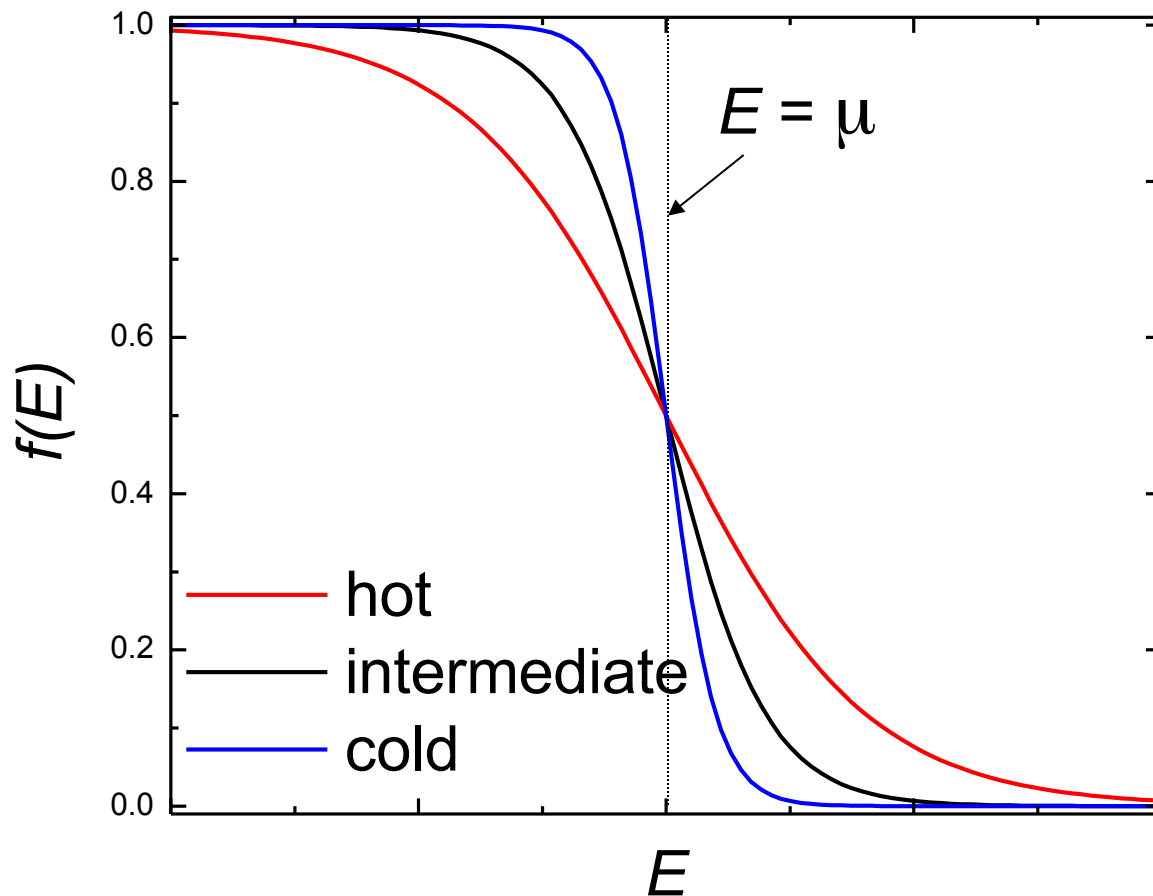


$$Q_{m,\tau} = \int_t^{t+\tau} i_m V_m dt$$

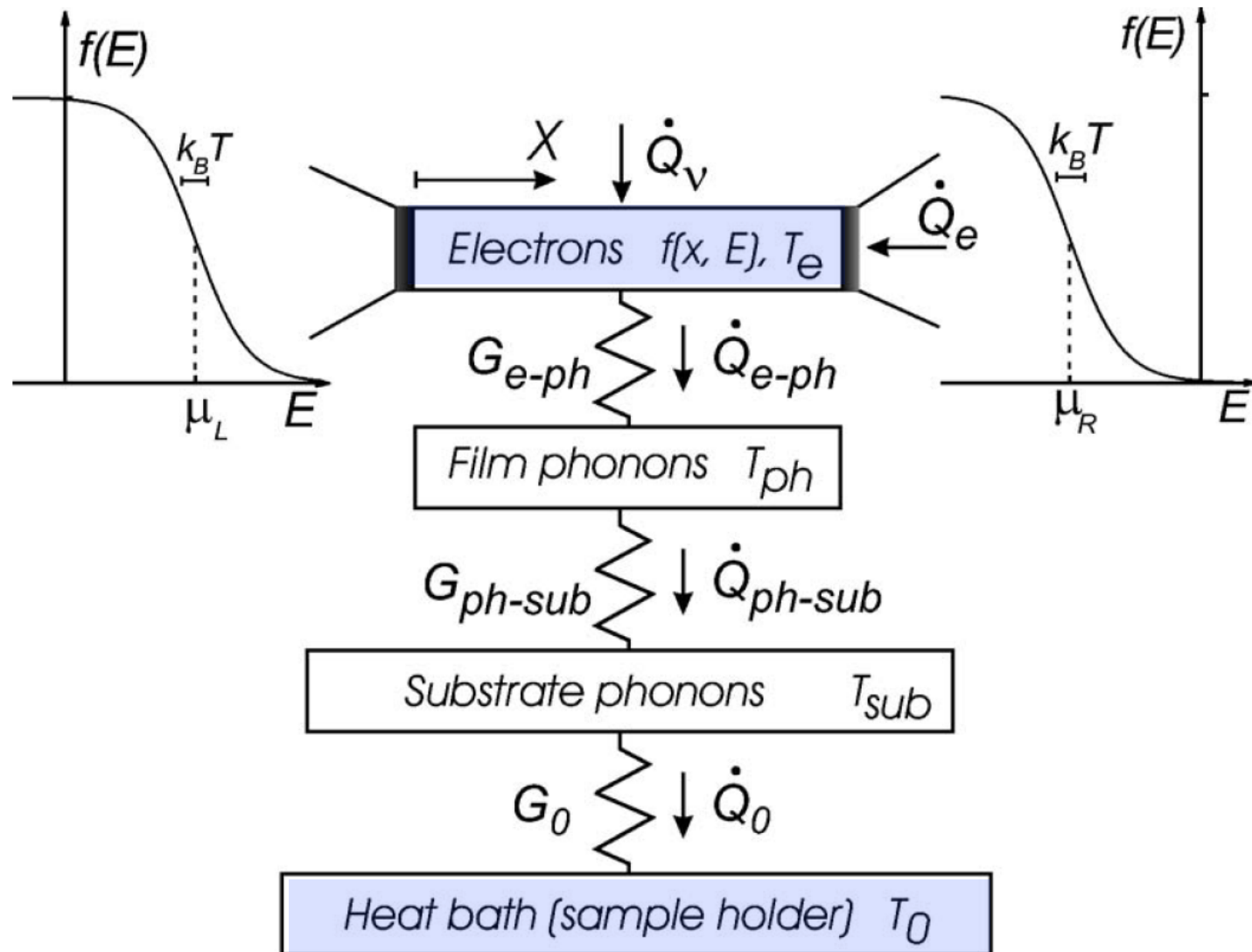
S. Ciliberto et al., PRL 110, 180601 (2013)

Temperature in an electronic device

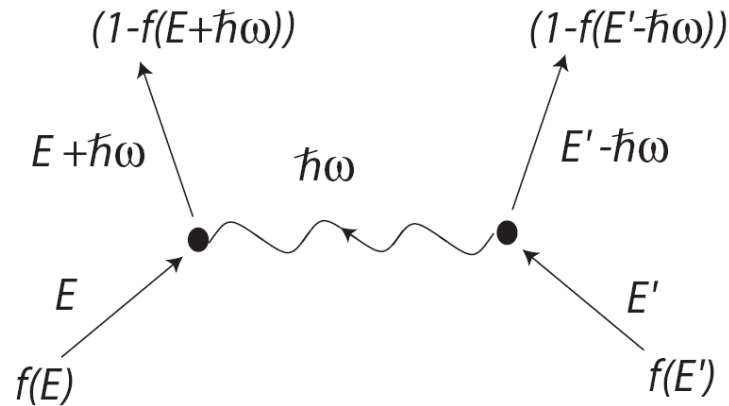
$$f(E) = \frac{1}{1 + e^{(E-\mu)/k_B T}}$$



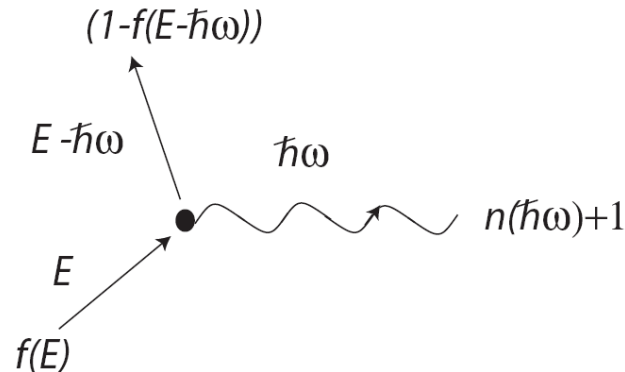
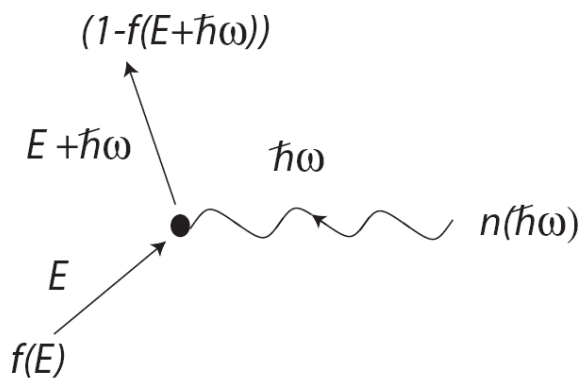
Generic thermal model for an electronic conductor



Electron-electron and electron-phonon relaxation



e-e relaxation drives
the system towards
quasi-equilibrium



e-p relaxation drives
the system towards
equilibrium

The energy distribution of electrons in a small metal conductor

The distribution is determined by energy relaxation:

Equilibrium – Thermometer measures the temperature of the "bath"

Quasi-equilibrium – Thermometer measures the temperature of the electron system which can be different from that of the "bath"

Non-equilibrium – There is no well defined temperature measured by the "thermometer"

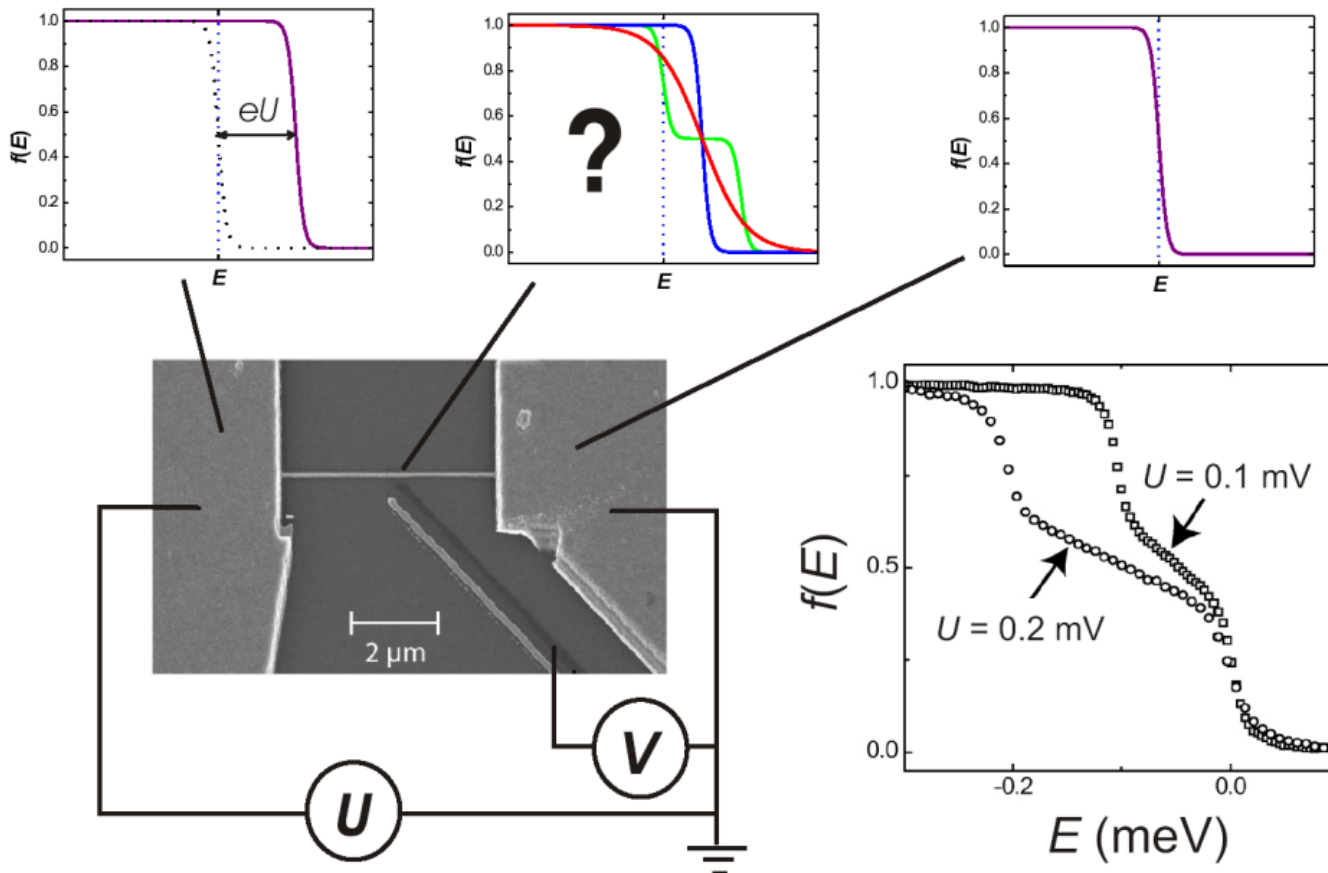


Illustration: diffusive normal metal wire
H. Pothier et al. 1997

$$\begin{aligned}\tau_{ee} &= 1\ \text{ns} \\ \tau_{ep} &= 10\ \mu\text{s}\end{aligned}$$

Electron-phonon relaxation in metals at low T

PHYSICAL REVIEW B

VOLUME 49, NUMBER 9

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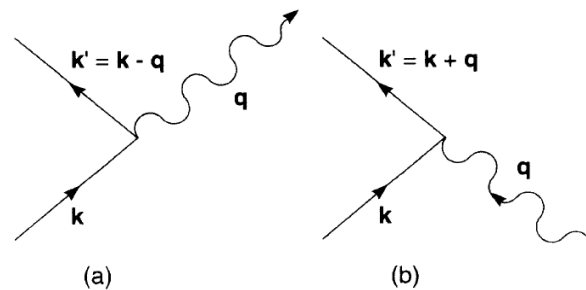
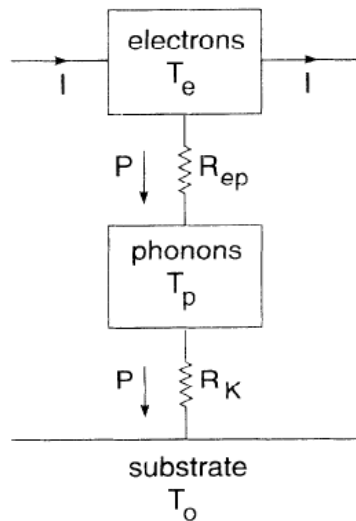
Hot-electron effects in metals

F. C. Wellstood,^{*} C. Urbina,[†] and John Clarke

Department of Physics, University of California, Berkeley, California 94720

and Materials Sciences Division, Lawrence Berkeley Laboratory, Berkeley, California 94720

(Received 21 July 1993)

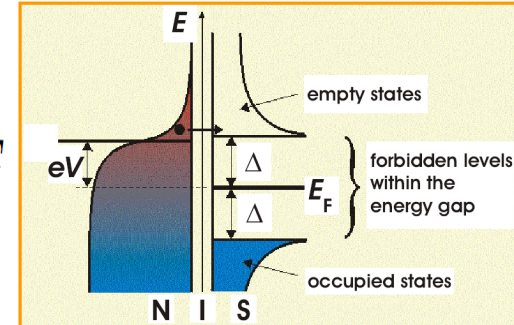


$$\dot{Q}_{ep} = \Sigma \Omega (T_e^5 - T_p^5)$$

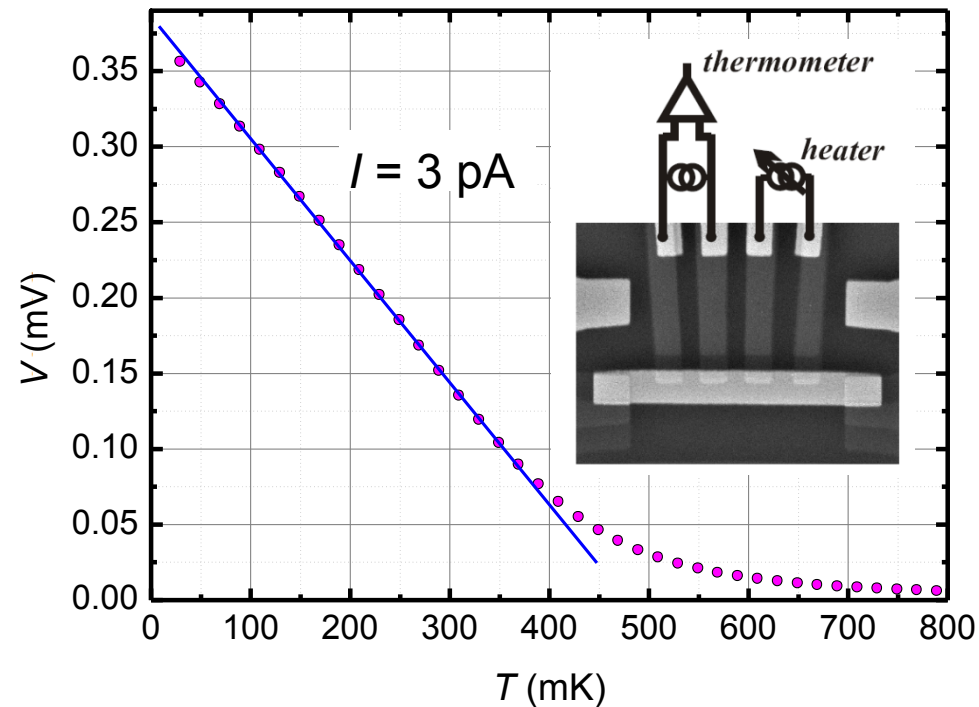
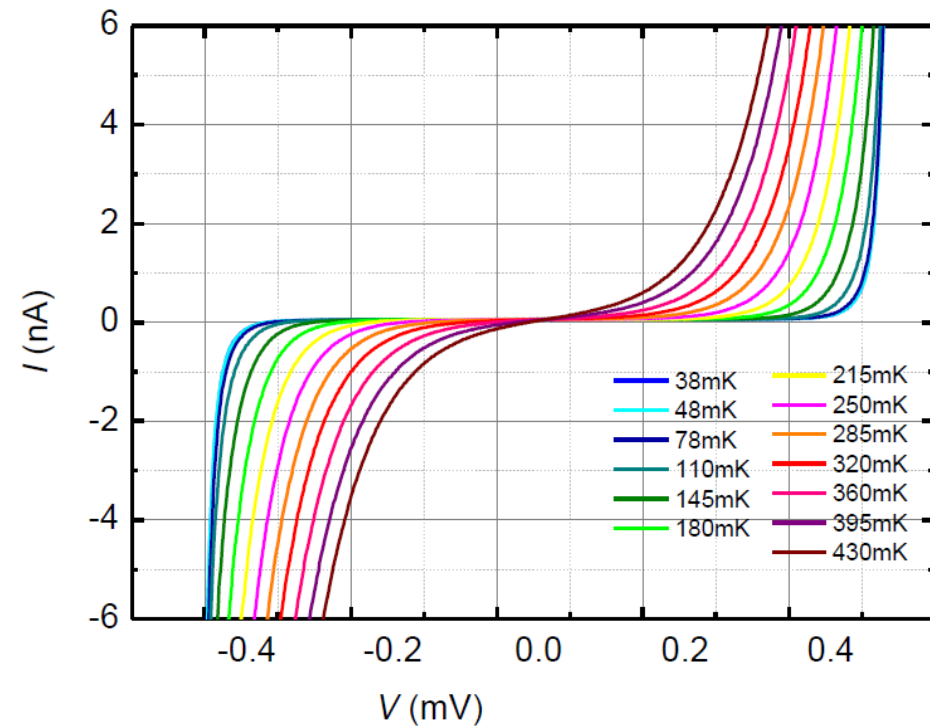
FIG. 2. Emission and absorption of phonons of wave vector q by an electron of wave vector k .

NIS-thermometry

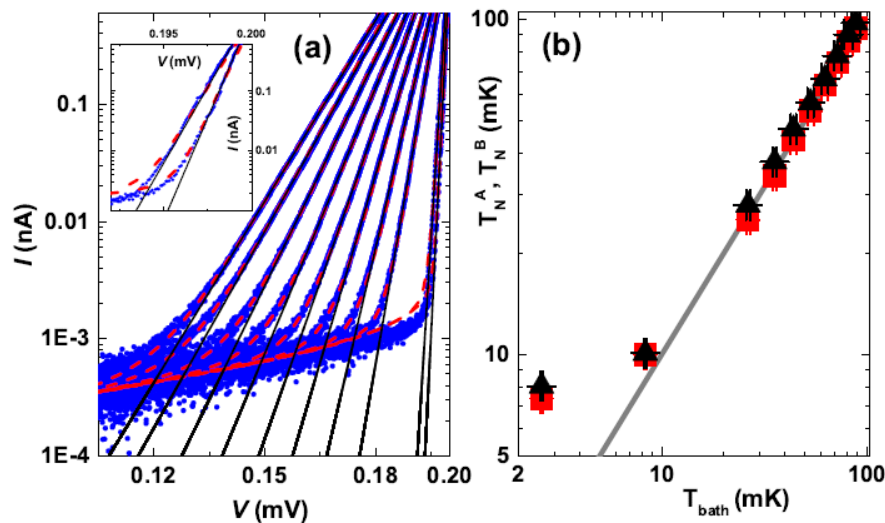
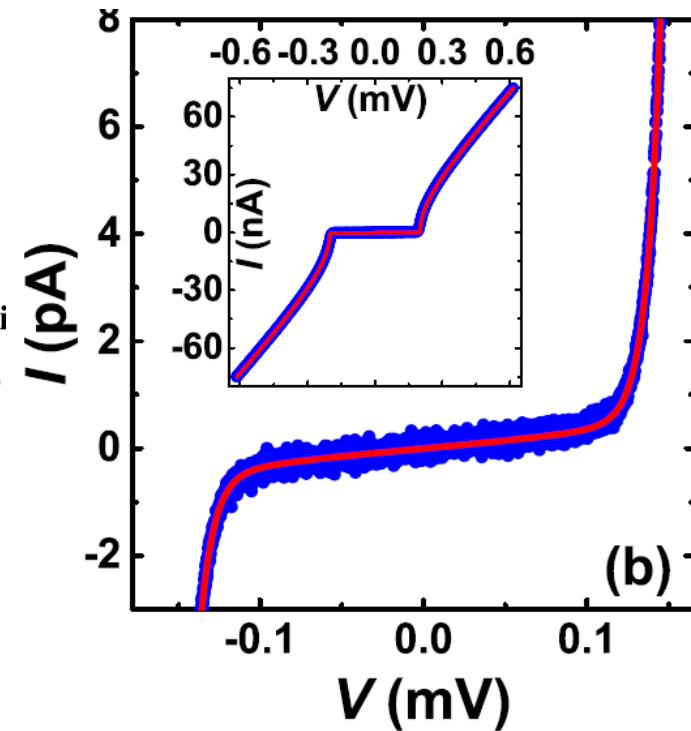
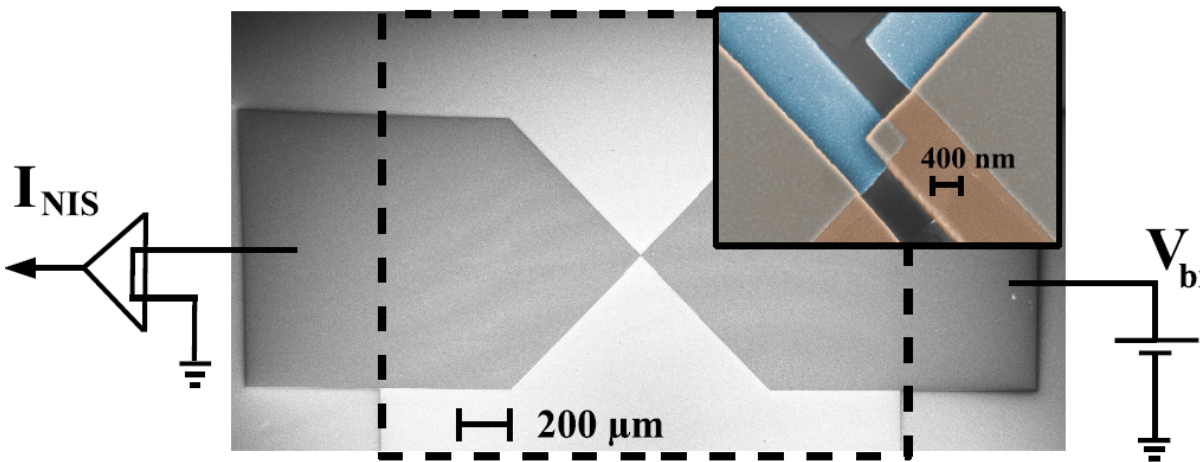
$$I = \frac{1}{2eR_T} \int n_S(E) [f_N(E - eV) - f_N(E + eV)] dE$$



Probes electron temperature of N island (and not of S!)



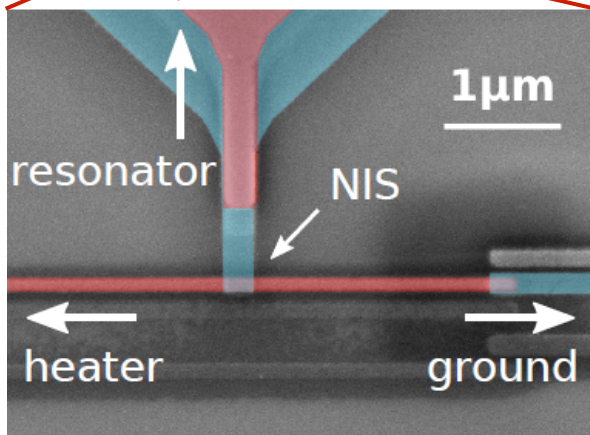
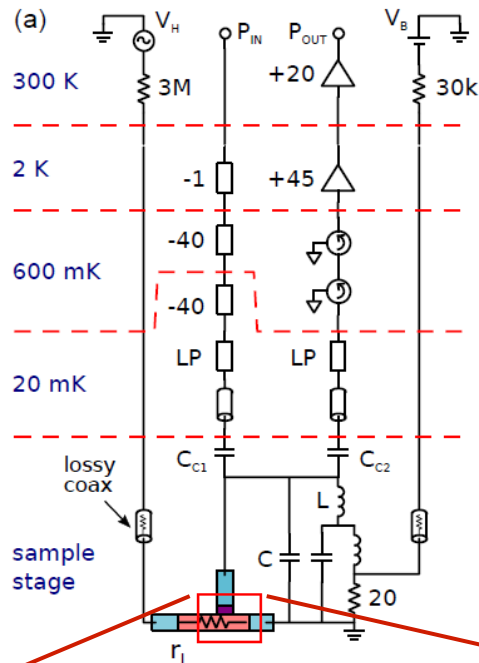
NIS-thermometry at mK temperatures



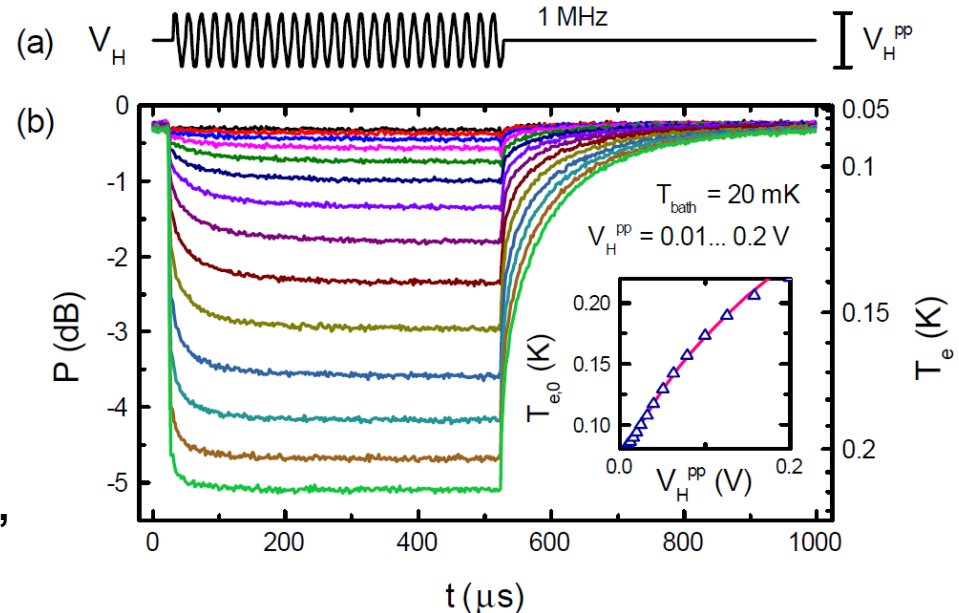
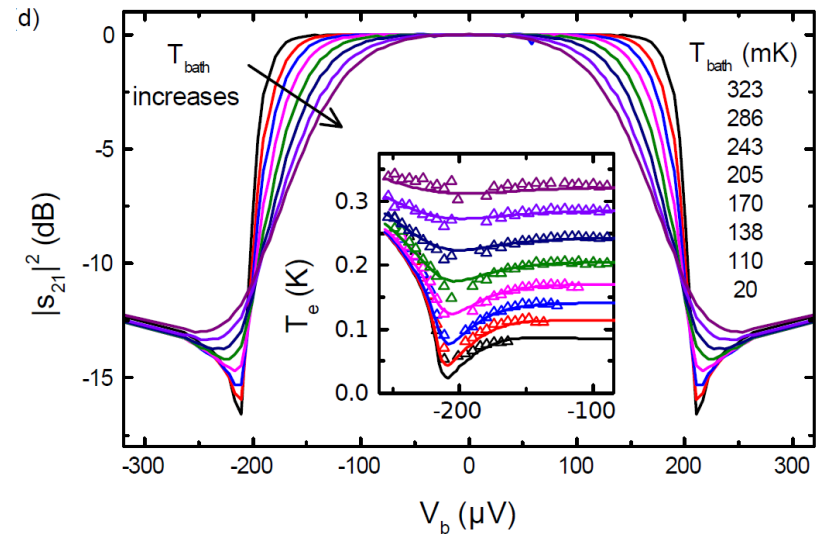
A. V. Feshchenko et al. (Aalto - Basel Univ. collaboration), Phys. Rev. Appl. 4, 034001 (2015).

Potential to operate down to 1 mK

Fast thermometry



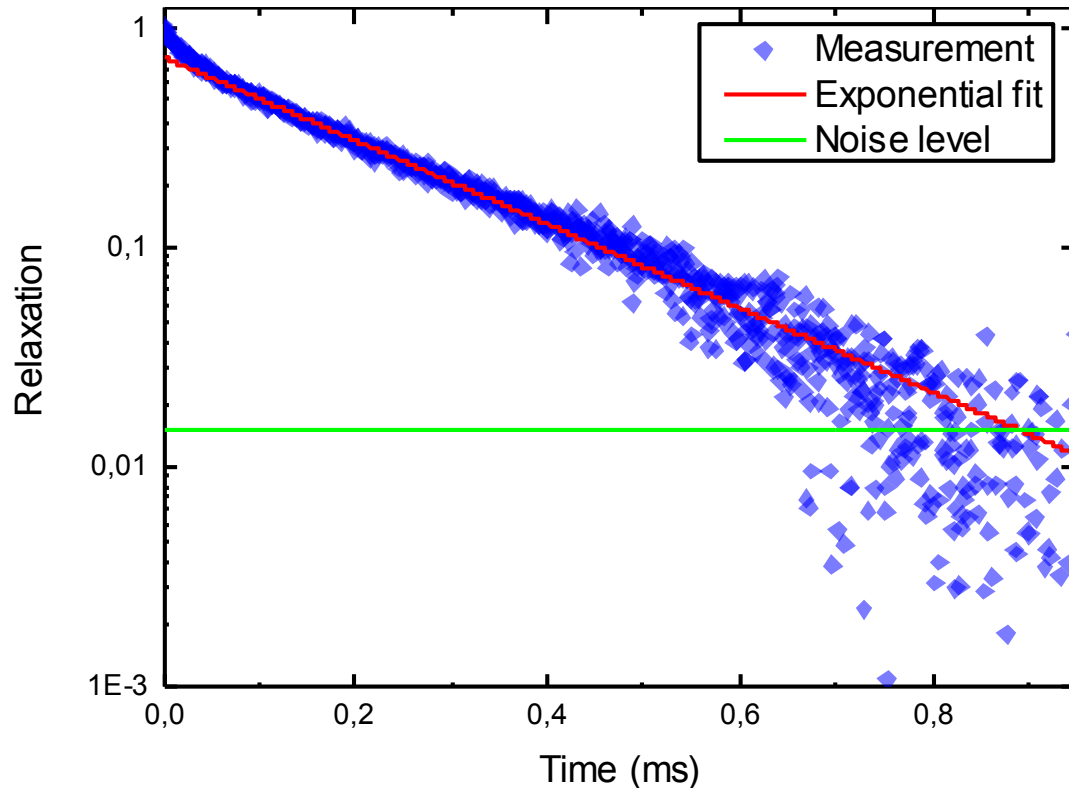
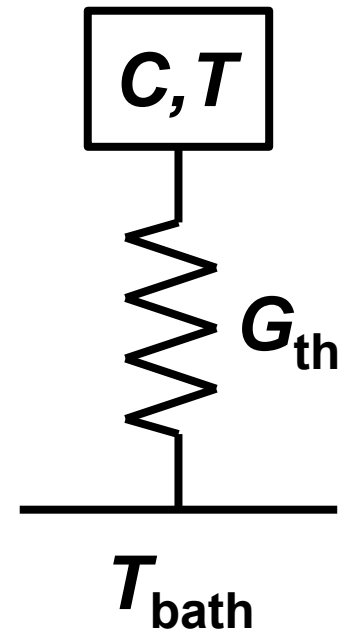
Transmission read-out at 600 MHz of a NIS junction



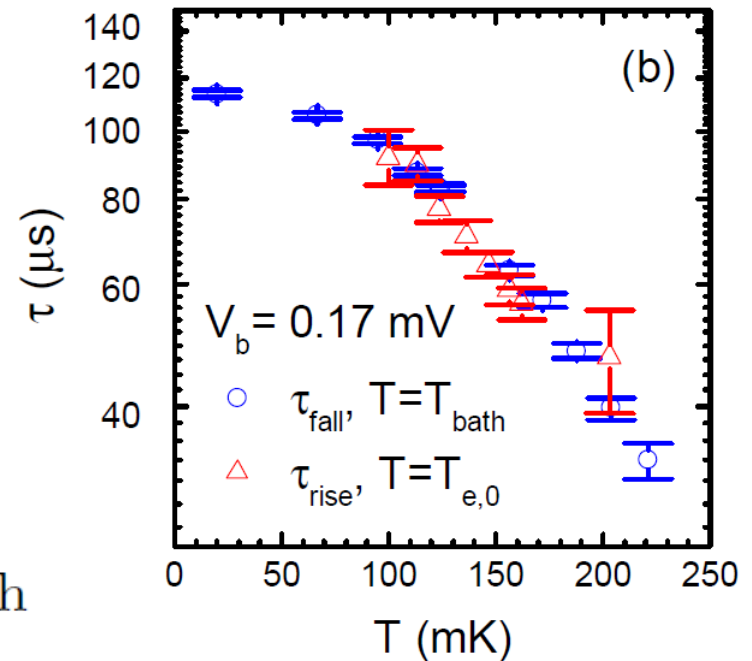
S. Gasparinetti et al., Phys. Rev. Applied 3, 014007 (2015); K. L. Viisanen et al., New J. Phys. 17, 055014 (2015).

Fast thermometry

Time-resolved measurements

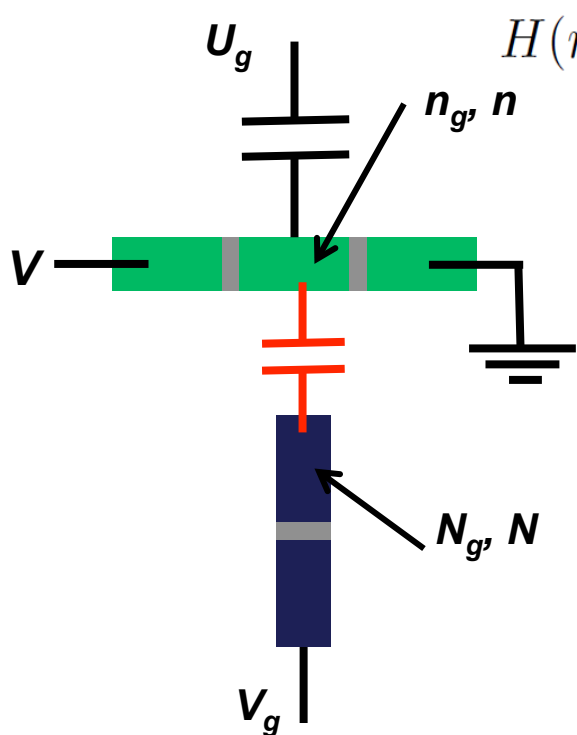


$$C \frac{d\delta T}{dt} = -G_{th} \delta T \quad \tau = C/G_{th}$$

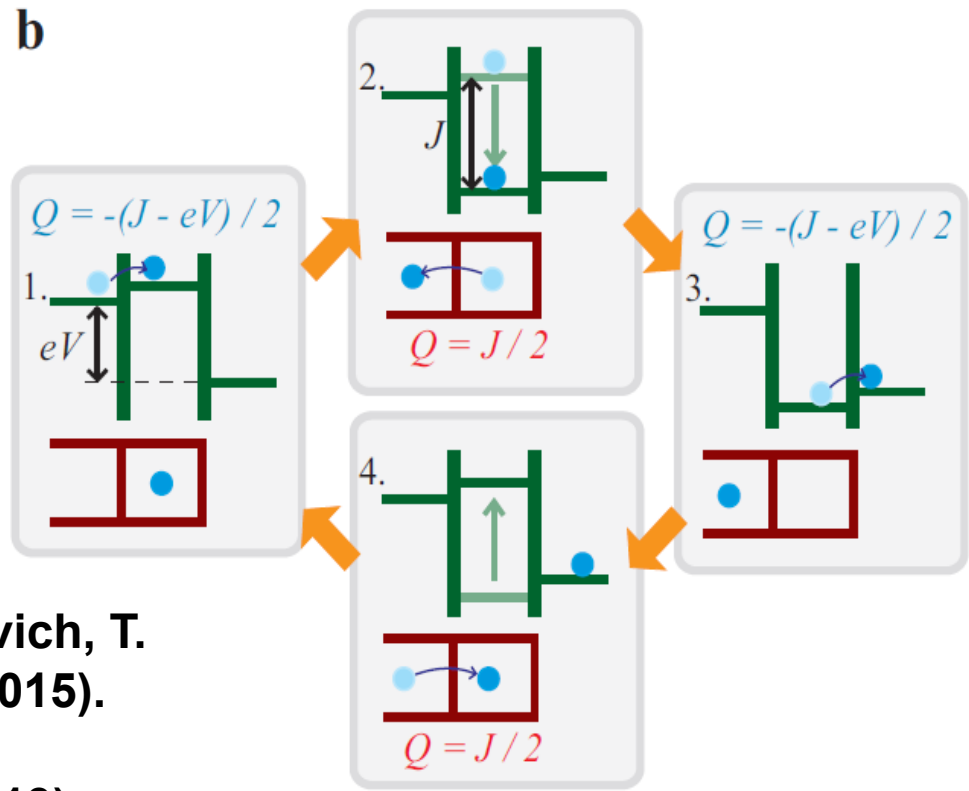


Autonomous Maxwell's demon

System and Demon: all in one
Realization in a circuit:



$$H(n, N) = E_s(n - n_g)^2 + E_d(N - N_g)^2 + 2J(n - n_g)(N - N_g)$$

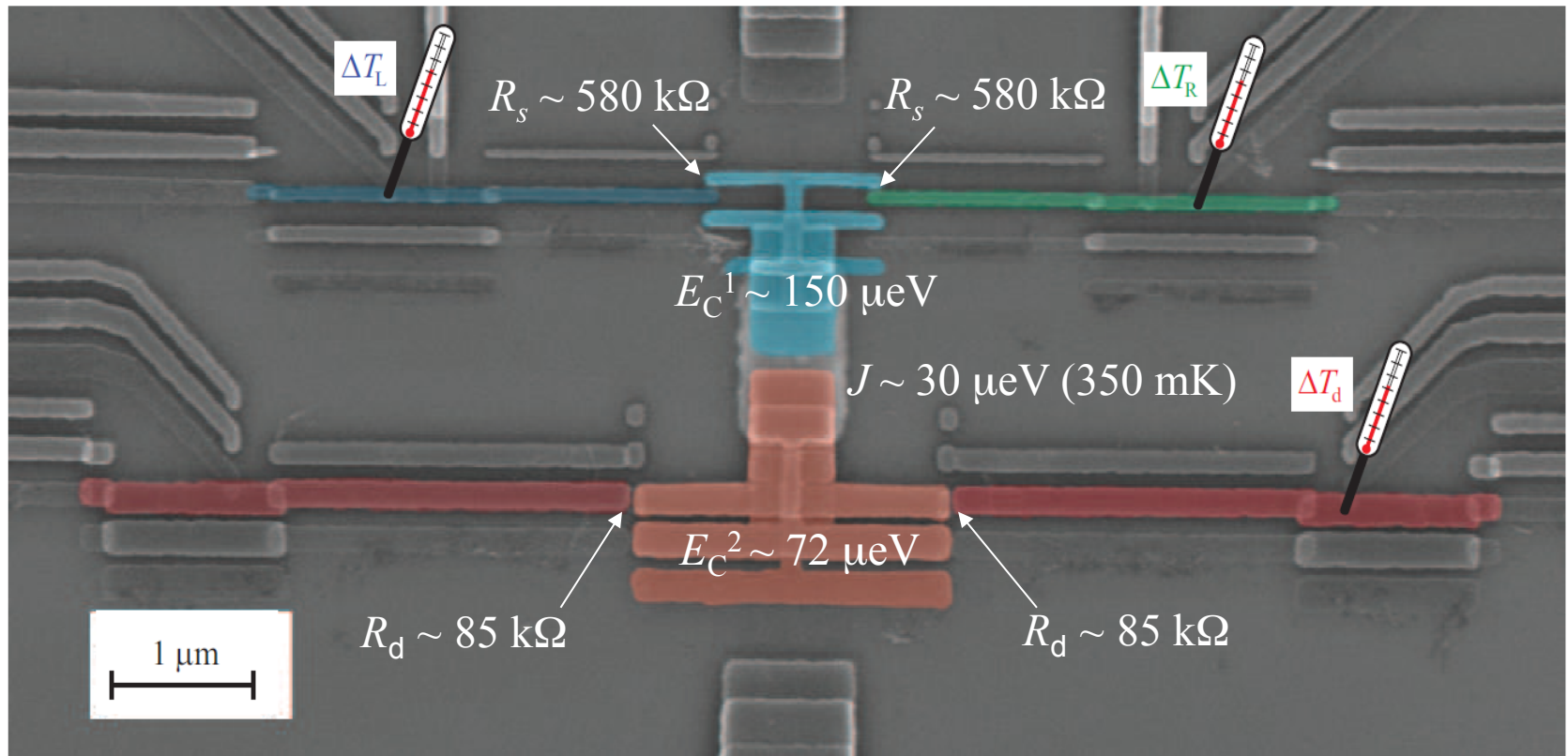


J. V. Koski, A. Kutvonen, I. M. Khaymovich, T. Ala-Nissila, and JP, PRL 115, 260602 (2015).

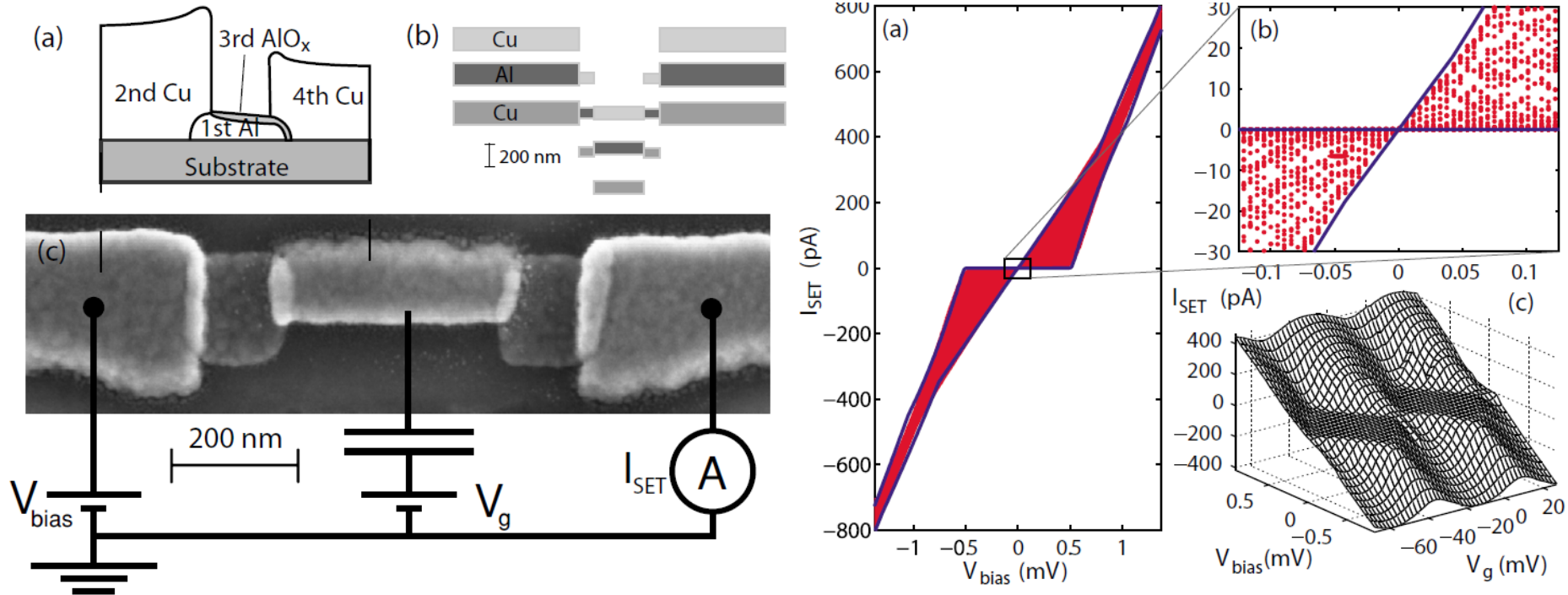
P. Strasberg et al., PRL 110, 040601 (2013).

Autonomous Maxwell's demon – information-powered refrigerator

Image of the actual device



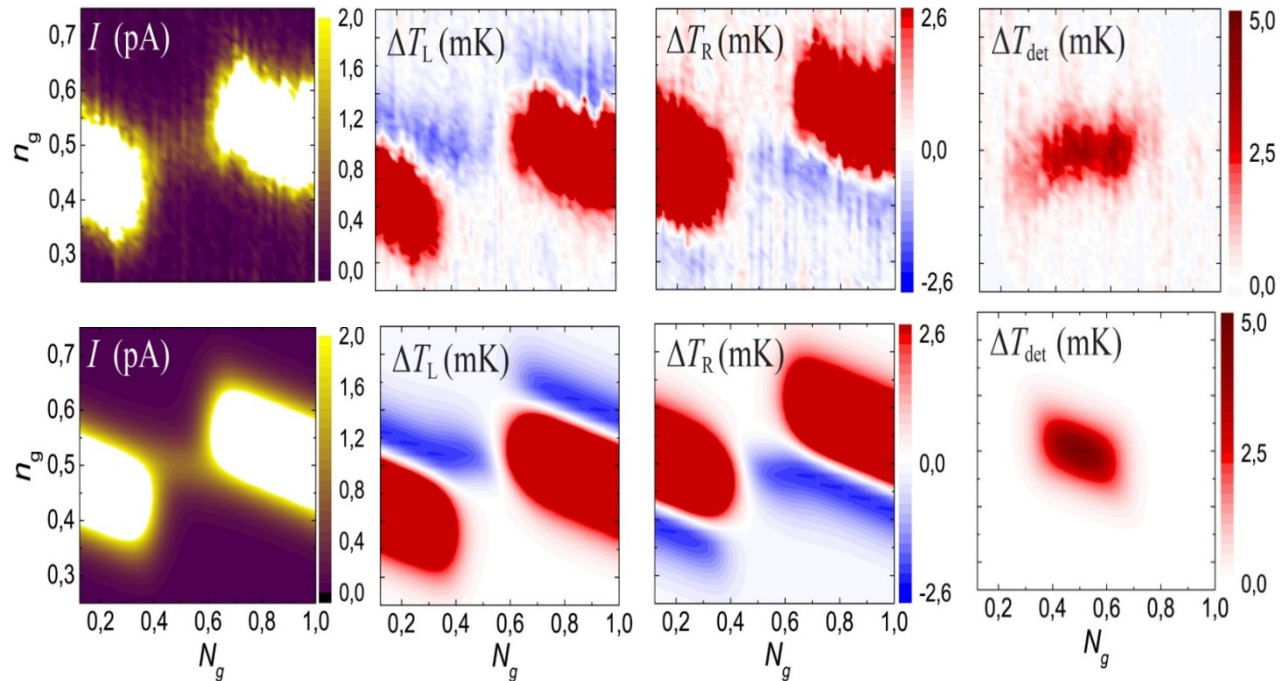
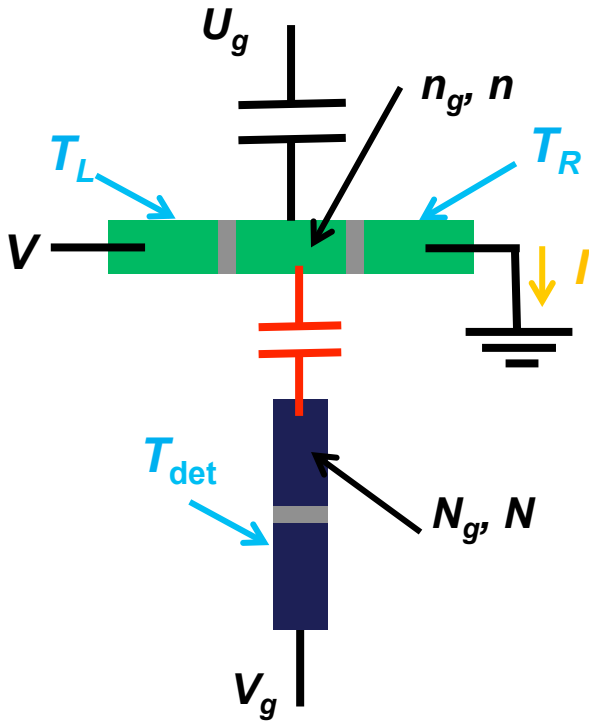
Technique to produce fully normal tunnel junctions and SETs



Laterally proximized tunnel junctions:

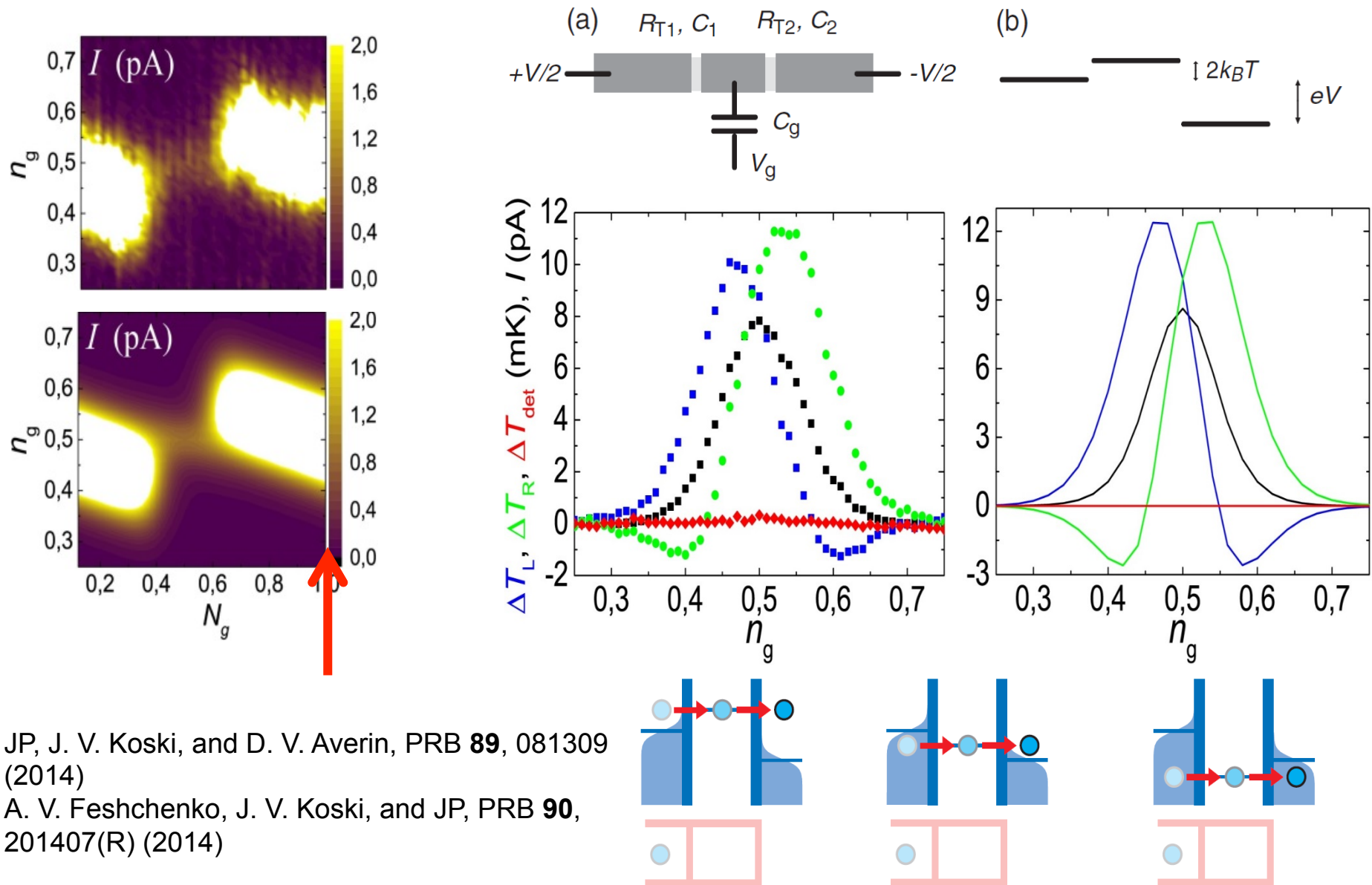
J. Koski et al., Appl. Phys. Lett. **98**, 203501 (2011).

Current and temperatures at different gate positions



$V = 20 \mu\text{V}$, $T = 50 \text{ mK}$

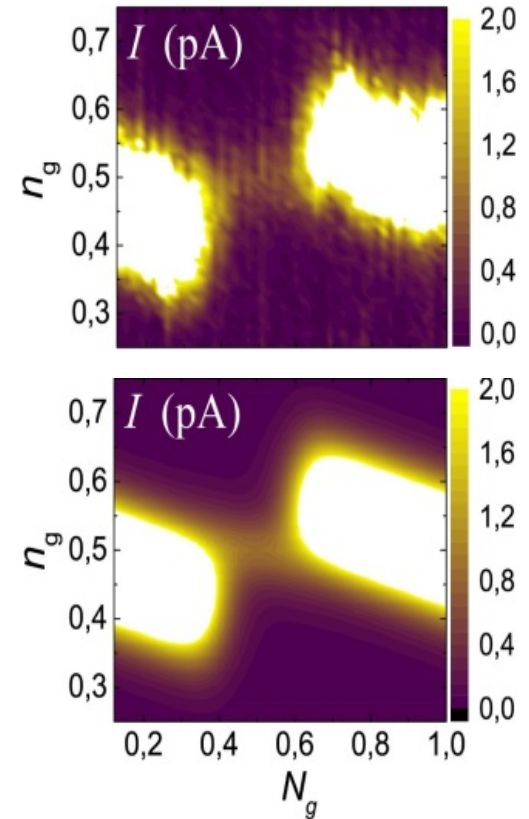
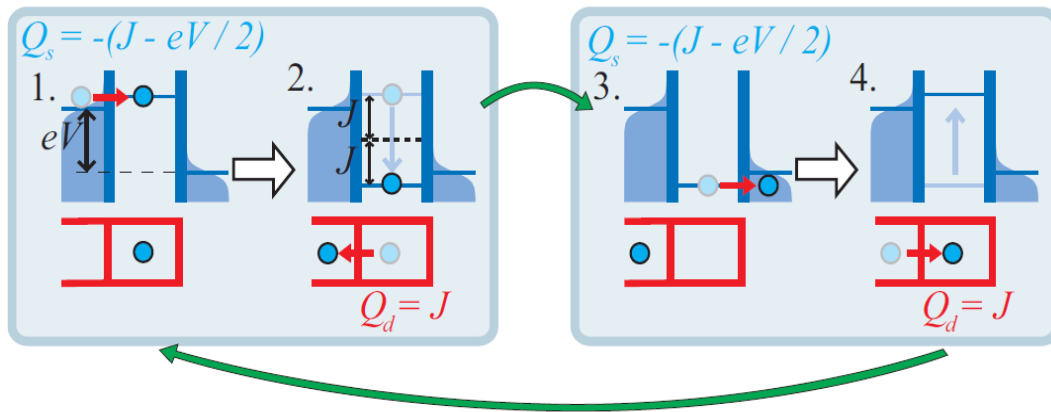
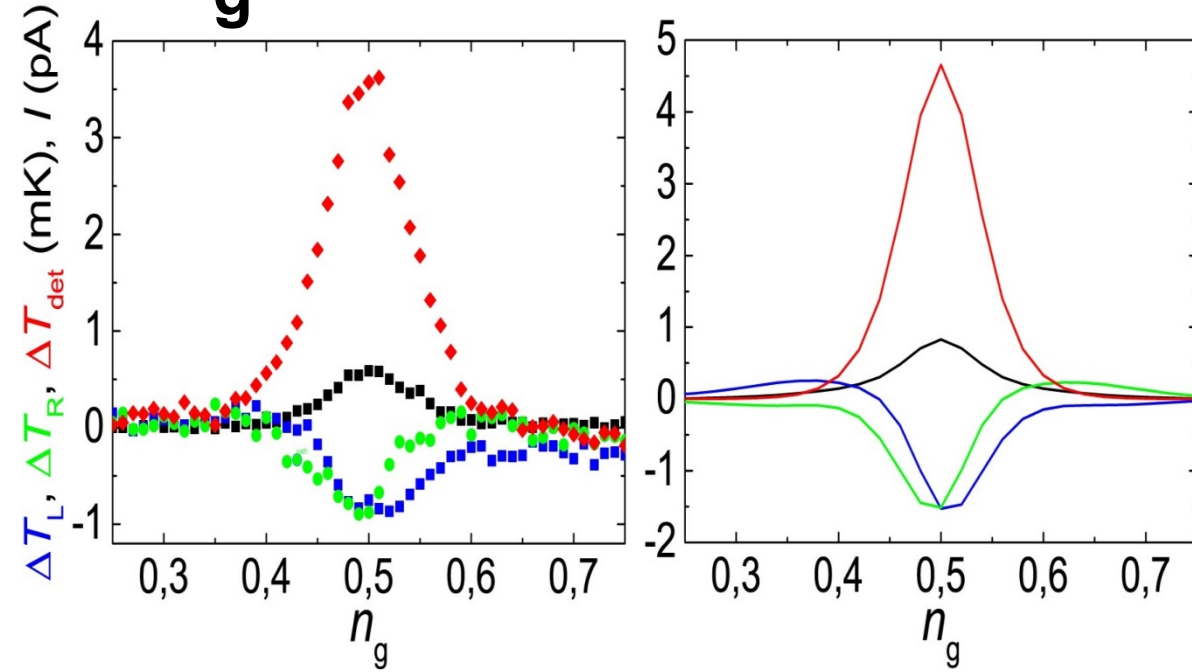
$N_g = 1$: No feedback control ("SET-cooler")



JP, J. V. Koski, and D. V. Averin, PRB **89**, 081309 (2014)

A. V. Feshchenko, J. V. Koski, and JP, PRB **90**, 201407(R) (2014)

$N_g = 0.5$: feedback control (Demon)

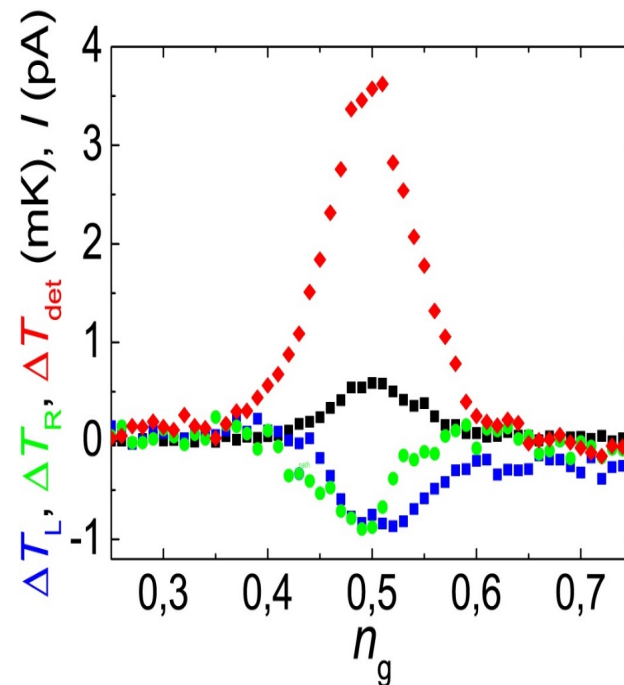
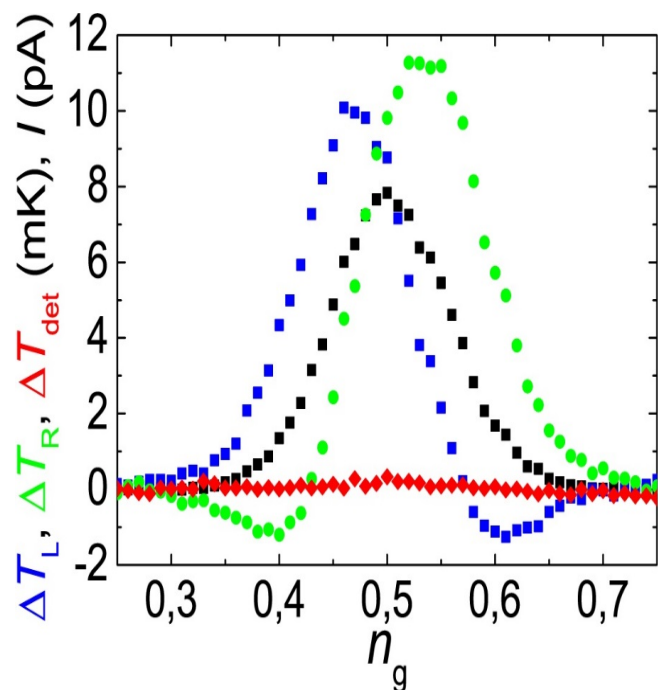


**Both T_L and T_R drop: entropy of the System decreases;
 T_{det} increases: entropy of the Demon increases**

Summary of the autonomous

SET cooler demon experiment

Demon



current I $\rightarrow$

