

Majoranas in semiconductor nanowires

Leo Kouwenhoven

Önder Gül, Hao Zhang, Michiel de Moor, Fokko de Vries, Jasper van Veen
David van Woerkom, Kun Zuo, Vincent Mourik, Srijit Goswami,
Maja Cassidy, Attila Geresdi

wires: Diana Car, Sébastien Plissard & Erik Bakkers

materials: S. Conesa Boj, M. Quintero-Perez, K. Watanabe, T. Taniguchi

theory: Michał Nowak, Michael Wimmer, Anton Akhmerov



INTRO

Majorana Fermion

“a particle that is equal to its own anti-particle”

$$\gamma^\dagger = \gamma$$

creating a Majorana is the same operation
as annihilating it:

meaning zero energy, zero
spin and zero charge.

okay for Bosons but not for Fermions
(electron \neq positron)

*How to detect a particle
with everything zero?*



1938

(Wilczek, *Majorana Returns*, Nature Physics, 2009)

Wilczek, *Majorana modes materialize*, Nature, News&Views, 2012)

Ettore
Rom
'59"

ANSA^{en} Science & Technology

[General News](#)[Politics](#)[Business](#)[Science&Technology](#)[LifeStyle +](#)[Sport](#)

TRENDING > Renzi • Budget • Grillo euro exit • Pope Francis • [Expo2015](#)

ANSA.it • English • Science & Technology • [1938 disappearance physicist Majorana 'alive 1955-59'](#)

1938 disappearance physicist Majorana 'alive 1955-59'

Rome prosecutors close case of mysterious disappearance

Redazione ANSA

ROME

04 February 2015

16:48

NEWS

Suggerisci

Facebook

Twitter

Google+

Altri

A+ A A-

Stampa

125

Tweet

79

8+1

37

in LinkedIn



Particle superpositions

“equal superposition of an electron and a hole”

$$c^\dagger \left| \rightarrow \right\rangle = \left| \bullet \rightarrow \right\rangle$$

$c^\dagger = \gamma_1 + i\gamma_2$ creates an electron

$$c \left| \bullet \rightarrow \right\rangle = \left| \rightarrow \right\rangle$$

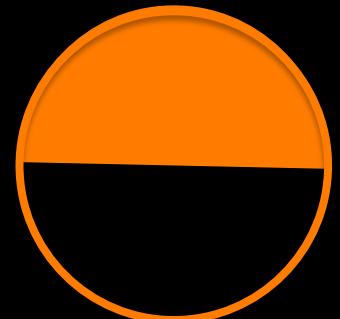
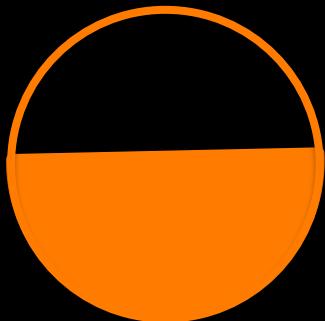
$c = \gamma_1 - i\gamma_2$ creates a hole

$$\gamma_1 = \frac{1}{2} (c^\dagger + c) \text{ is Majorana-1}$$

$$\gamma_2 = \frac{1}{2i} (c^\dagger - c) \text{ is Majorana-2}$$

satisfying $\gamma_1^\dagger = \gamma_1$ and $\gamma_2^\dagger = \gamma_2$

One Majorana is “half” a Fermion



$$c^\dagger = \gamma_1 + i\gamma_2$$

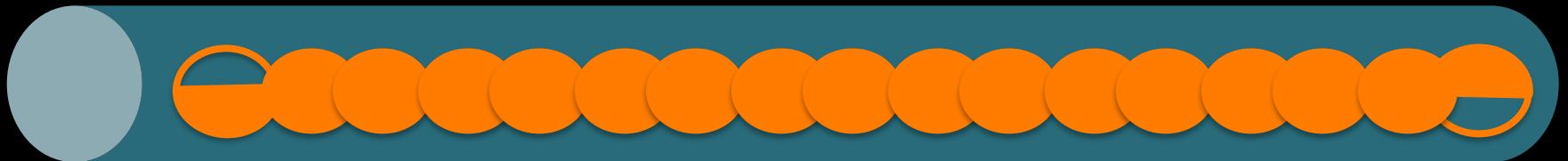


non-local separation provides topological protection

...,Kitaev, Read, Fu, Kane, Das Sarma, Beenakker, Alicea,...
(see review Nick Read, Physics Today July 2012.)



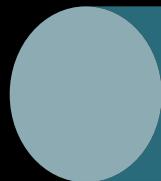
from spatial to particle superposition



from spatial to particle superposition

$$\gamma_1^t = \gamma_1$$

$$\gamma_2^t = \gamma_2$$



Superconductor in magnetic field

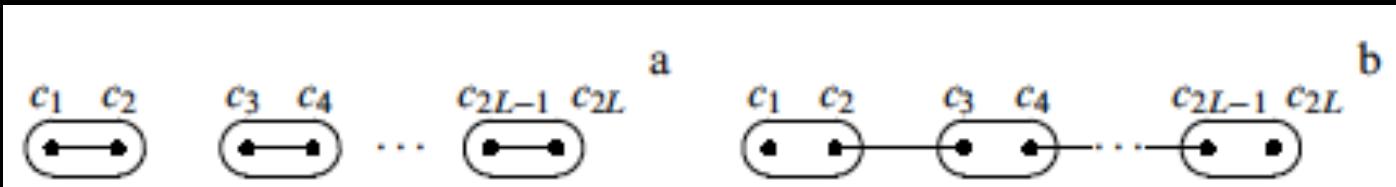


Figure 2. Two types of pairing.

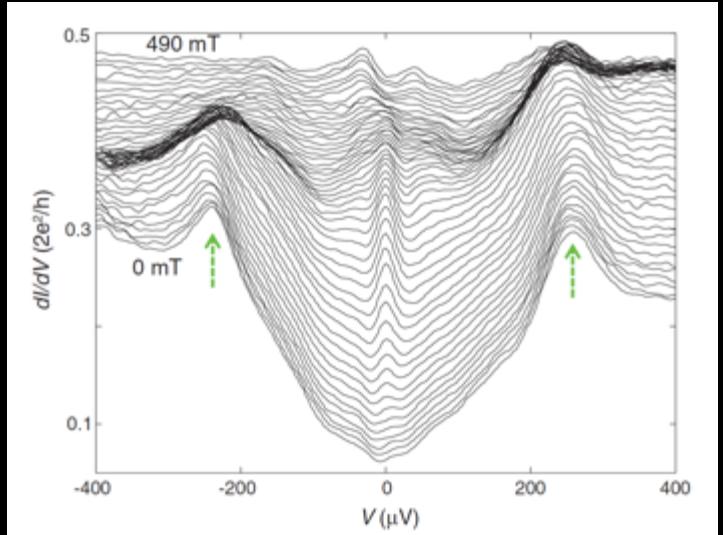
Kitaev 2001

INTRO-2

Majorana recipe


$$\gamma_1^\dagger = \gamma_1$$
$$\gamma_2^\dagger = \gamma_2$$

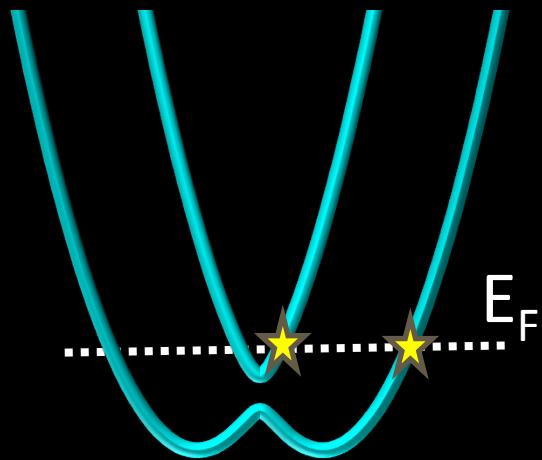
Lutchyn, Sau, Das Sarma, PRL 2010
Oreg, Refael, von Oppen, PRL 2010



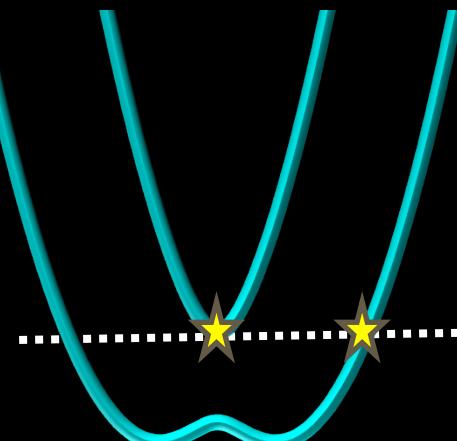
- one-dimensional quantum wire zero bias anomalies
 - spin-orbit interaction →
 - superconductivity
 - tune μ & apply magnetic field
- Kouwenhoven, *Delft* (2012)
Xu, *Lund*
Heiblum, *Weizmann*
Harlingen, *Urbana-Champaign*
Marcus, *Copenhagen*

SOI + Magnetism

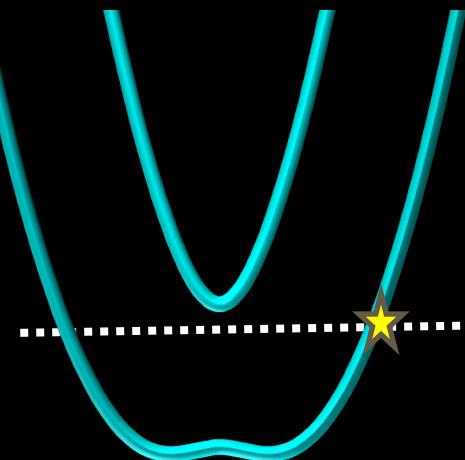
$B < B_{\text{critical}}$



$B = B_{\text{critical}}$



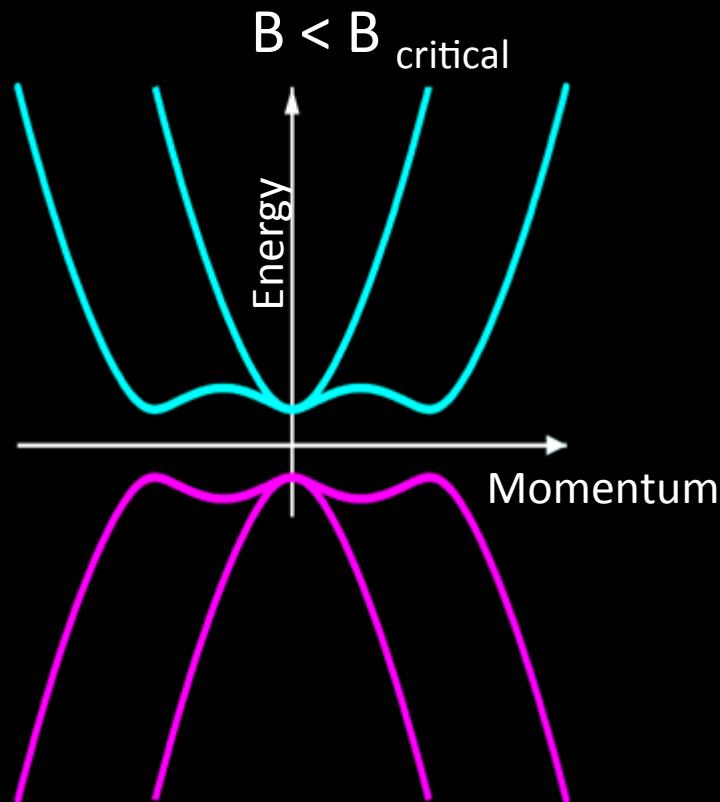
$B > B_{\text{critical}}$



TOPOLOGICALLY DISTINCT
Number of crossings different in parity



SOI + Magnetic field + superconductivity

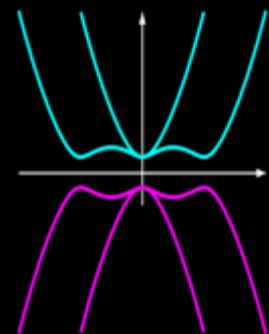


Y. Oreg et al. PRL(2010); R. M. Lutchyn et al PRL (2010);

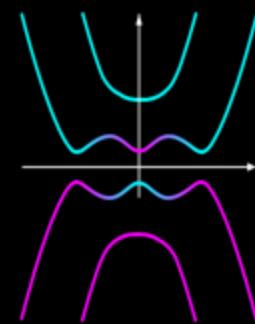


Topology & bandstructure in 1D

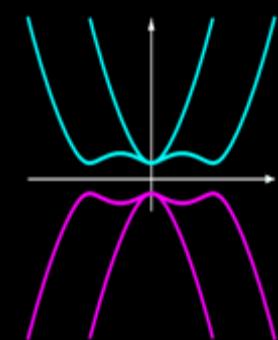
Non-Inverted
band structure



Inverted
band structure



Non-Inverted
band structure



TRIVIAL

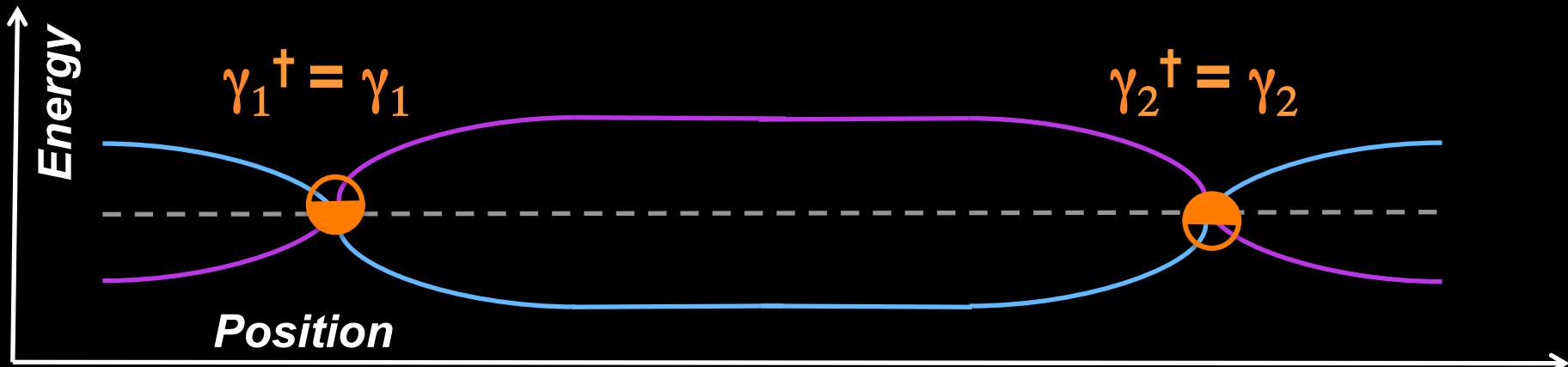
TOPOLOGICAL

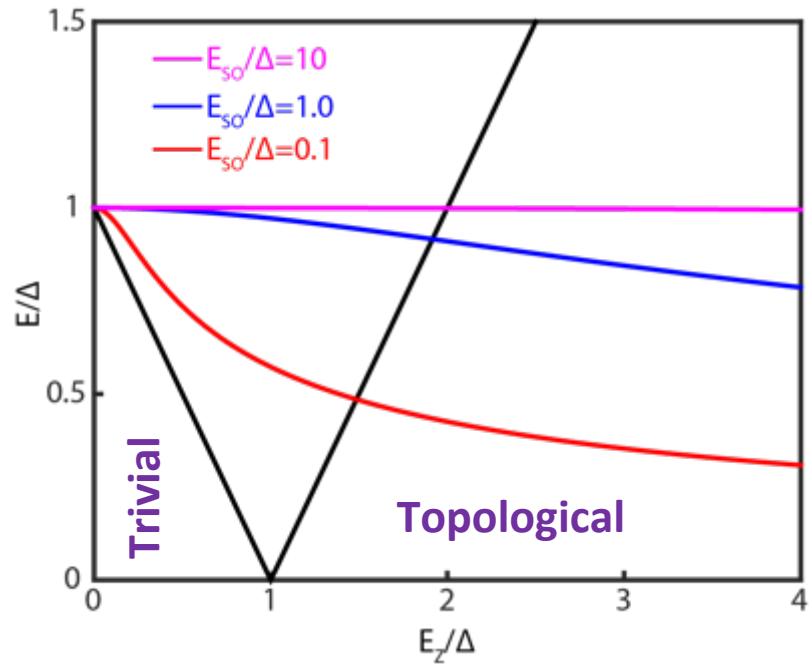
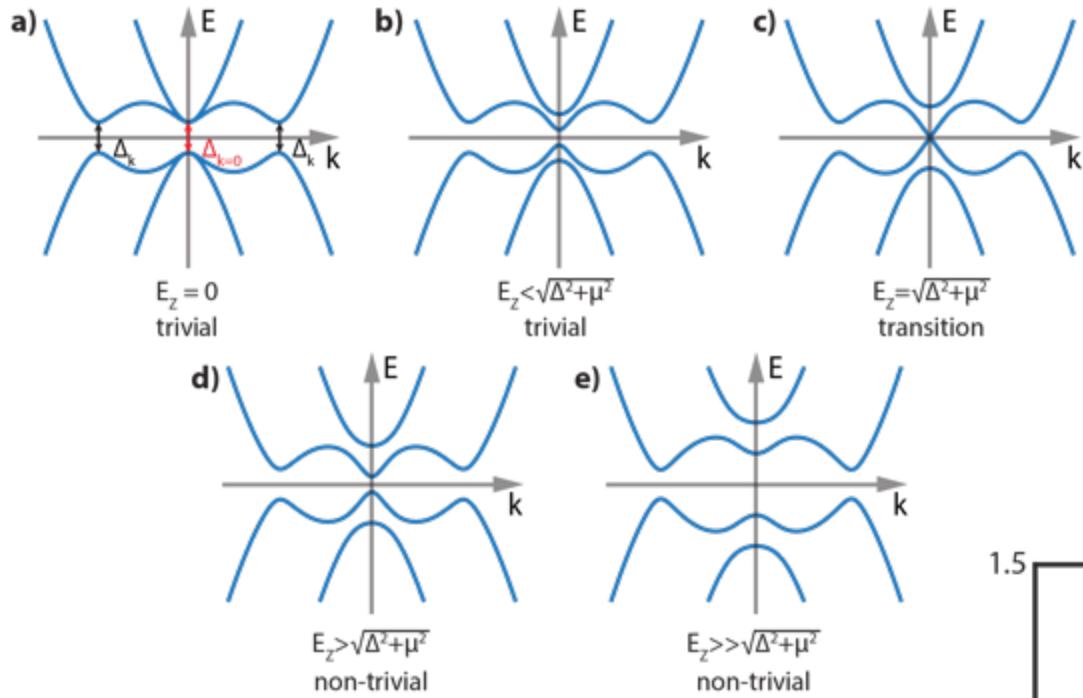
TRIVIAL

$$\gamma_1^\dagger = \gamma_1$$

$$\gamma_2^\dagger = \gamma_2$$

Position





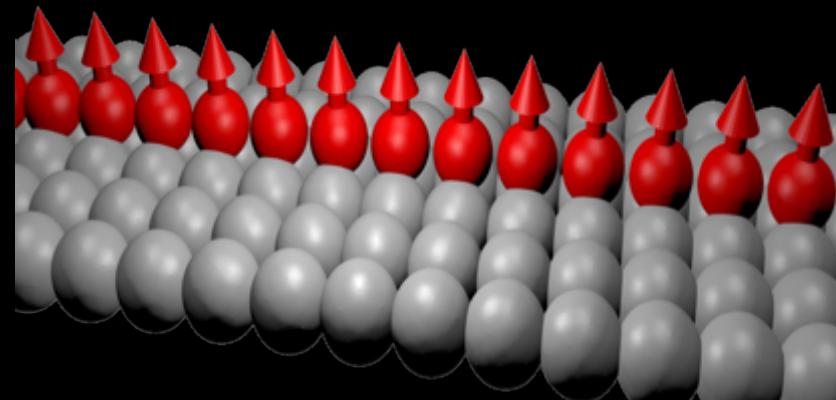
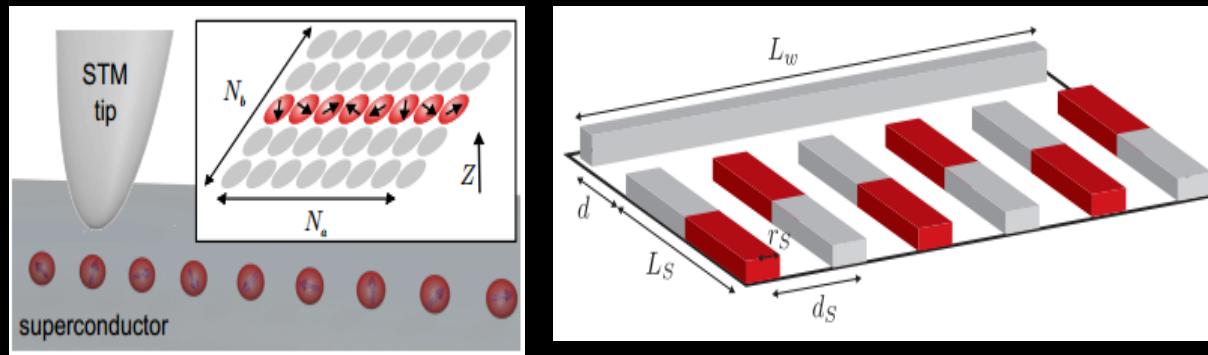
Thesis 2016, TU Delft repository
Kun Zuo & Vincent Mourik

effective SOI from magnetic texture

Majorana fermions emerging from magnetic nanoparticles on a superconductor without spin-orbit coupling

T.-P. Choy, J. M. Edge, A. R. Akhmerov, and C. W. J. Beenakker

Instituut-Lorentz, Universiteit Leiden, P.O. Box 9506, 2300 RA Leiden, The Netherlands



Kjaergaard, Wölms, Flensberg (PRB 2012)

Klinovaja, Stano, Loss (PRL 2012)

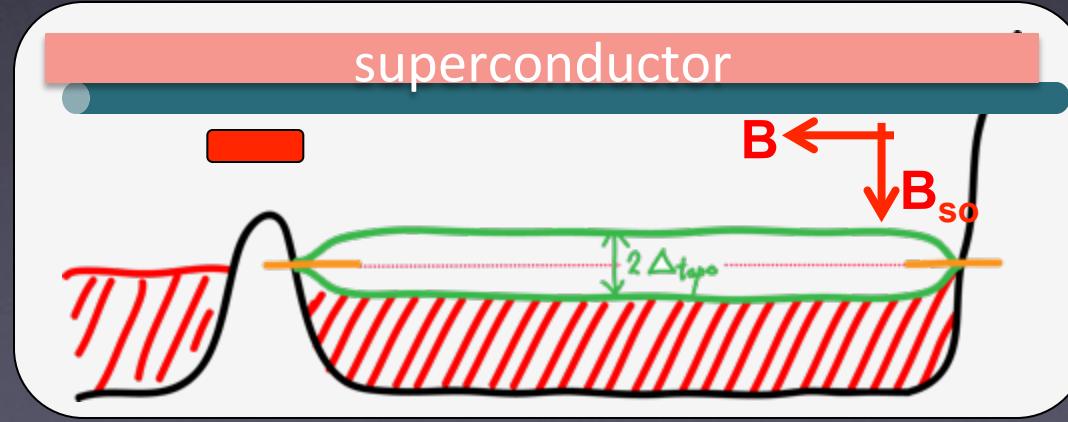
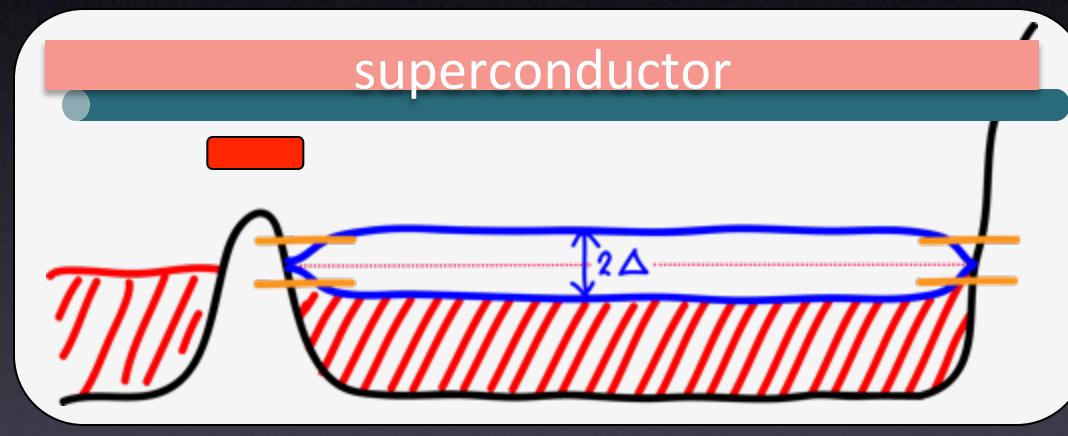
Nakosai, Tanaka, Nagaosa (PRB 2013)

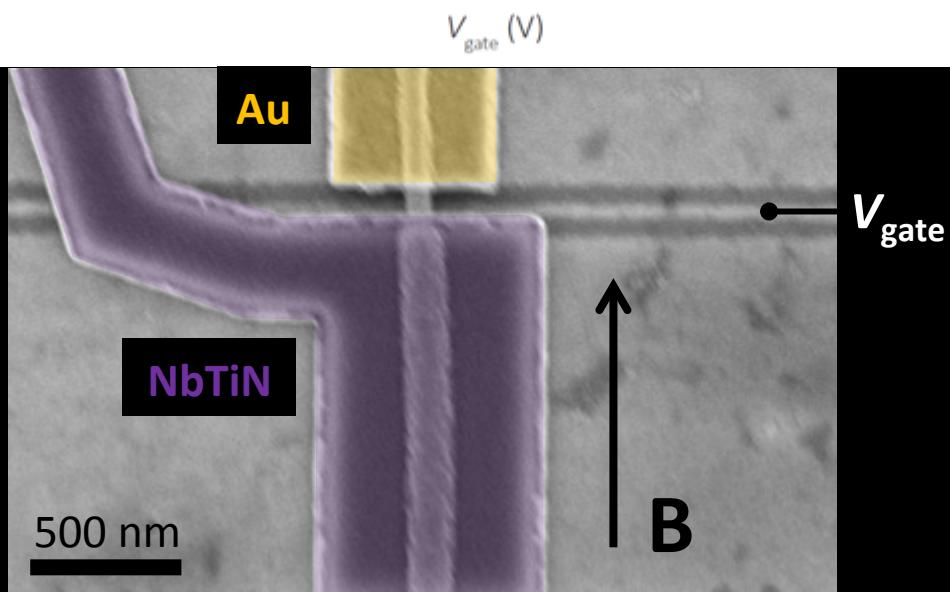
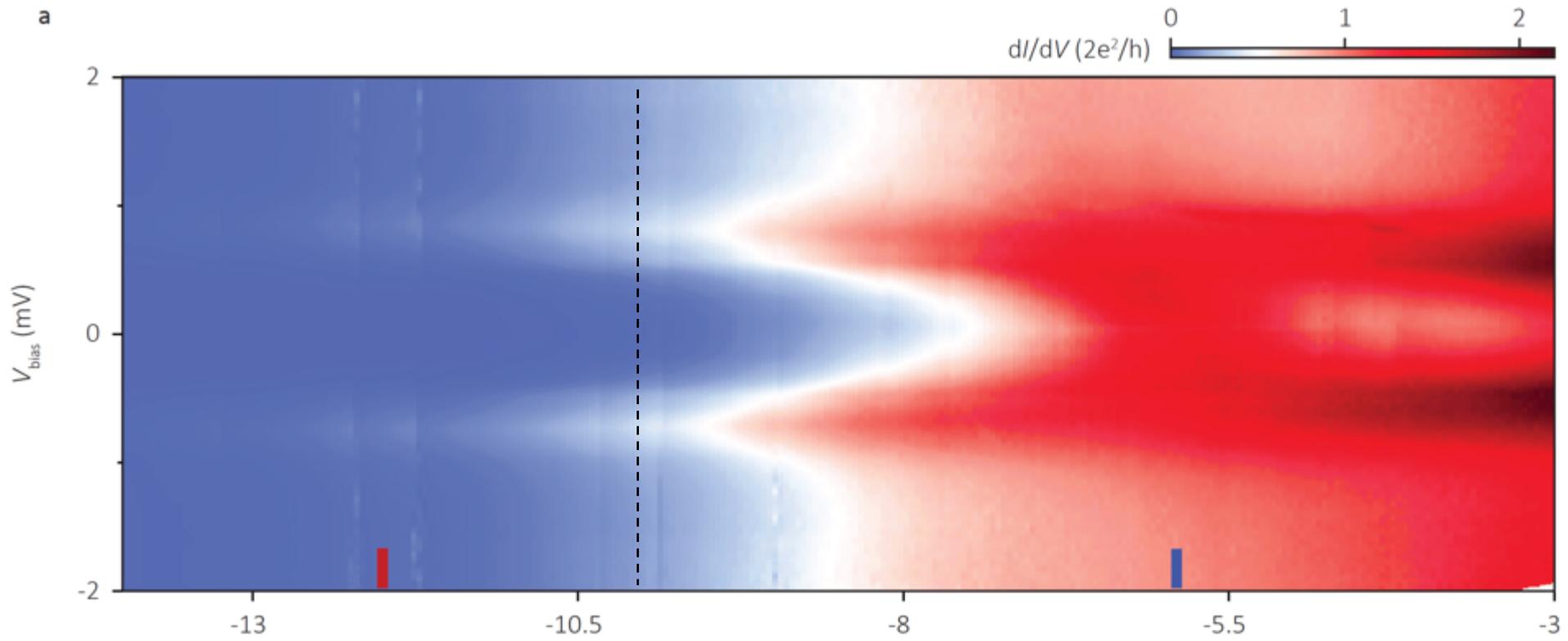
Martin and Morpurgo (PRB 2012)

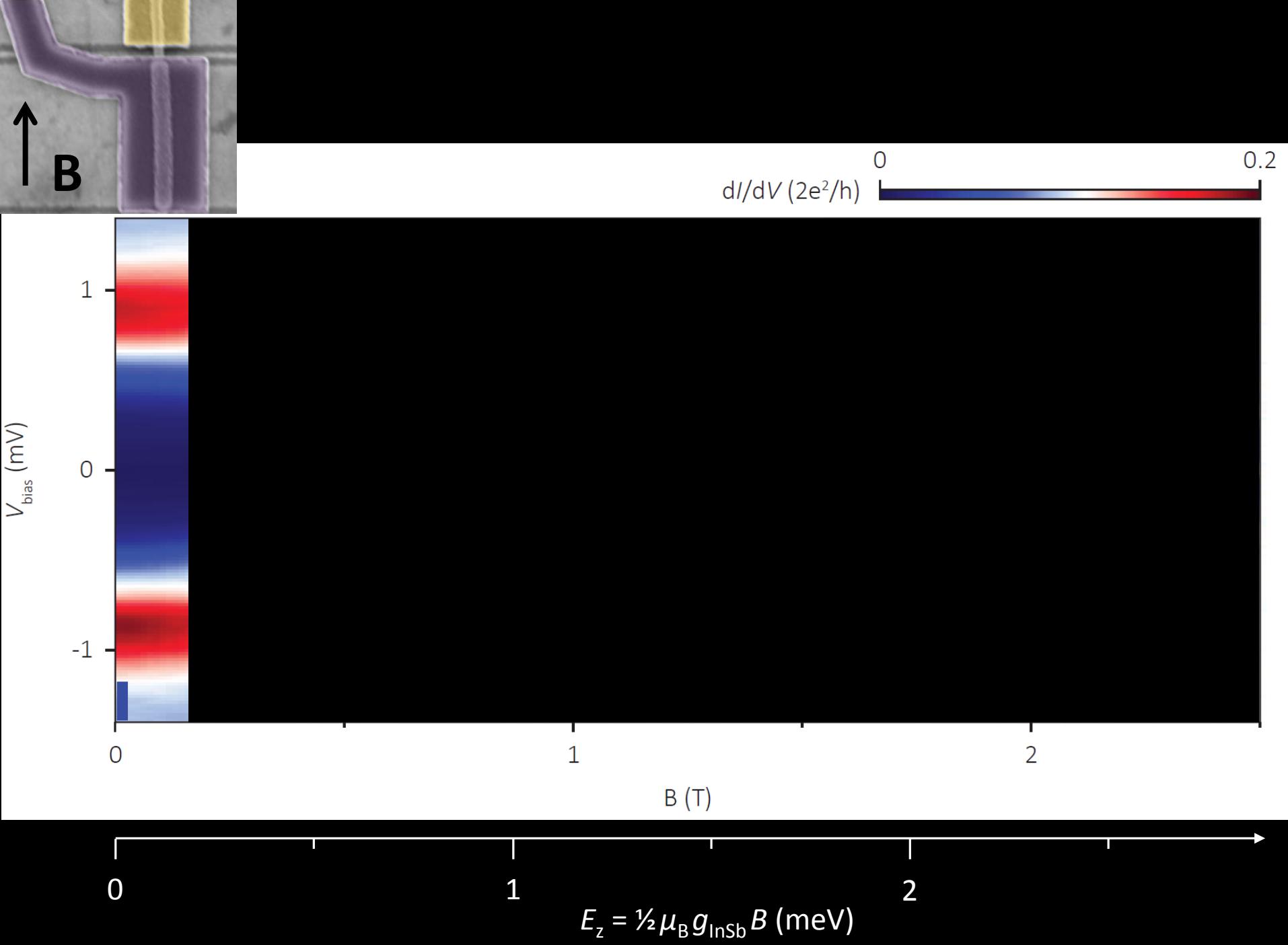
Glazman (PRB 2013), and others ..

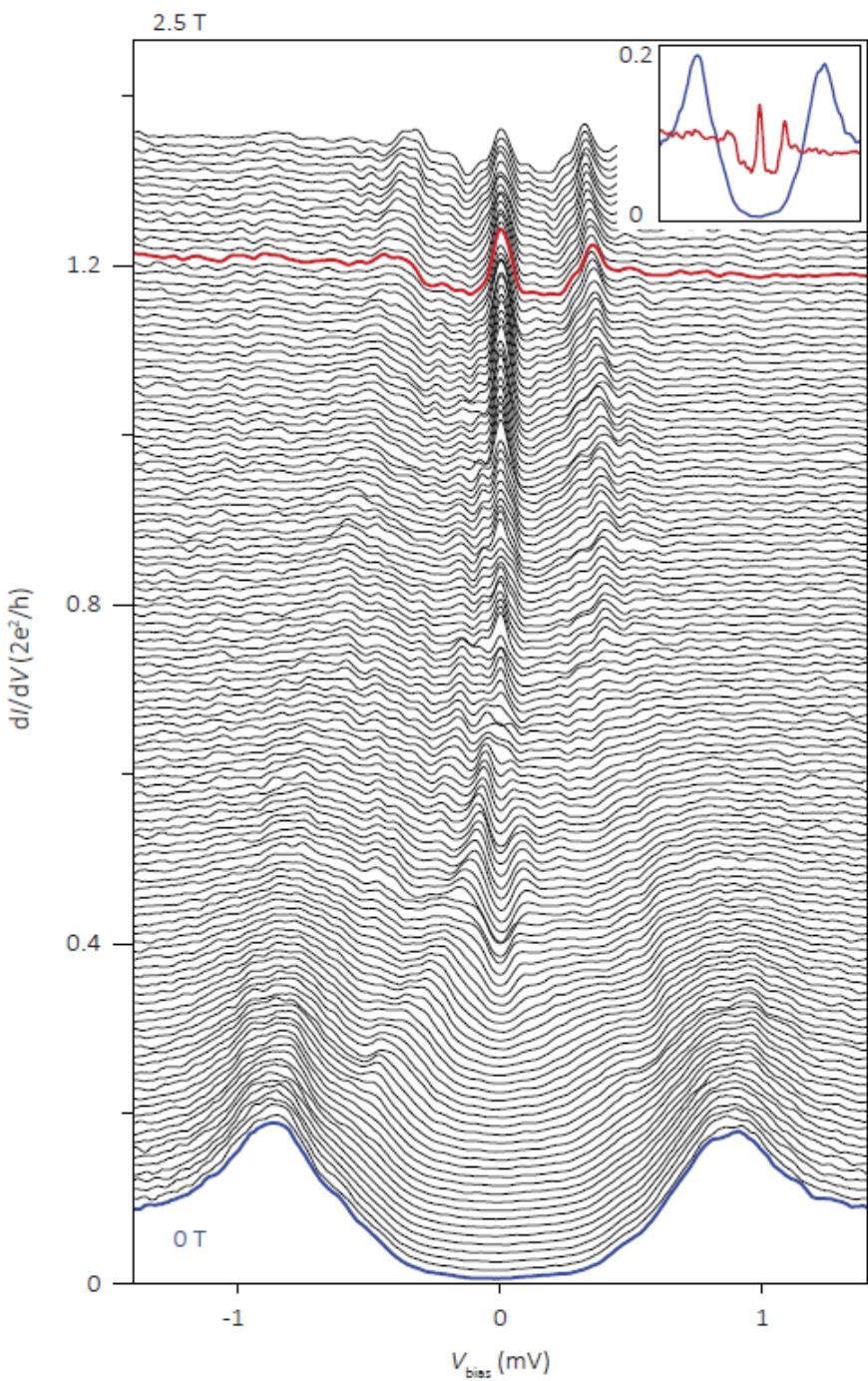
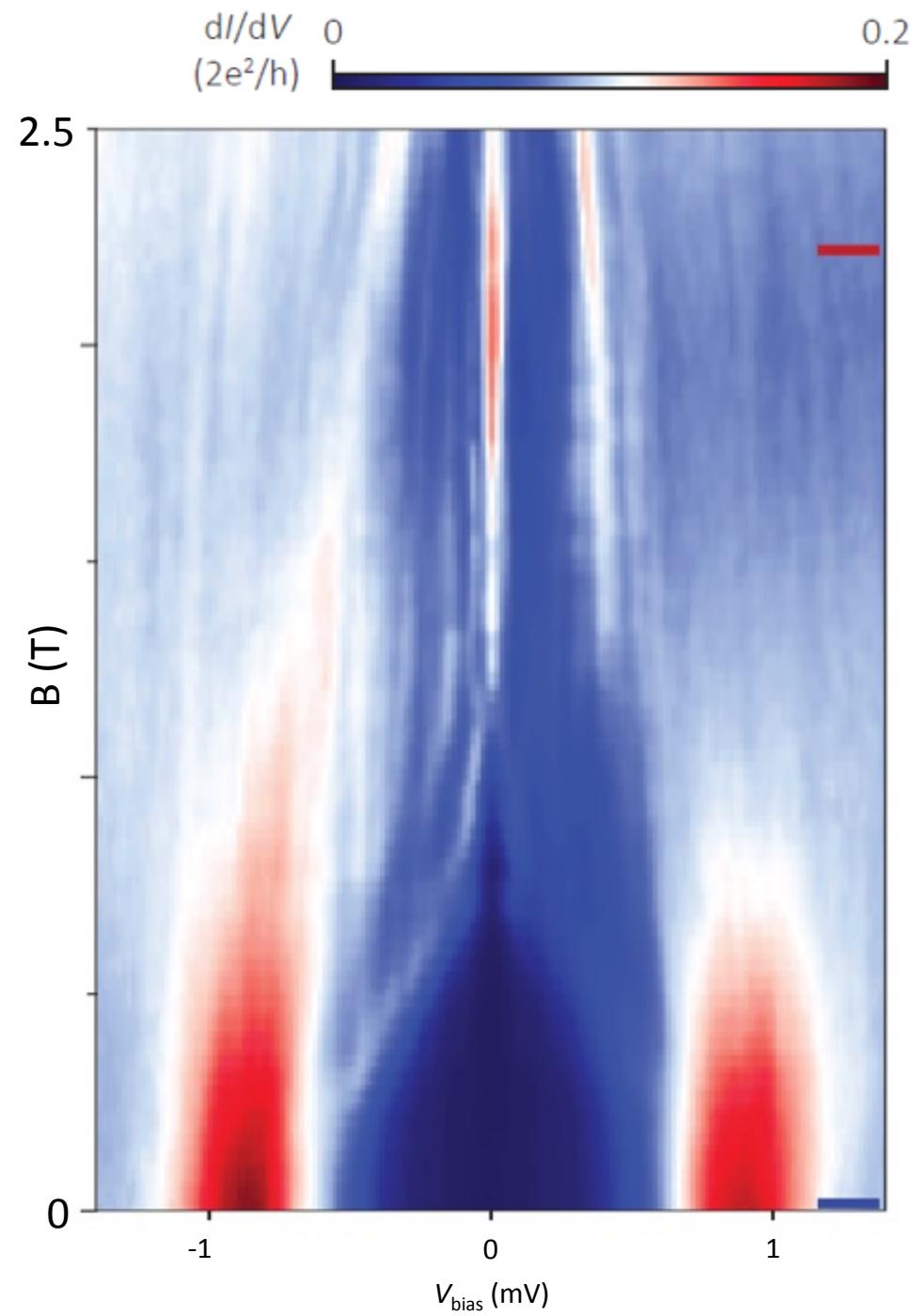
Nadj-Perge, Yazdani,
et al. (2014)

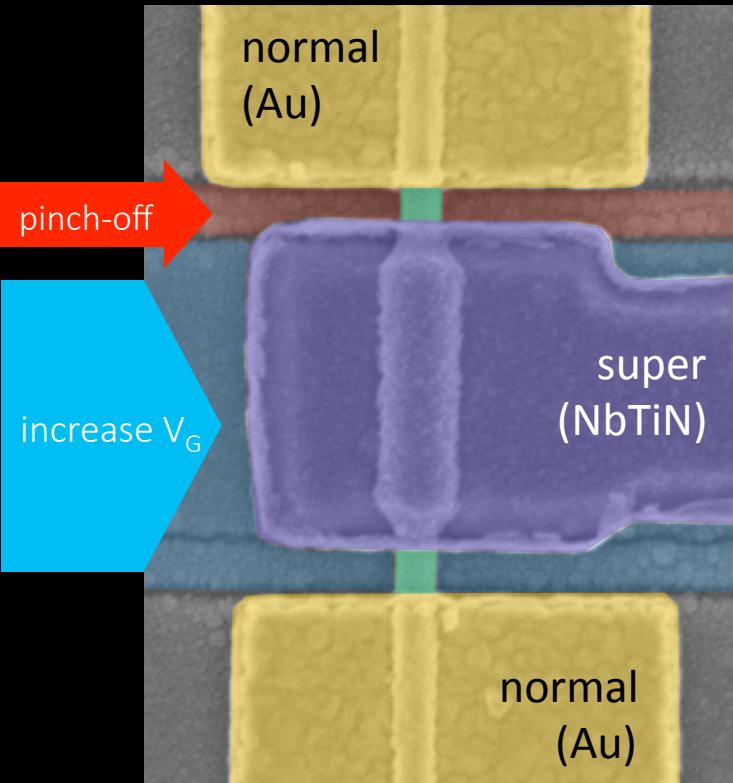
ZBP



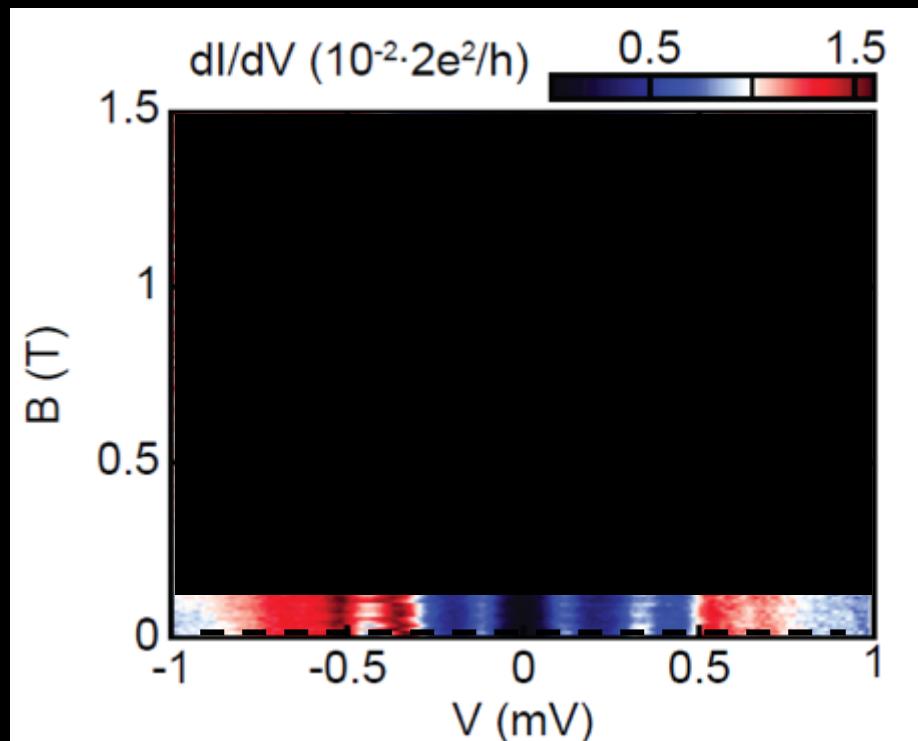


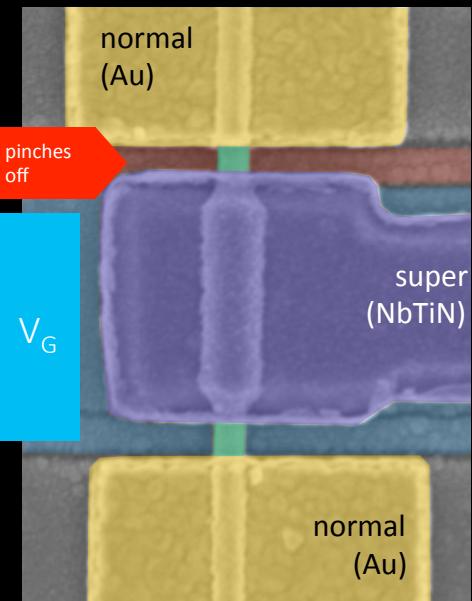






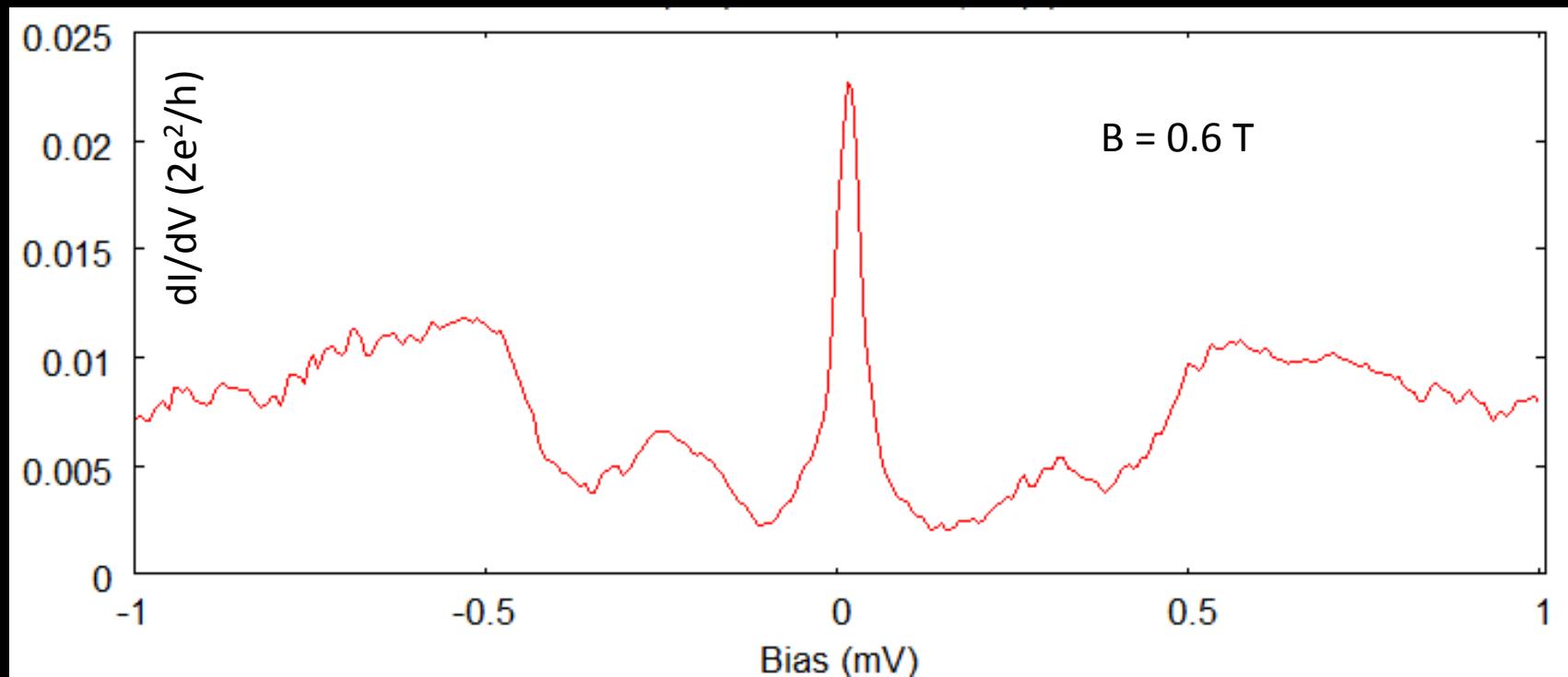
↑ $B \parallel$ wire



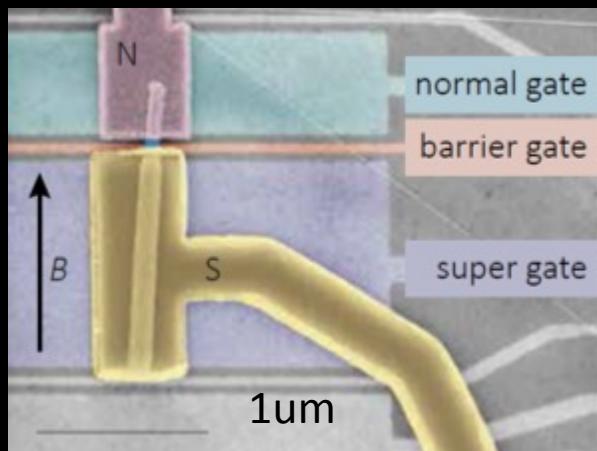


↑ $B \parallel$ wire

ZBP exceeds normal state conductance

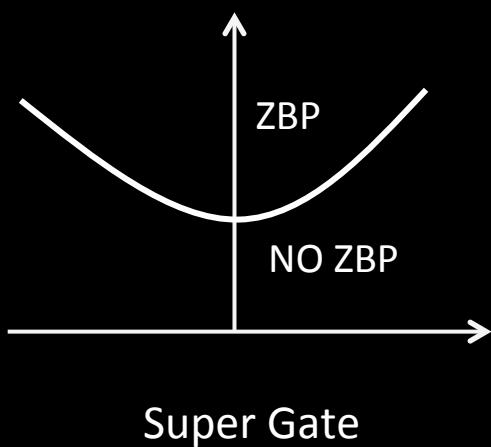


A typical Majorana Device:

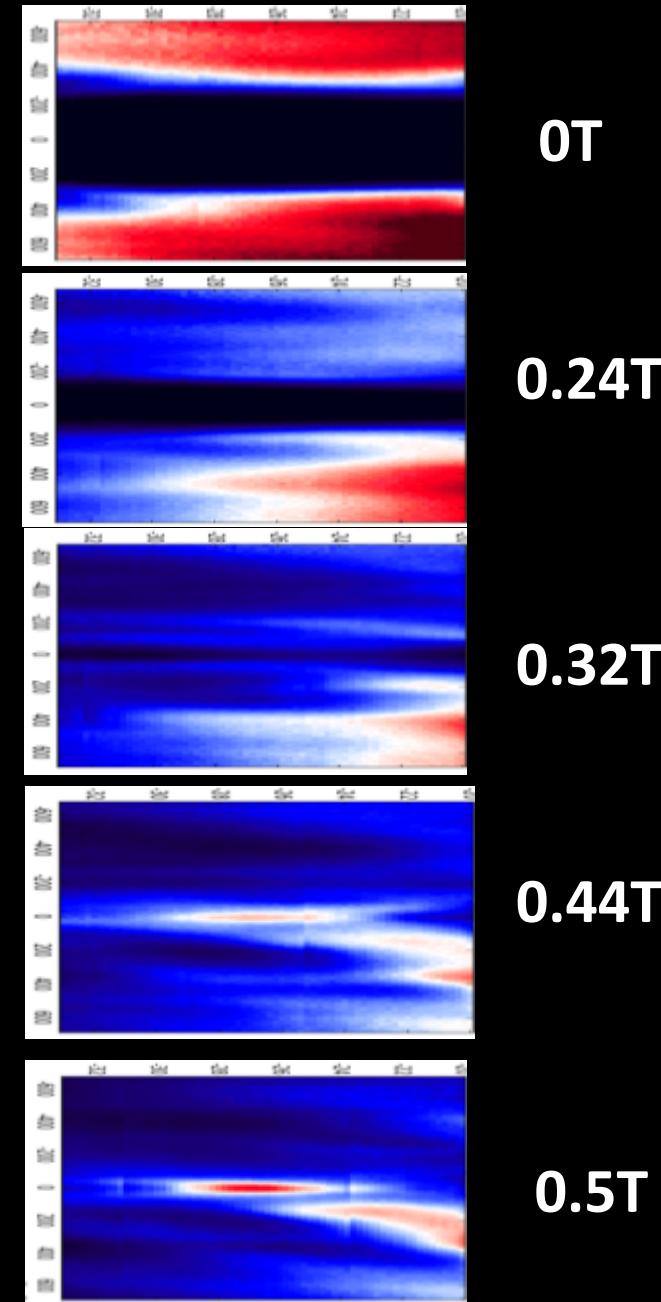


**Majorana condition
(single sub-band regime)**

↓
B
↓
Super gate

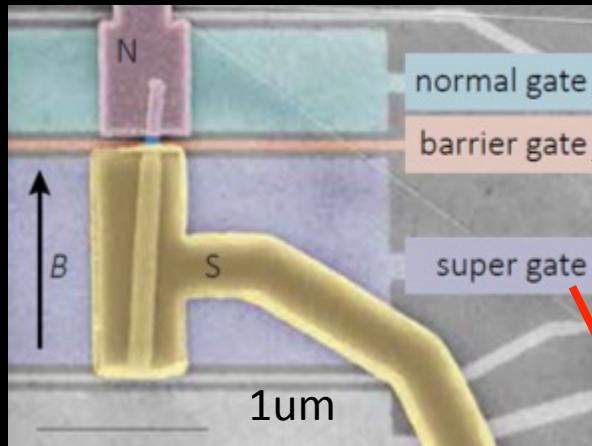


B (T)
↓

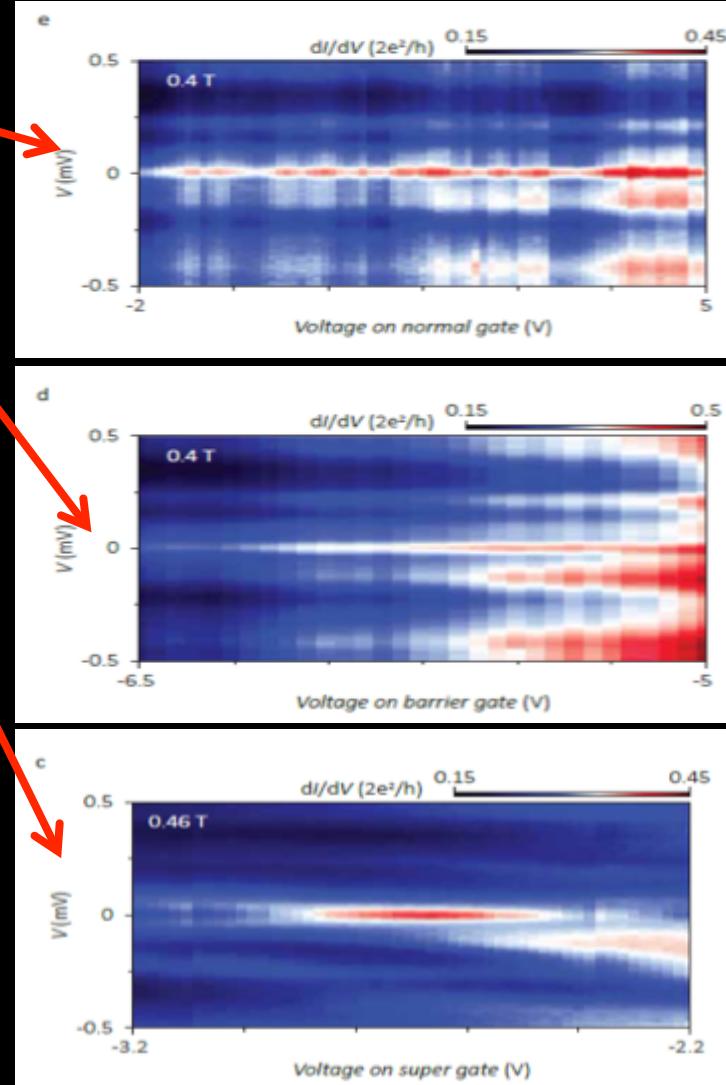
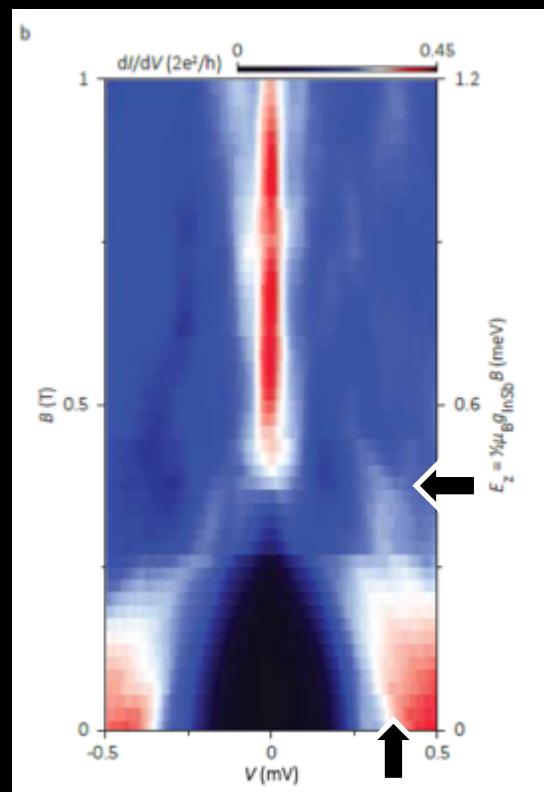


Super gate

A typical Majorana Device:

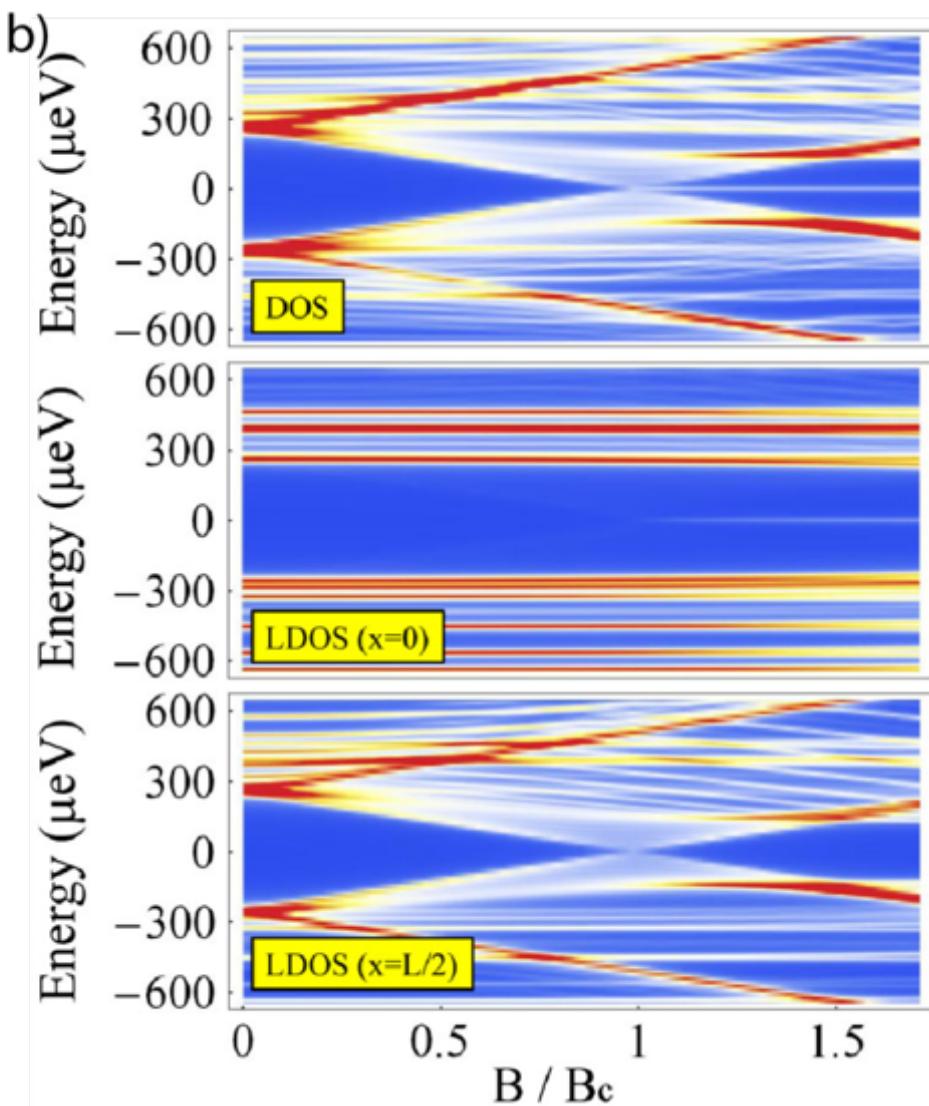
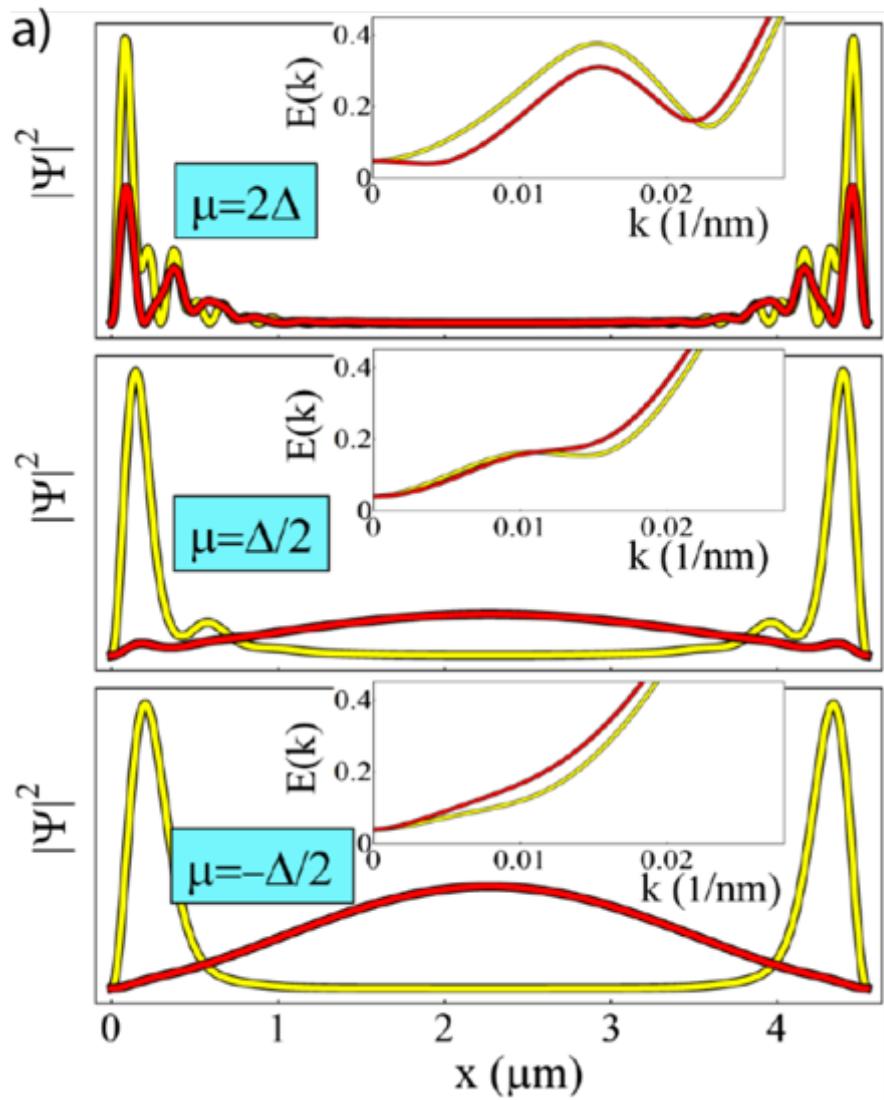


Where does the ZBP live?

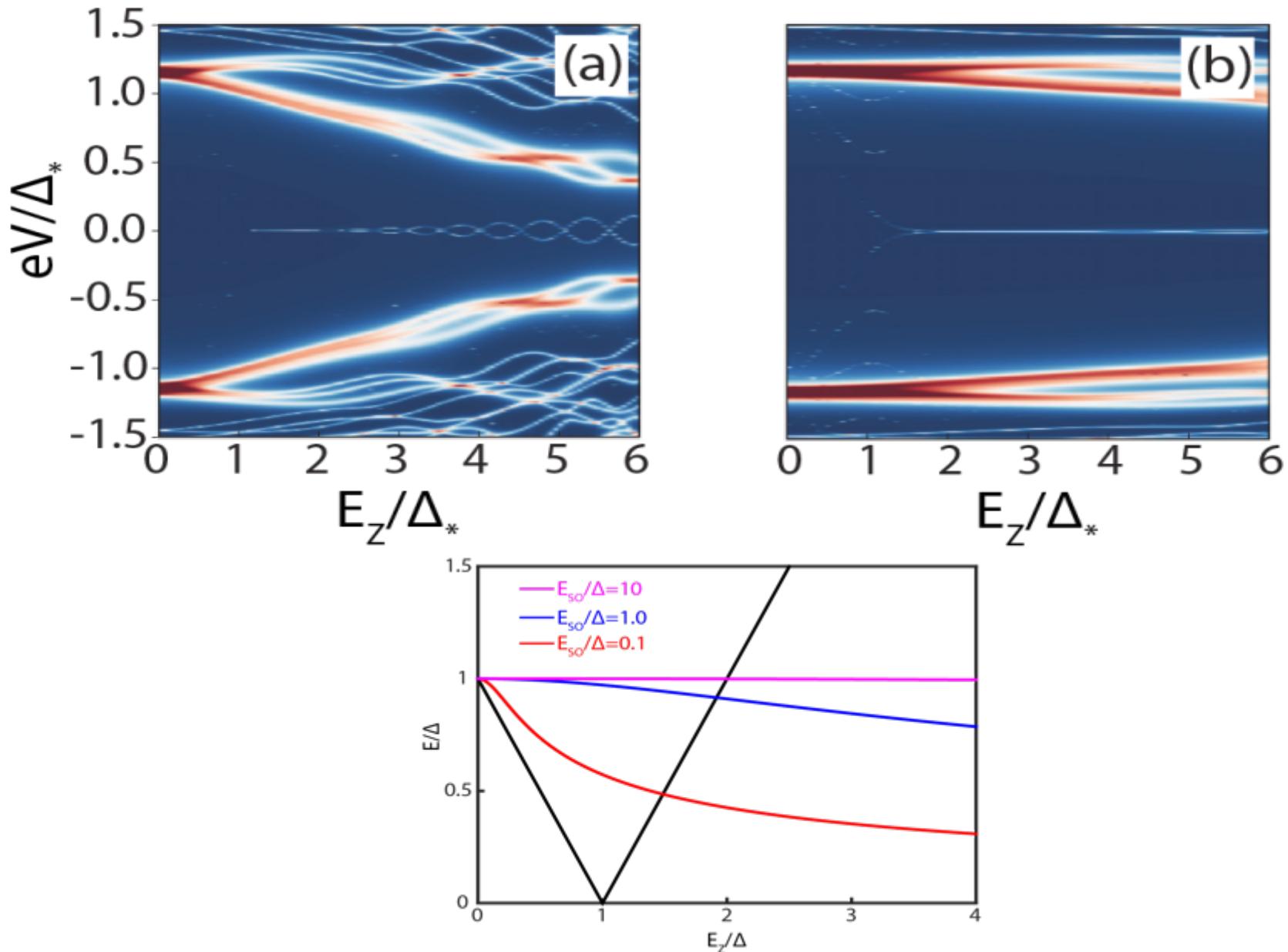


The ZBP lives underneath the superconductor

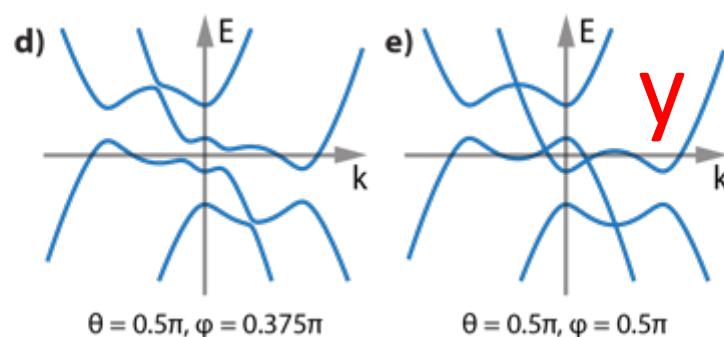
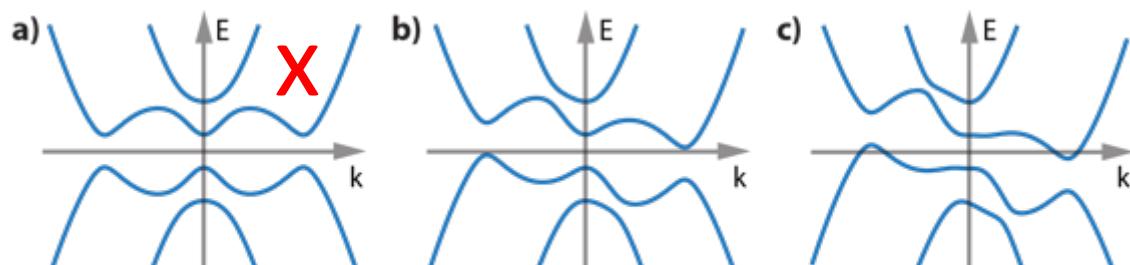
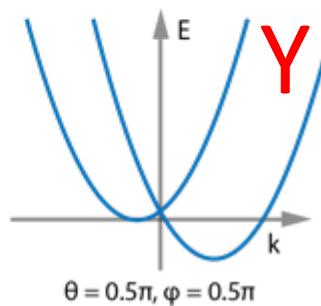
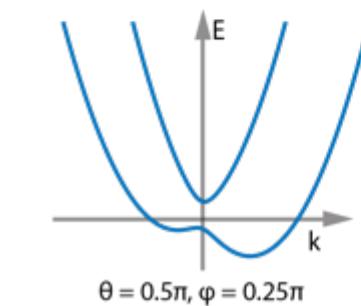
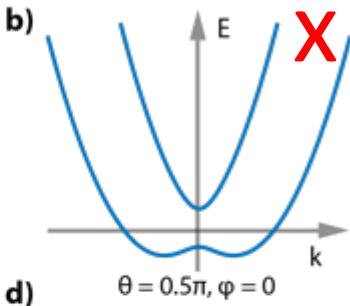
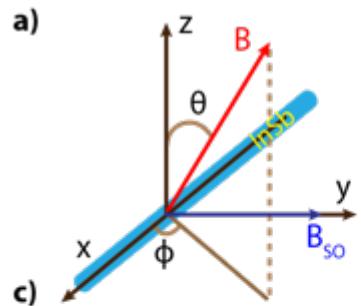
Numerics Stanescu, Tewari, Sau & Das Sarma

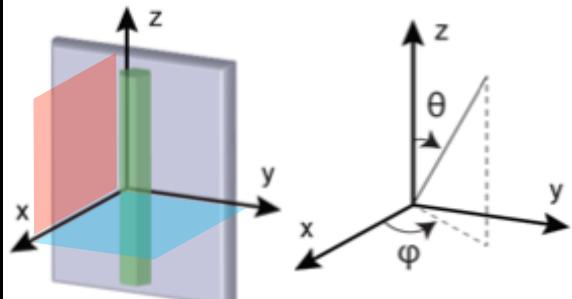


Numerics Rainis, Trifunovic, Klinovaja & Loss

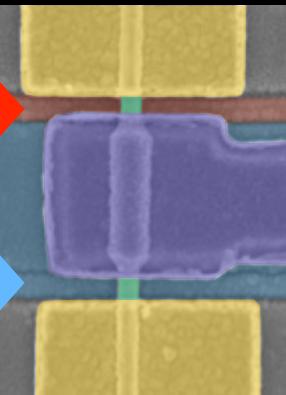


Angle





spin-orbit along y

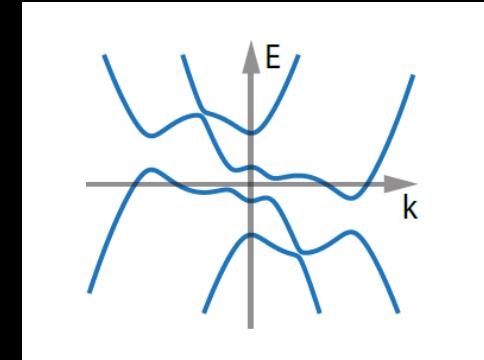
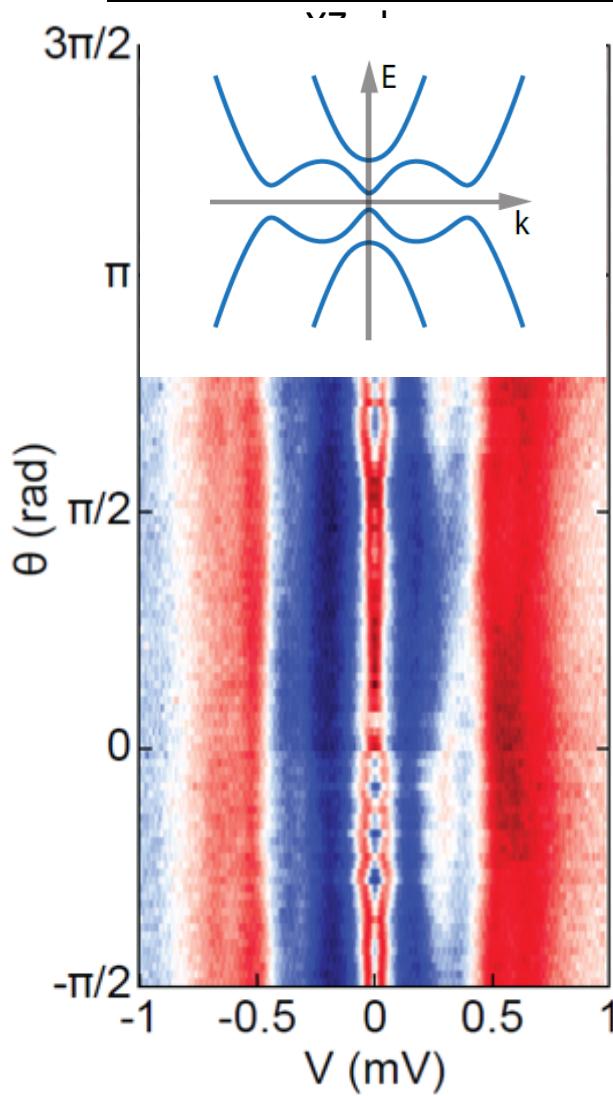


$B = 0.42 \text{ T}$

rotation of magnetic field

XZ plane \perp spin-orbit

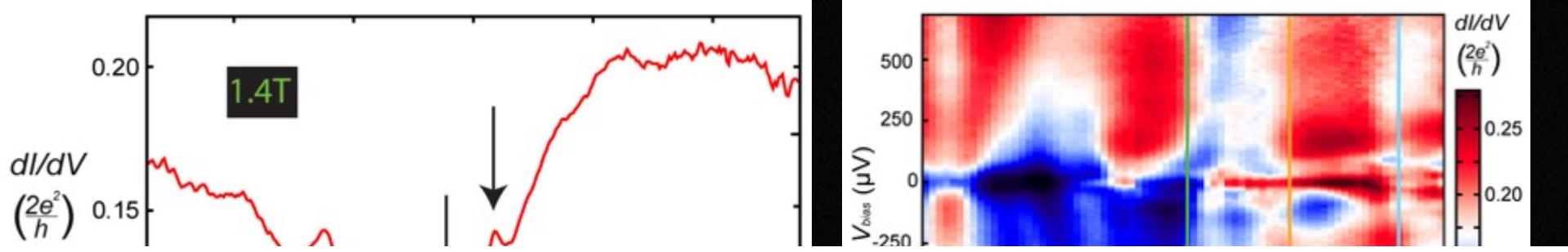
XY plane // spin-orbit



ZBP

SPLITTING

2 Nov 2012



A Majorana smoking gun for the superconductor-semiconductor hybrid topological system

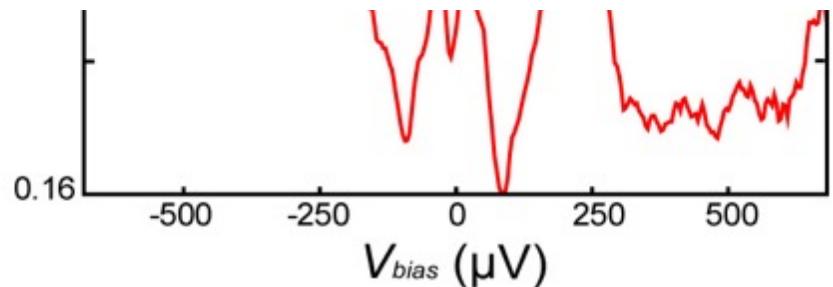
S. Das Sarma¹, Jay D. Sau², and Tudor D. Stanescu³

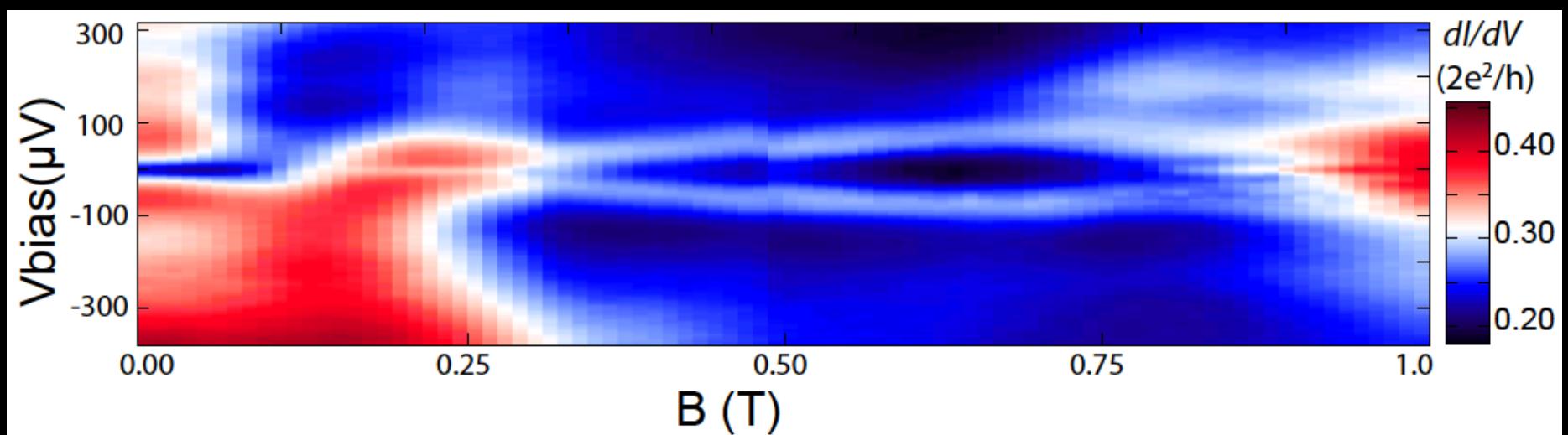
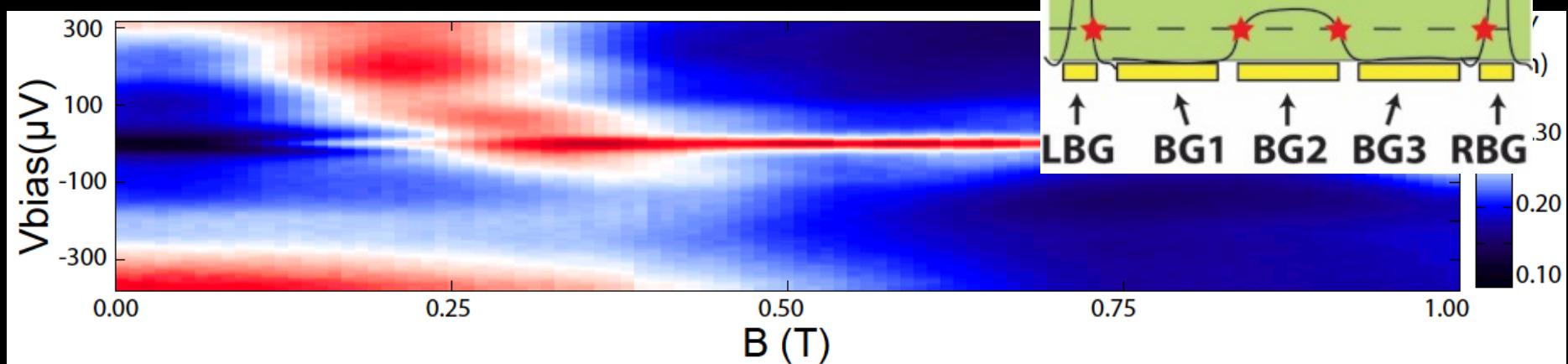
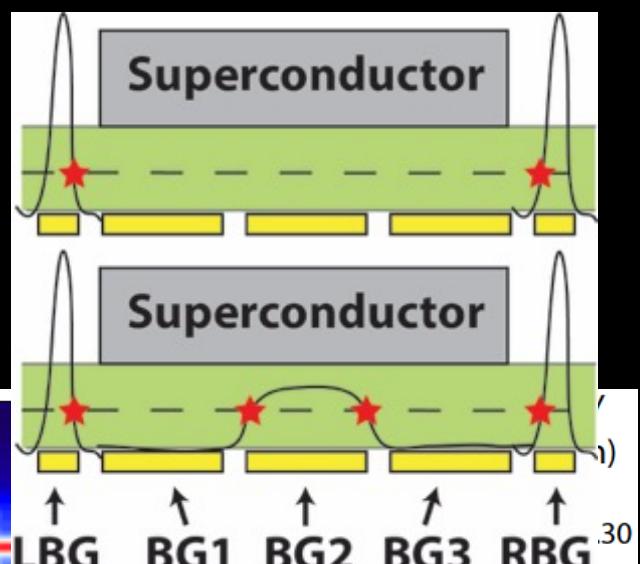
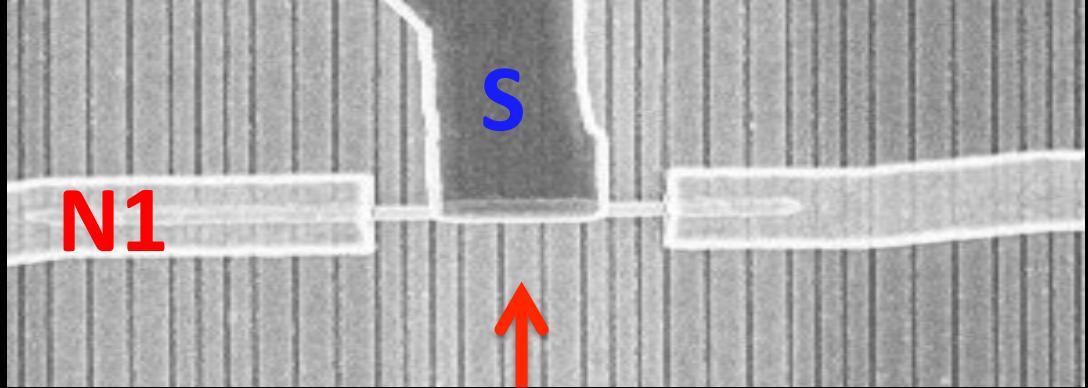
¹*Condensed Matter Theory Center, Department of Physics,
University of Maryland, College Park, Maryland 20742-4111, USA*

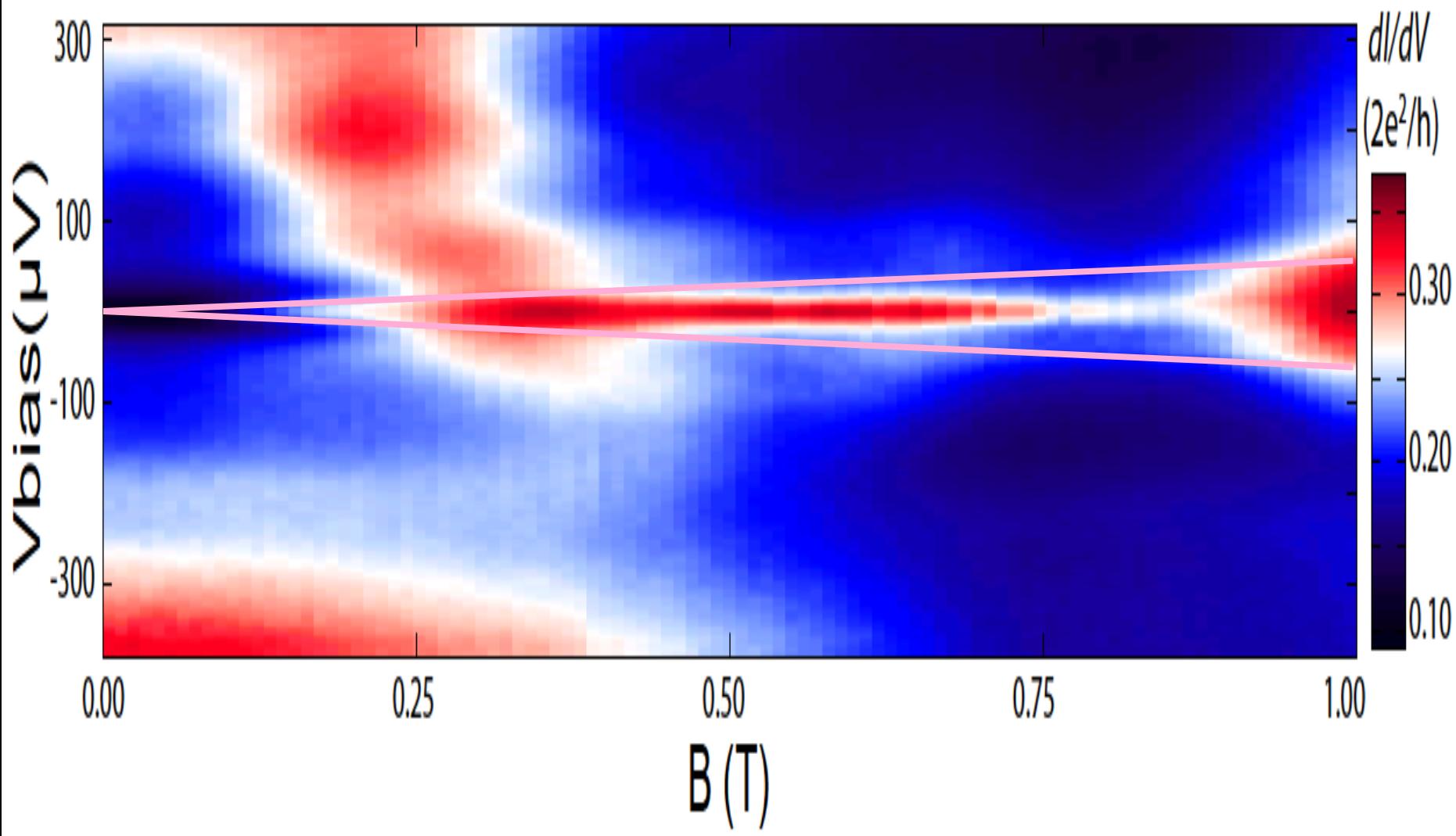
²*Department of Physics, Harvard University, Cambridge, Massachusetts 02138, USA*

³*Department of Physics, West Virginia University, Morgantown, WV 26506, USA*

Recent observations of a zero bias conductance peak in tunneling transport measurements in superconductor–semiconductor nanowire devices provide evidence for the predicted zero-energy Majorana modes, but not the conclusive proof for their existence. We establish that direct observation of a splitting of the zero bias conductance peak can serve as the smoking gun evidence for the existence of the Majorana mode. We show that the splitting has an oscillatory dependence on the Zeeman field (chemical potential) at fixed chemical potential (Zeeman field). By contrast, when the density is constant rather than the chemical potential – the likely situation in the current experimental set-ups – the splitting oscillations are generically suppressed. Our theory predicts the conditions under which the splitting oscillations can serve as the smoking gun for the experimental confirmation of the elusive Majorana mode.

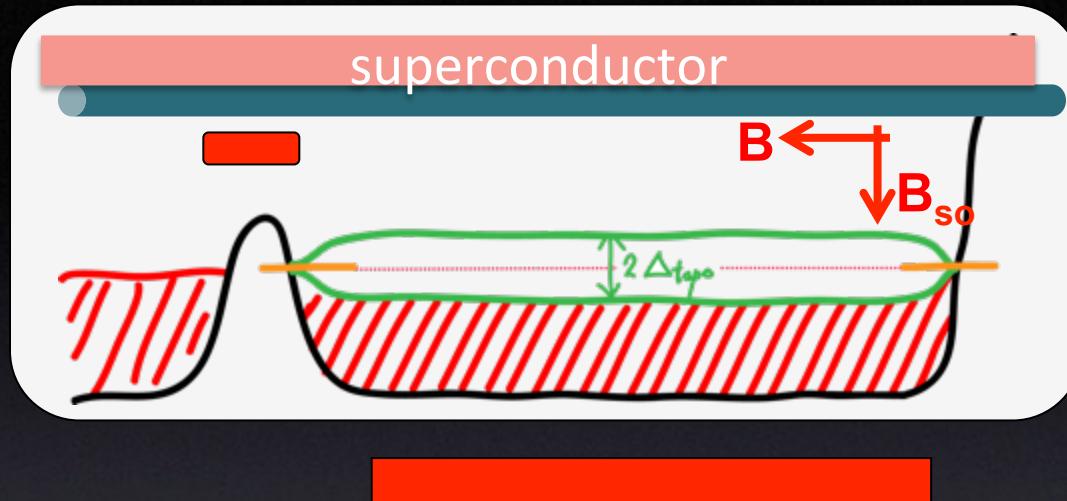




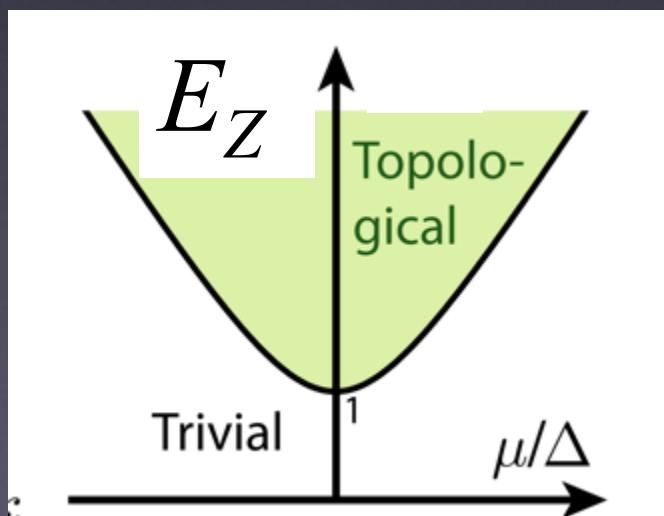


if Kondo-1/2, then $g\mu_B B < 60 \mu\text{eV}$ and thus $g < 1 \ll g_{\text{InSb}}/50$

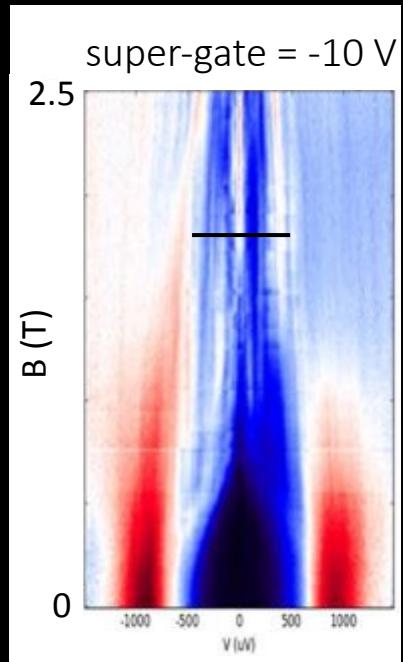
PHASE DIAGRAM

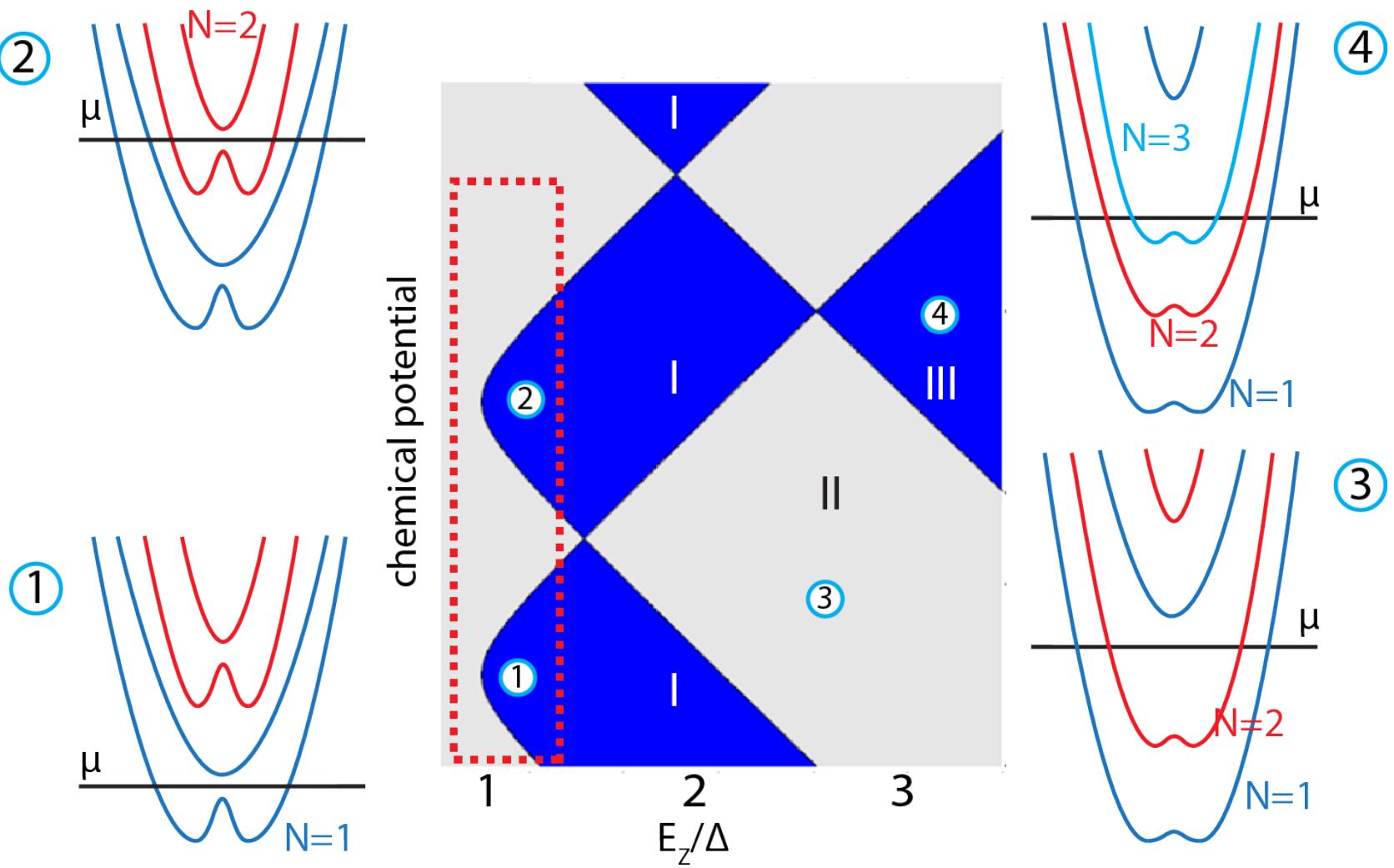


super-Vg

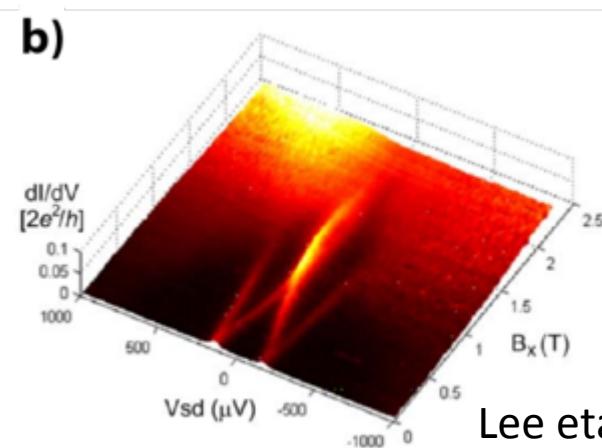
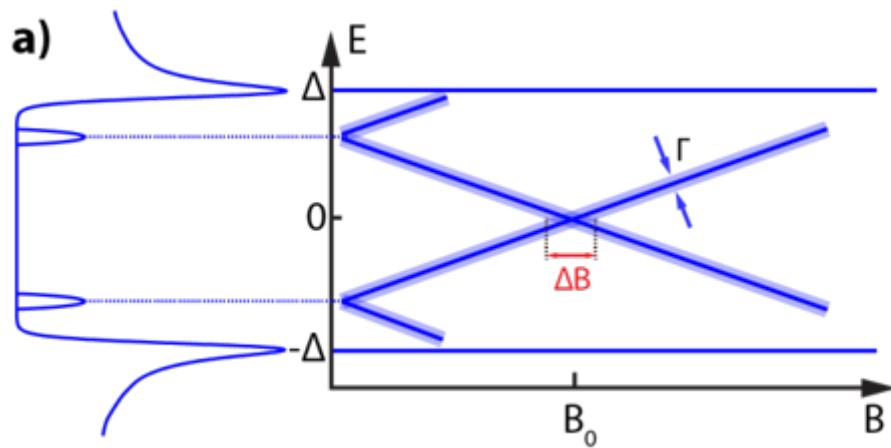


$$E_Z > \sqrt{\Delta^2 + \mu^2}$$

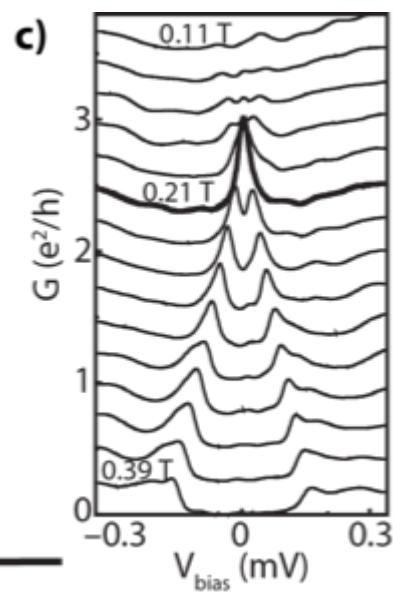
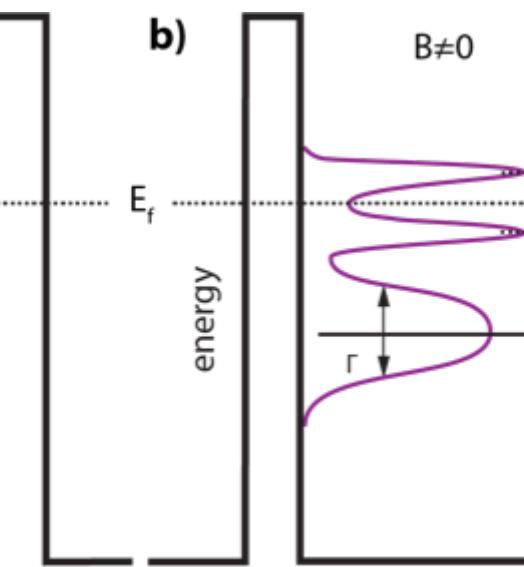
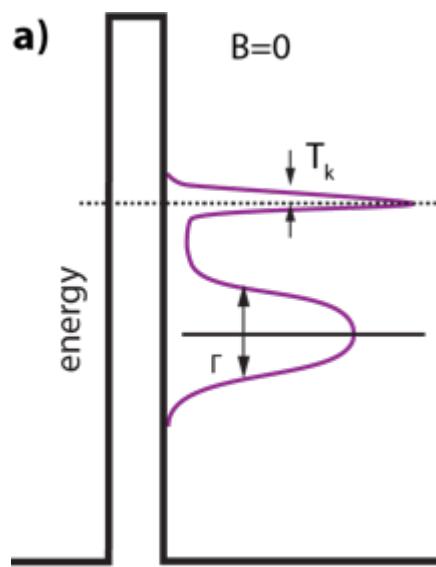




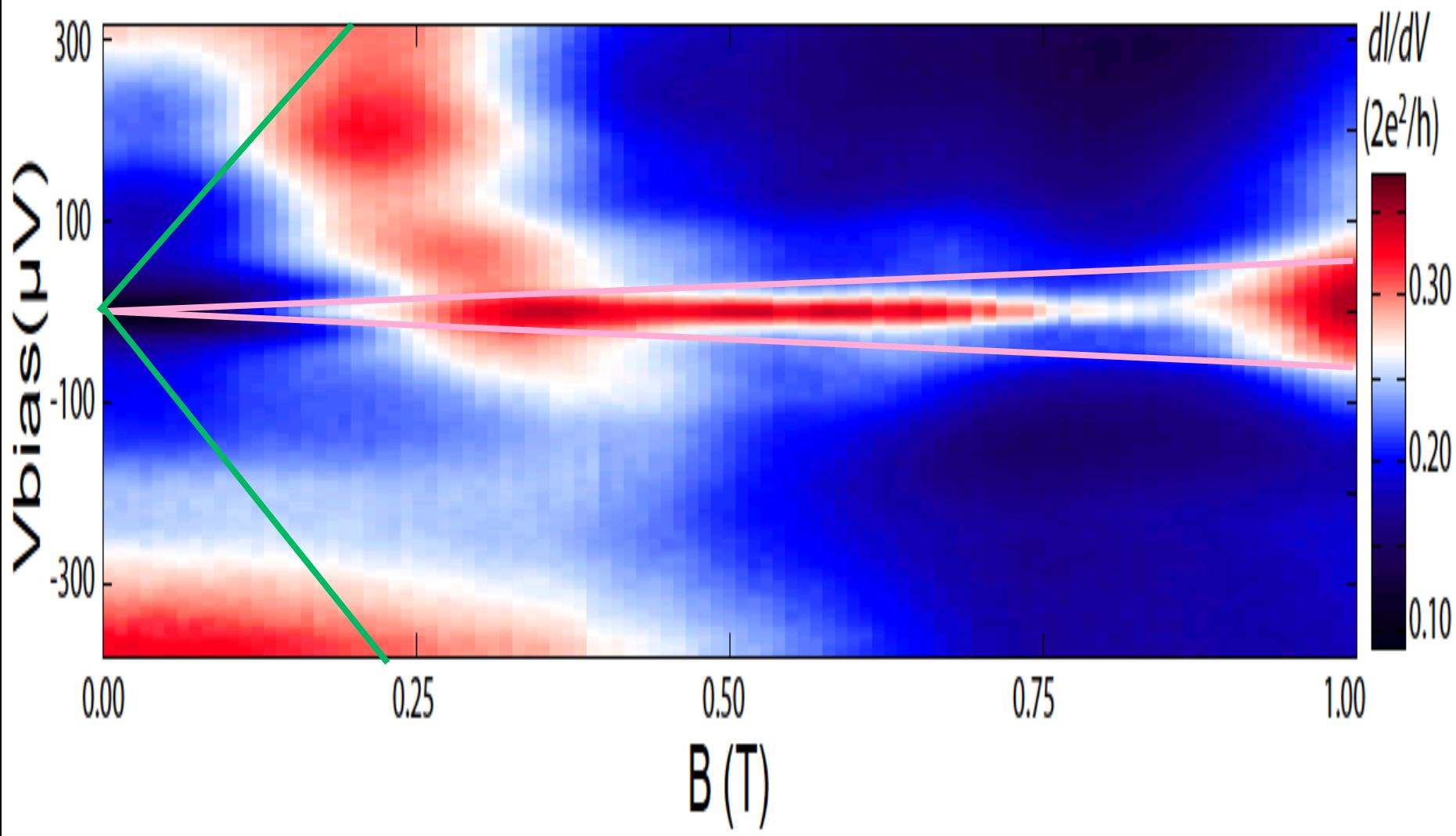
Level-crossing
&
KONDO



Lee et al.

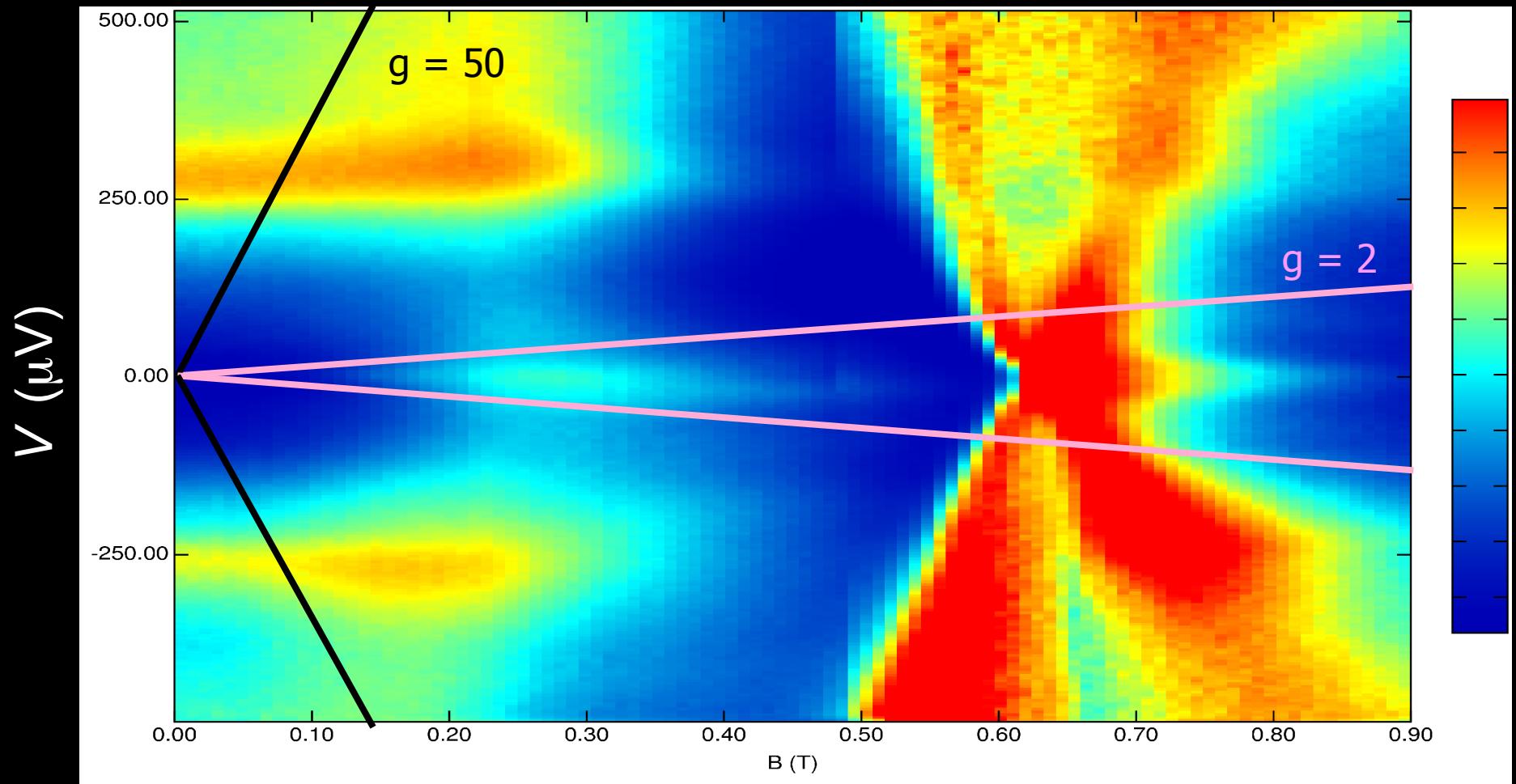


ST-Kondo, Sasaki et al.



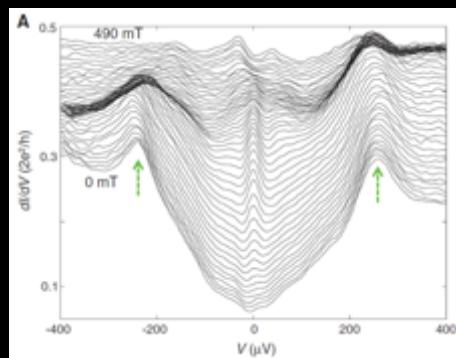
if Kondo-1/2, then $g\mu_B B < 60 \text{ }\mu\text{eV}$ and thus $g < 1 \ll g_{\text{InSb}}/50$

how would spin $\frac{1}{2}$ Kondo look?

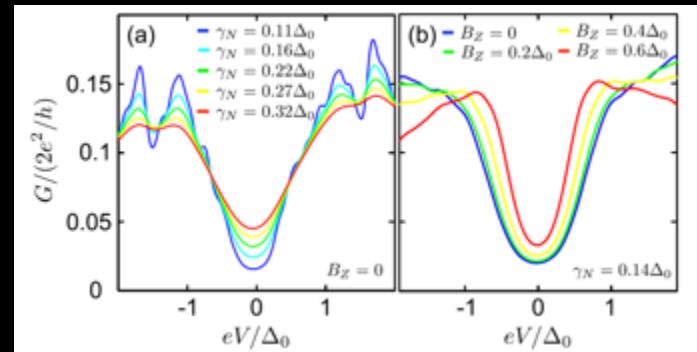


DISORDER

challenge: disorder

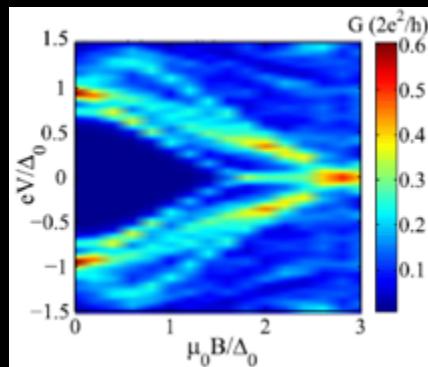


Mourik et al (2012)

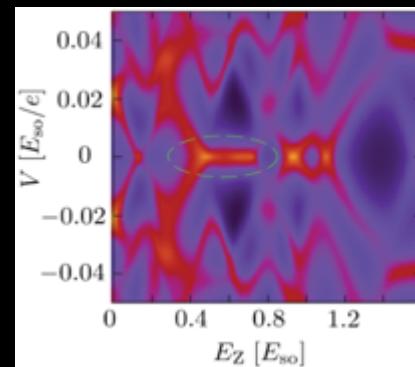


Takei et al. (2013), Chang et al (2015)

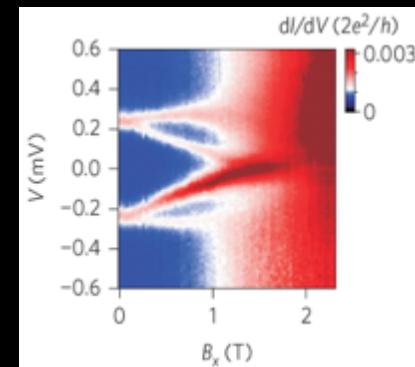
soft gap linked to disorder



Liu et al (2012)



Pikulin et al (2012)



Lee et al (2014)

Majorana signatures mimicked by disorder

Conclusions

Majorana recipe

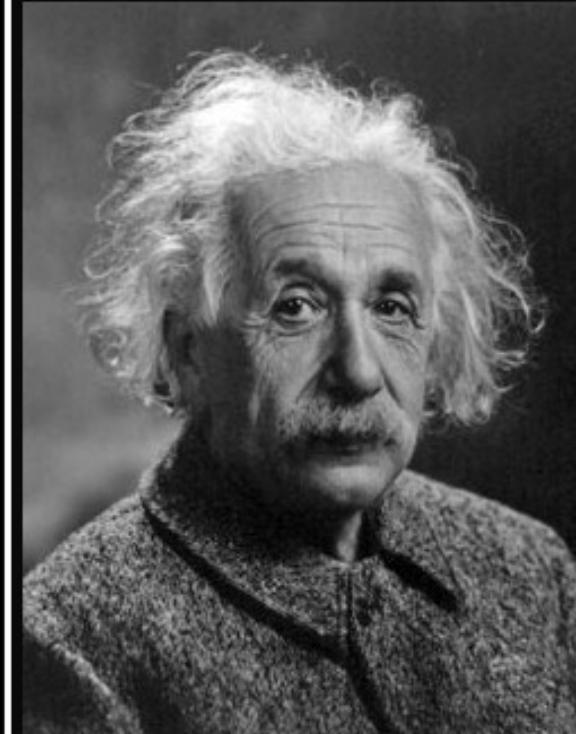
1. One-dimensional quantum wire
2. Spin-orbit interaction
3. Superconductivity
4. Apply magnetic field

*(Lutchyn, Sau, Das Sarma, PRL 2010,
Oreg, Refael, von Oppen, PRL 2010)*

Unique recipe for Majoranas.

Take one ingredient out and Majoranas are gone.

ZBP rigid over large ranges in B and Vg .



No amount of experimentation can ever prove me right; a single experiment can prove me wrong.

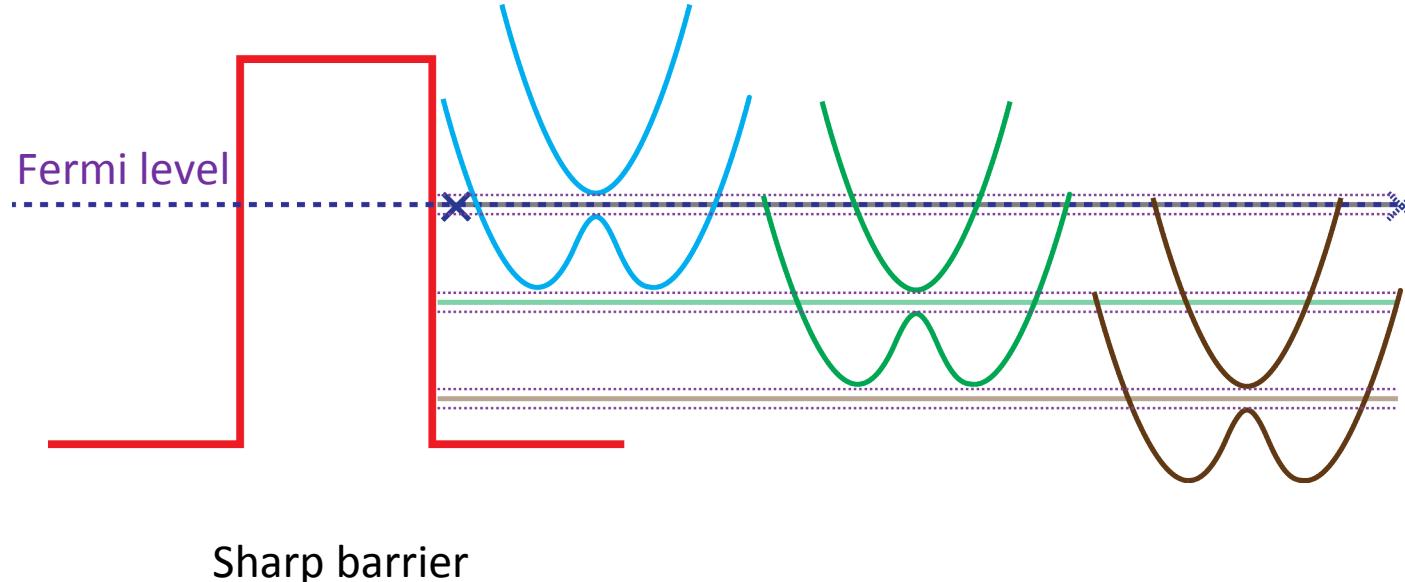
(Albert Einstein)

izquotes.com

Theories are excluded that are based on:

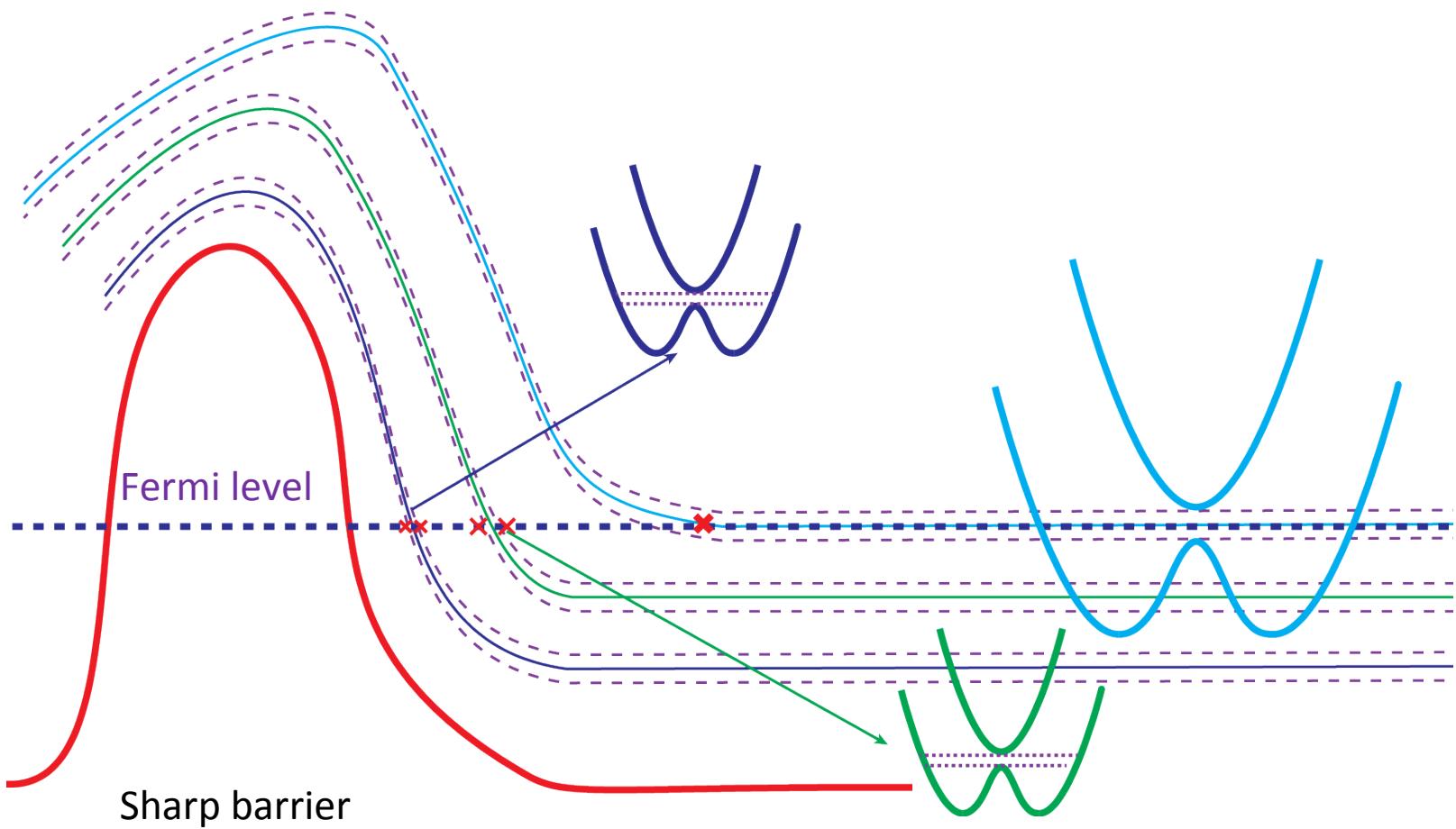
- an effect from the tunnel barrier
- $T \sim 1 \Rightarrow$ no Coulomb blockade
- ballistic wires \Rightarrow no localization effects
- Kondo only possible if $g < 0.5 = g_{InSb}/100$

BARRIER



1. With a sharp tunnel barrier, only the top populated subband will be topological (holds Majorana's) and the other two bands as shown in the plot are all trivial.
2. As the tunnel barrier is sharp, the three bands have the same transmission coefficient, therefore Majorana's as ZBP in the tunneling measurement will be clear.

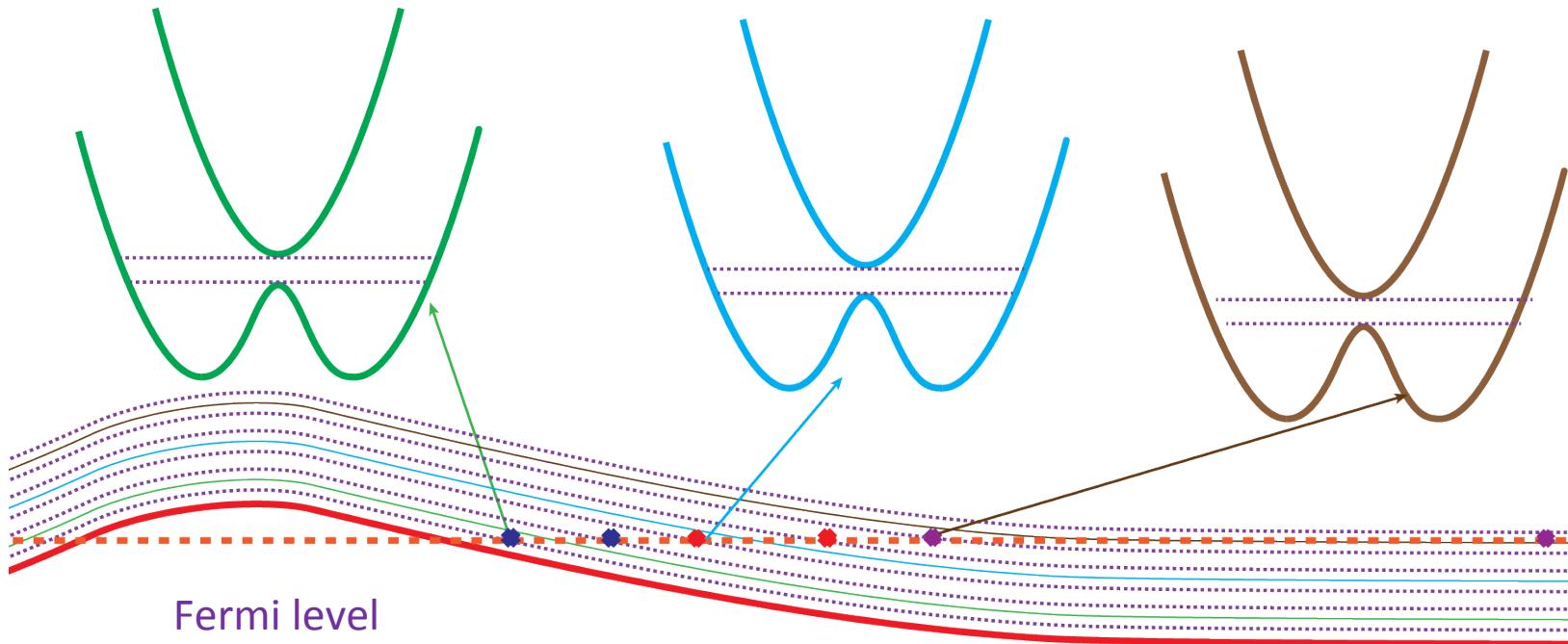
Dashed lines indicate the spin orbit gap



If we have a smooth barrier, the situation would be different

1. For smooth barrier, what is the relevant length scale?
2. What are the coherence length of Majorana's, could we have multiple Majorana's, like indicated in the figure.

Dashed lines indicate the spin orbit gap

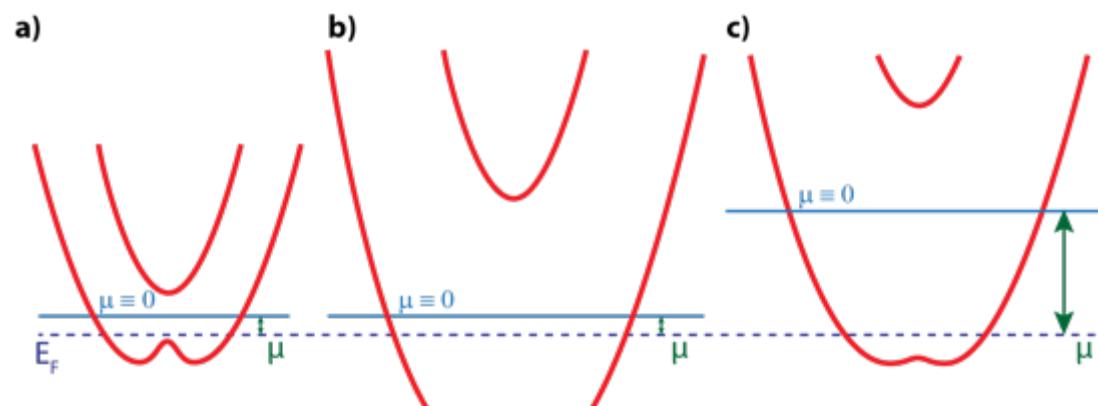


An more exaggerated case, tunnel barrier is very smooth, could be relevant for T. Stanscu's soft gap paper, the two more possible Majorana's pairs are well separated (still depend on the coherent length of Majorana's in the system):

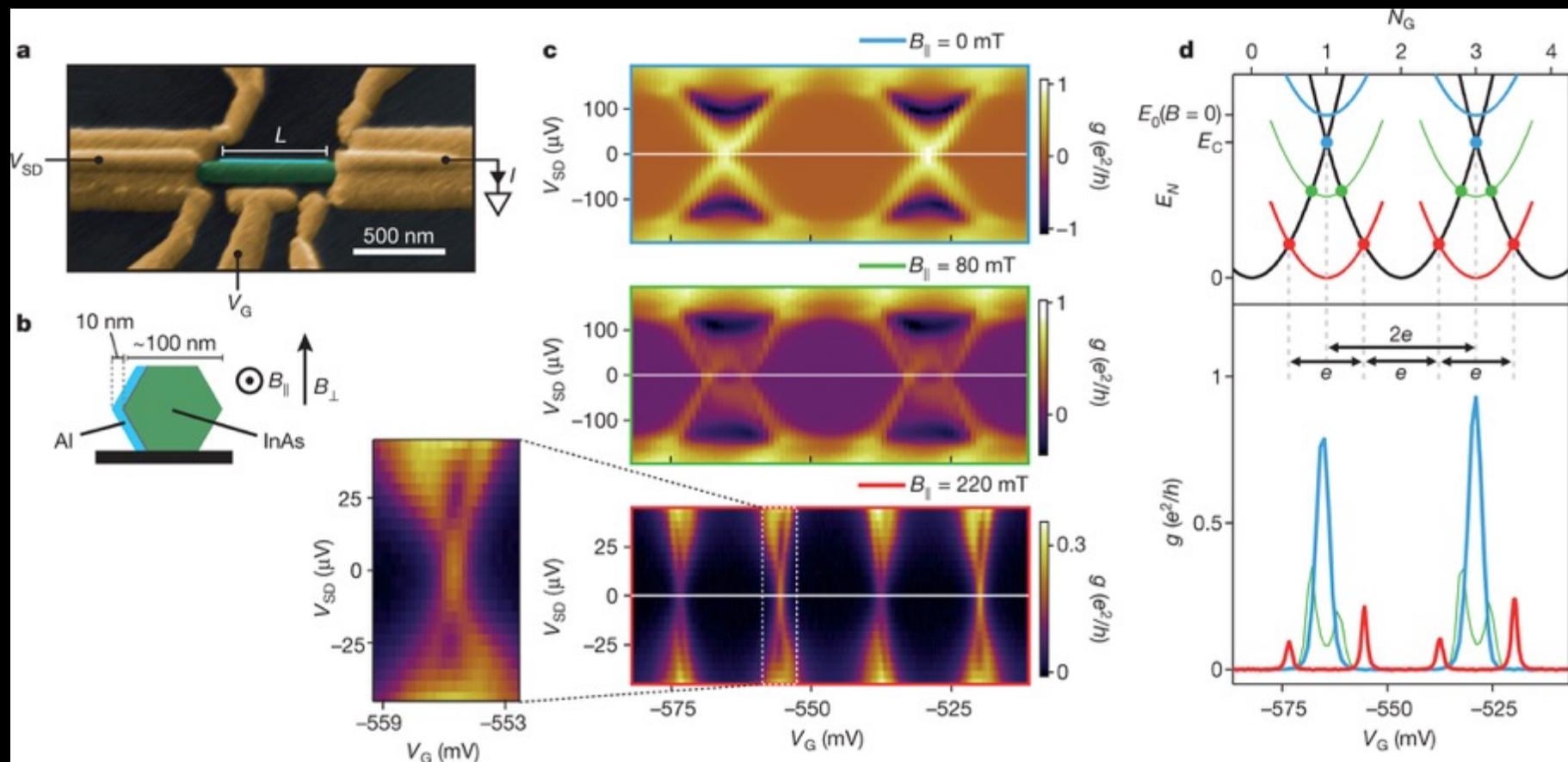
1. Could we discriminate them?
2. What is the visibility difference in tunneling measurements of the Majorana's at the top subband vs. lower subbands?
3. How about the stability in B and chemical potential of these Majorana's at lower subbands compared to top subband Majorana's?

Dashed lines indicate the spin orbit gap

rest



Majorana island device



S M Albrecht *et al.* *Nature* **531**, 206–209 (2016) doi:10.1038/nature17162

nature

