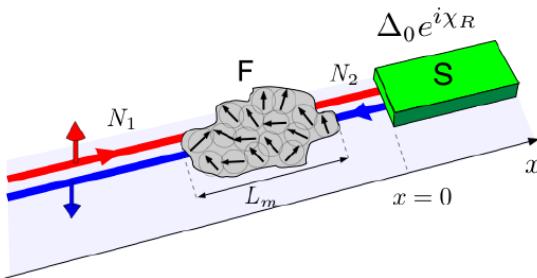




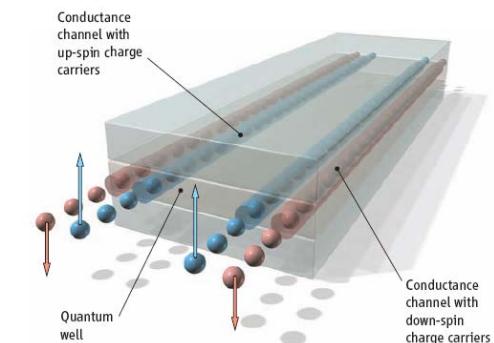
# Superconducting hybrids based on QSH systems

## Lecture 1



11th Capri Spring School  
*Transport in Nanostructures*

April 13-17, 2015  
Capri, Italy



# Björn Trauzettel

Pablo Burset (Uni Würzburg)  
François Crépin (Uni Würzburg)  
Fabrizio Dolcini (PolyTech Torino)  
Florian Geissler (Uni Würzburg)  
Ewelina Hankiewicz (Uni Würzburg)  
Naoto Nagaosa (Tokyo University)  
Yukio Tanaka (Nagoya University)  
Grigory Tkachov (Uni Würzburg)



Alexander von Humboldt  
Stiftung / Foundation

**DFG** Deutsche  
Forschungsgemeinschaft

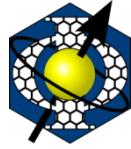
 HELMHOLTZ  
ASSOCIATION



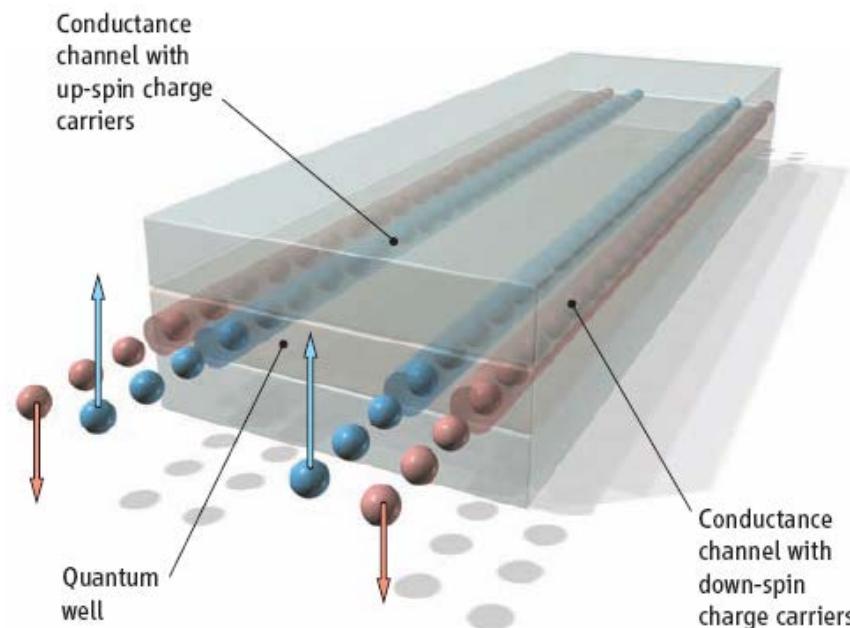


# Outline

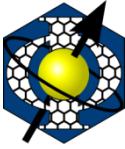
- **Recap:** Quantum spin Hall (QSH) systems
- **Introduction:** Hybrid structures based on QSH systems – Bogoliubov-De Gennes Formalism
- **Application:** Transport signatures in NS junctions



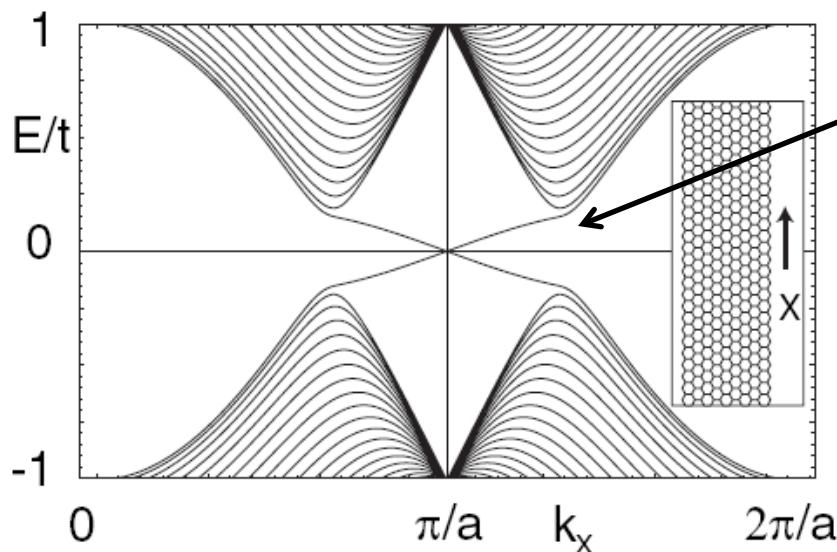
# What are QSH systems?



first symmetry-protected topological state of matter



# Prediction in graphene



spin filtered edge states:  
topologically non-trivial  
w/ respect to TRS

$$H = v(p_x \sigma_x \tau_x + p_y \sigma_y) + \Delta_{SO} \sigma_z \tau_z s_z$$

Kane & Mele PRL 2005

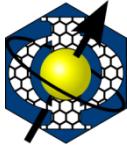
$$\Delta_{SO} \approx 24 \mu\text{eV}$$

Gmitra, Konschuh, Ertler, Ambrosch-Draxl & Fabian PRB 2009

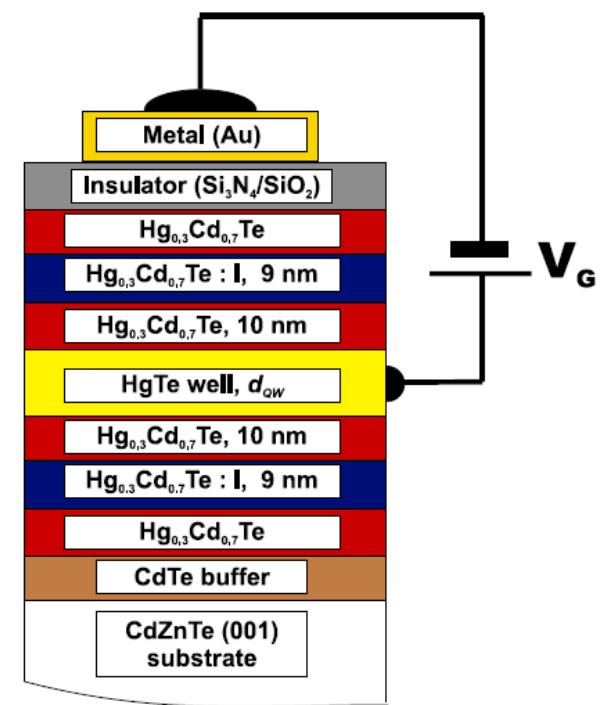
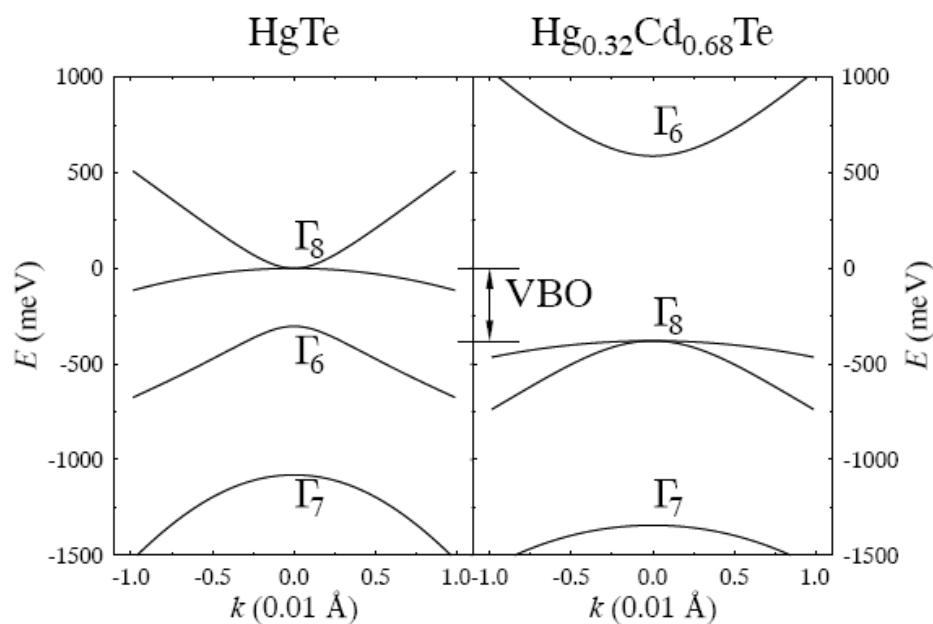


# SOI too small ...

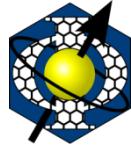
How about a better system?



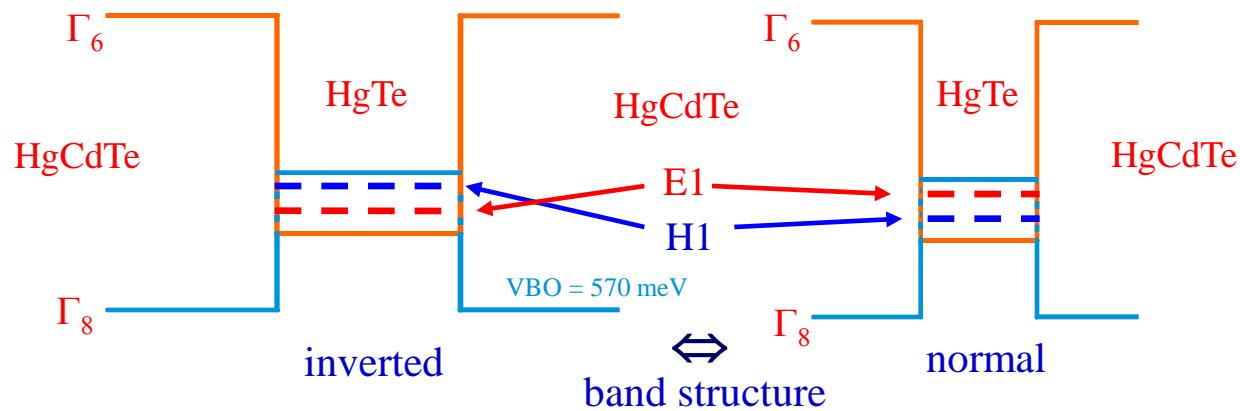
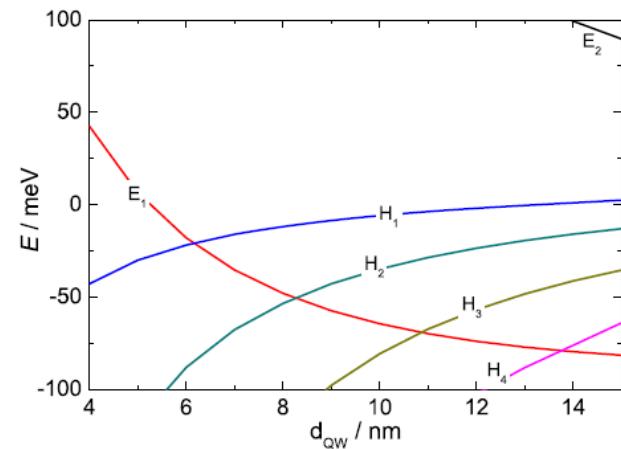
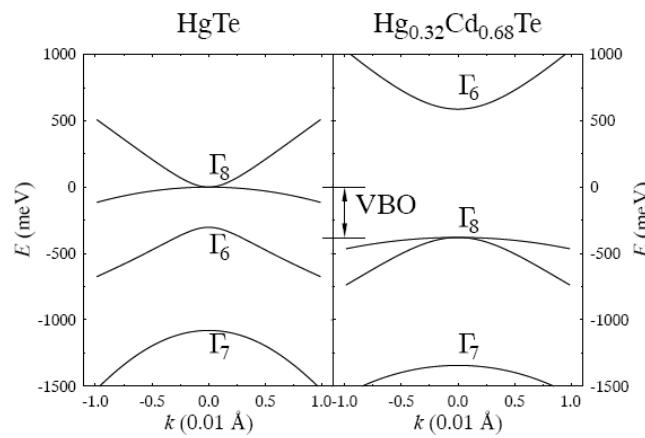
# HgTe/CdTe QW I

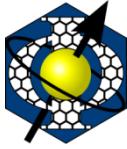


König et al. JPSJ 2008



# HgTe/CdTe QW II





# Effective Hamiltonian

$$H = \begin{pmatrix} h(\vec{k}) & 0 \\ 0 & h^*(-\vec{k}) \end{pmatrix} \quad \text{with} \quad h(\vec{k}) = \varepsilon(\vec{k}) + d_a(\vec{k})\sigma^a$$

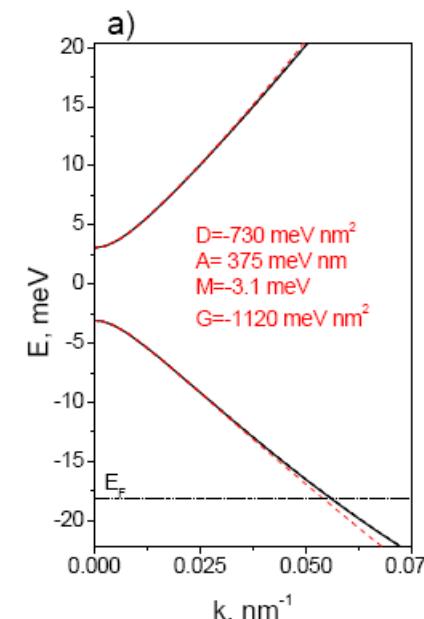
Bernevig, Hughes, Zhang Science 2006

$$\varepsilon(\vec{k}) = C - Dk^2$$

$$\vec{d}(\vec{k}) = (Ak_x, -Ak_y, M - Gk^2)$$

with basis states:

$$\{|E+\rangle, |H+\rangle, |E-\rangle, |H-\rangle\}$$

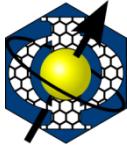


comparison  
to 8-band  
Kane model

±: degenerate Kramers partners

Novik et al. PRB 2005

Schmidt, Novik, Kindermann & BT PRB 2009



# Topology in the QSHE I

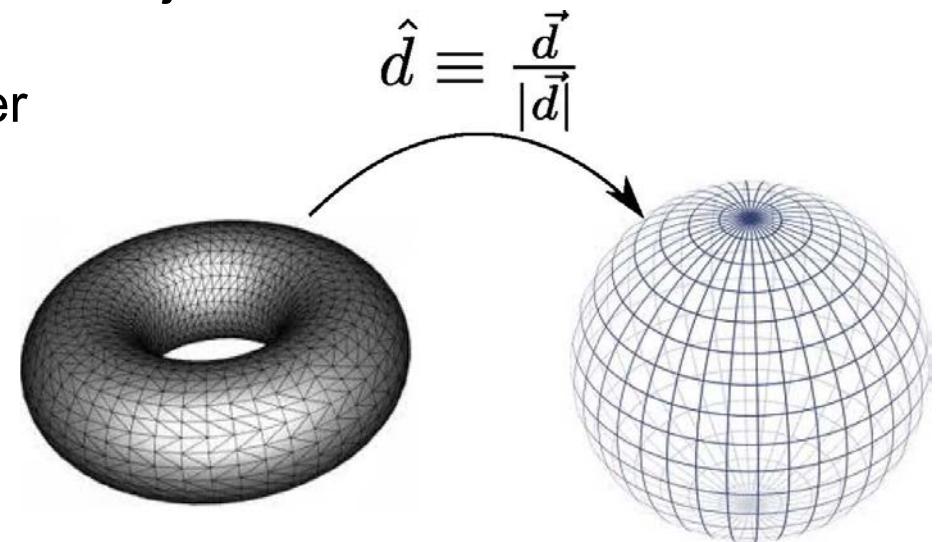
$$H = \begin{pmatrix} h(\vec{k}) & 0 \\ 0 & h^*(-\vec{k}) \end{pmatrix} \quad \text{with} \quad h(\vec{k}) = \varepsilon(\vec{k}) + d_a(\vec{k}) \sigma^a$$

TRS preserved

Thouless, Kohmoto, Nightingale & Den Nijs

TKNN invariant for one Kramers partner  
(quantum anomalous Hall effect):

$$n_{\uparrow} = \frac{1}{4\pi} \int d\mathbf{k} (\partial_{k_x} \hat{d} \times \partial_{k_y} \hat{d}) \cdot \hat{d}$$





# Topology in the QSHE I

$$H = \begin{pmatrix} h(\vec{k}) & 0 \\ 0 & h^*(-\vec{k}) \end{pmatrix} \quad \text{with} \quad h(\vec{k}) = \varepsilon(\vec{k}) + d_a(\vec{k})\sigma^a$$

TRS preserved

TKNN invariant for one Kramers partner:

$$n_\uparrow = \frac{1}{4\pi} \int d\mathbf{k} \left( \partial_{k_x} \hat{d} \times \partial_{k_y} \hat{d} \right) \cdot \hat{d}$$

zero Hall conductivity:

$$n = n_\uparrow + n_\downarrow = 0$$

quantized spin Hall conductivity:

$$n_\sigma = \frac{n_\uparrow - n_\downarrow}{2}$$



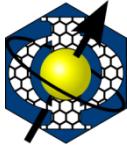
# Topology in the QSHE II

quantized spin Hall conductivity:

$$n_{\sigma} = \frac{n_{\uparrow} - n_{\downarrow}}{2}$$

->  $Z_2$  invariant:

$$\nu = n_{\sigma} \bmod 2$$



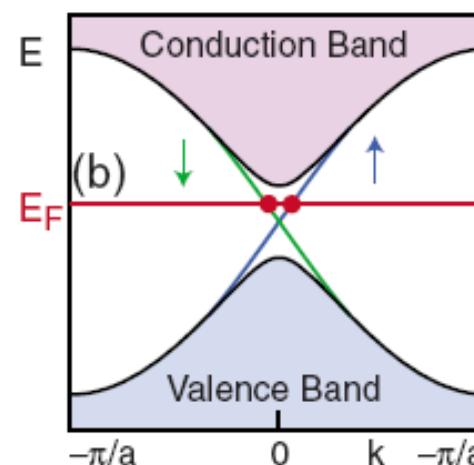
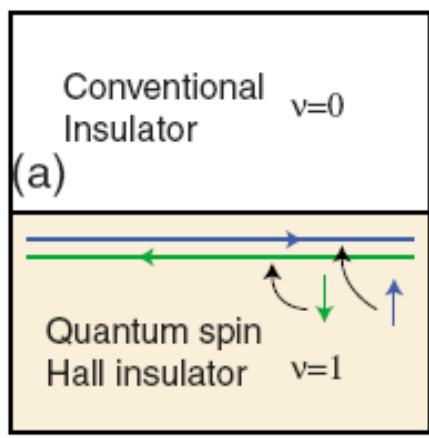
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quantized spin Hall conductivity:

$$n_{\sigma} = \frac{n_{\uparrow} - n_{\downarrow}}{2}$$

$\rightarrow Z_2$  invariant:

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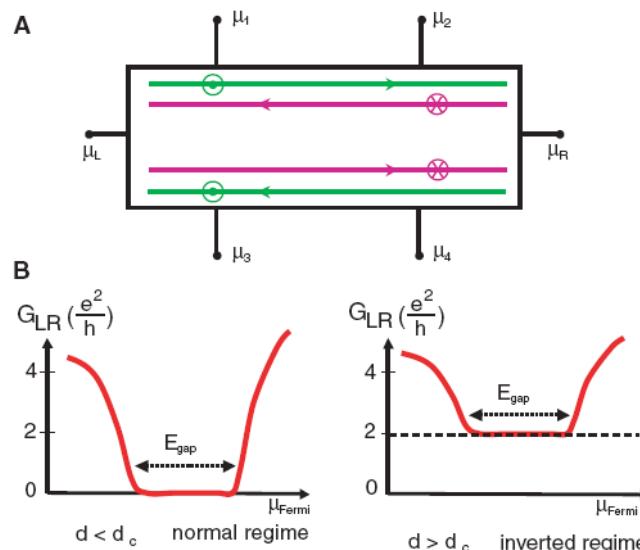
# of Kramers pairs at edge:

$$N_K = \Delta\nu \bmod 2$$

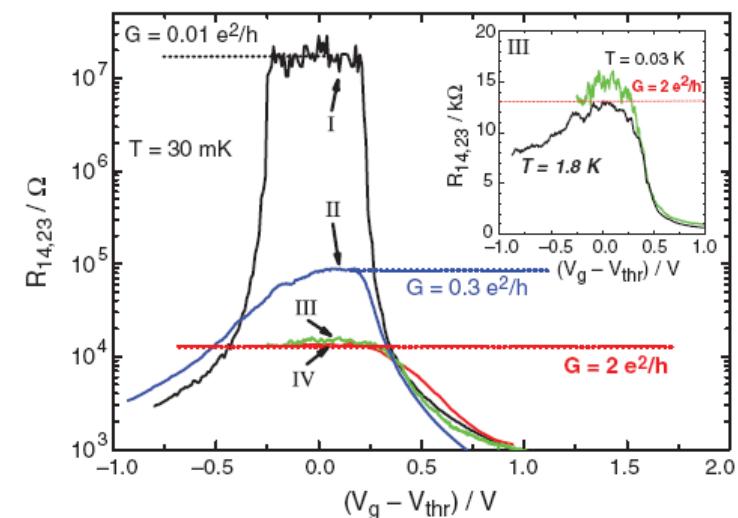


# Experimental evidence of edge states

Prediction:



Observation:



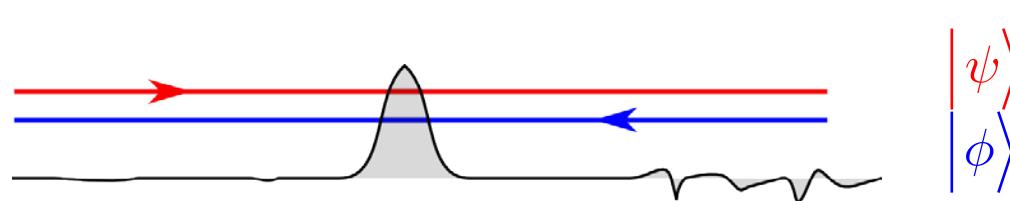
Bernevig, Hughes & Zhang Science 2006

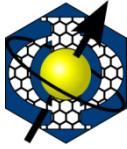
König et al. Science 2007



# Protection against elastic backscattering

$$\begin{aligned}\langle \psi | H | \phi \rangle &= \langle \phi | H | \psi \rangle^* = \langle T\phi | TH | \psi \rangle \\ &= \langle \psi | HT | \psi \rangle = \langle \psi | HT^2 | \phi \rangle = -\langle \psi | H | \phi \rangle\end{aligned}$$

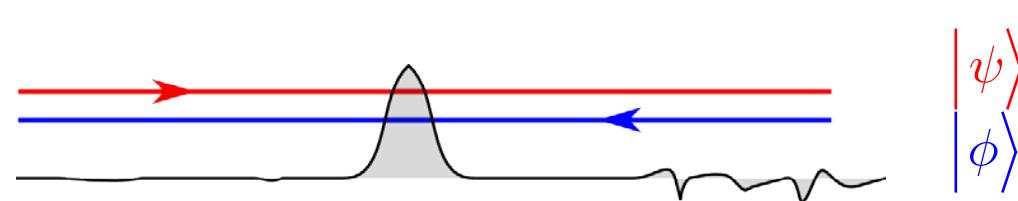


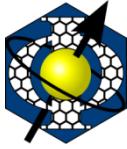


# Protection against elastic backscattering

$$\langle \lambda | \mu \rangle = \langle \mu | \lambda \rangle^*$$

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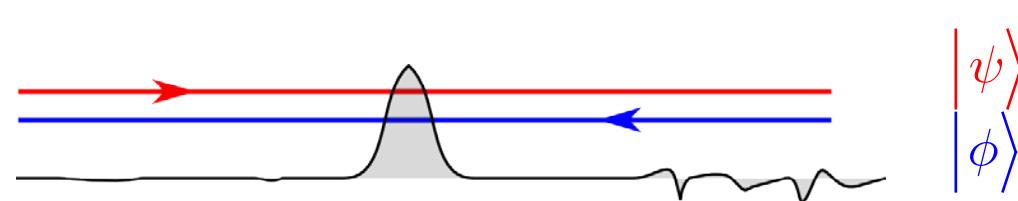


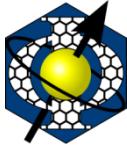
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$$\langle \lambda | \mu \rangle = \langle \mu | \lambda \rangle^*$$

$T$ : anti-unitary operator

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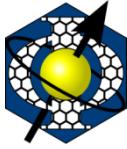
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$$[H, T] = 0 \text{ and } |\psi\rangle = T|\phi\rangle$$



# Protection against elastic backscattering

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$$[H, T] = 0 \text{ and } |\psi\rangle = T|\phi\rangle$$

$$T^2 = -1$$

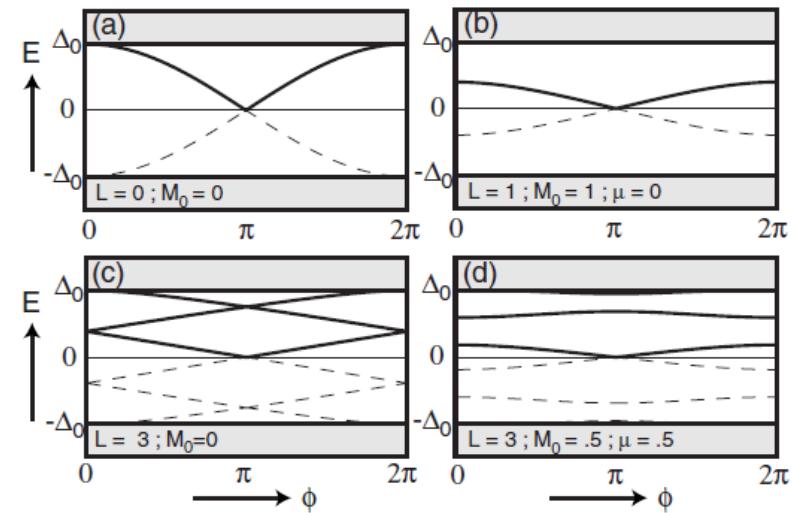
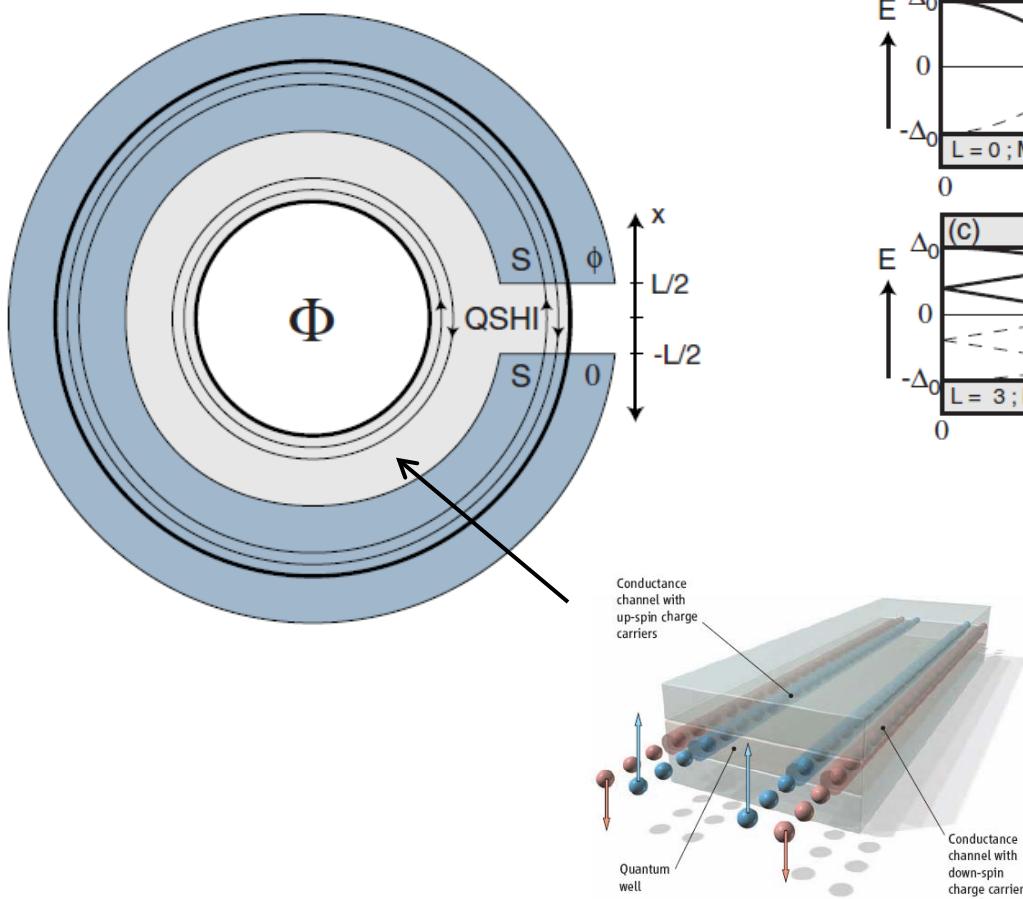


# Outline

- Recap: Quantum spin Hall (QSH) systems
- **Introduction:** Hybrid structures based on QSH systems – Bogoliubov-De Gennes Formalism
- Application: Transport signatures in NS junctions



# Pioneering prediction

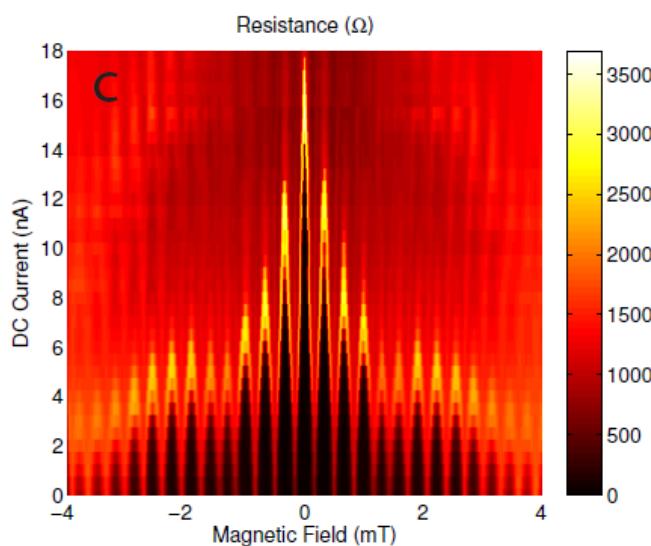


signatures of p-wave  
superconductivity?



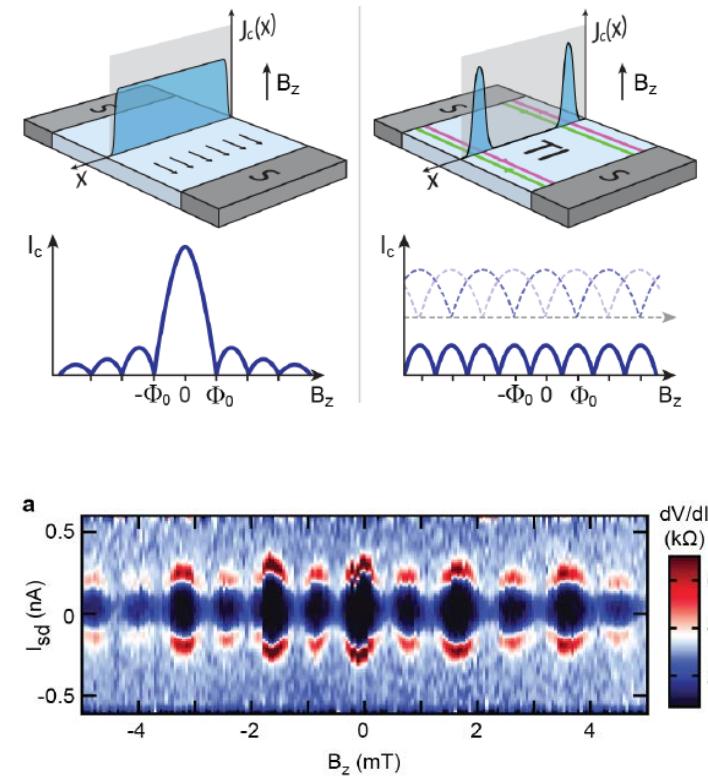
# Inspiring experiments

Hg(Cd)Te QWs

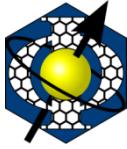


Hart et al. Nature Phys 2014

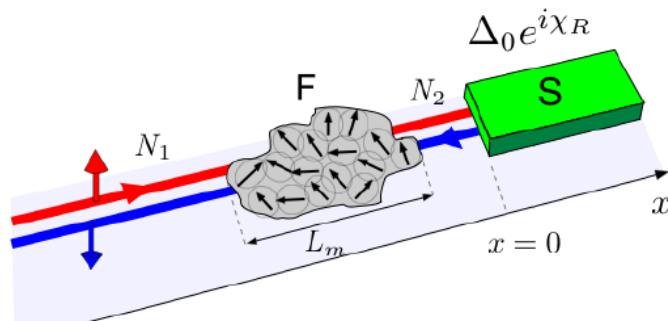
InAs/GaSb QWs



Pribiag et al. arXiv 2014

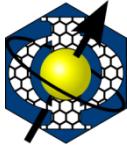


# Setup & Hamiltonian

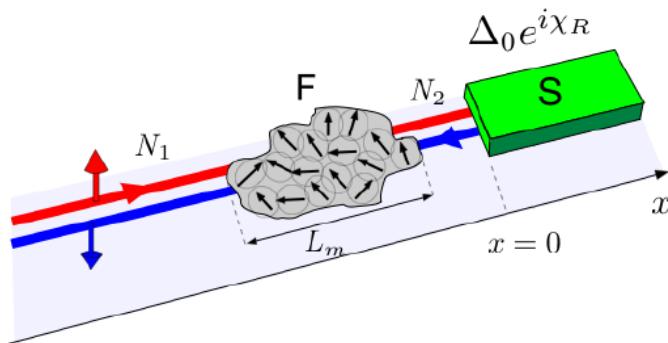


$$H = H_0 + H_{FM} + H_{SC}$$

$$H_0 = \int dx \left( \psi_{R\uparrow}^\dagger, \psi_{L\downarrow}^\dagger \right) \left[ v_F p_x \sigma_z - \mu \right] \begin{pmatrix} \psi_{R\uparrow} \\ \psi_{L\downarrow} \end{pmatrix}$$



# Setup & Hamiltonian



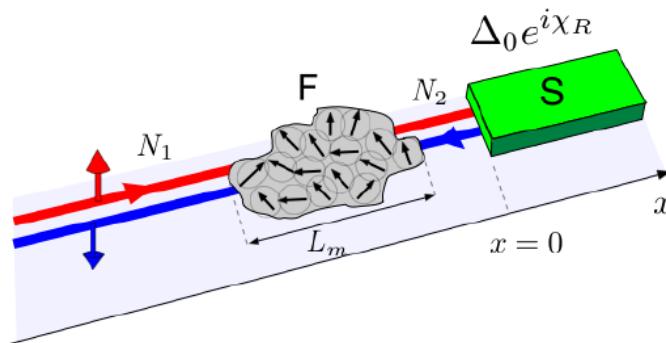
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$$H_{FM} = \int dx \left( \psi_{R\uparrow}^\dagger, \psi_{L\downarrow}^\dagger \right) \vec{m}(x) \cdot \vec{\sigma} \begin{pmatrix} \psi_{R\uparrow} \\ \psi_{L\downarrow} \end{pmatrix}$$



# Setup & Hamiltonian



$$H = H_0 + H_{FM} + H_{SC}$$

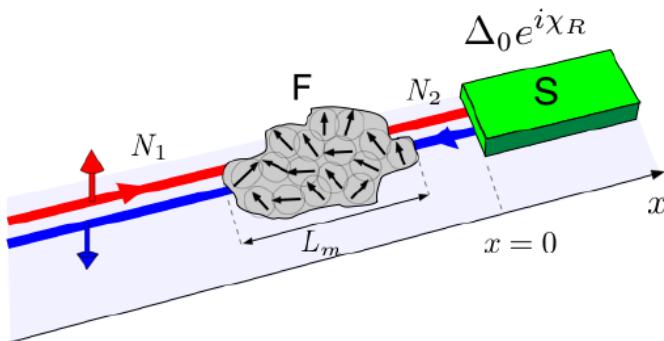
$$H_0 = \int dx \left( \psi_{R\uparrow}^\dagger, \psi_{L\downarrow}^\dagger \right) \left[ v_F p_x \sigma_z - \mu \right] \begin{pmatrix} \psi_{R\uparrow} \\ \psi_{L\downarrow} \end{pmatrix}$$

$$H_{FM} = \int dx \left( \psi_{R\uparrow}^\dagger, \psi_{L\downarrow}^\dagger \right) \vec{m}(x) \cdot \vec{\sigma} \begin{pmatrix} \psi_{R\uparrow} \\ \psi_{L\downarrow} \end{pmatrix}$$

$$H_{SC} = \int dx \left[ \Delta(x) \psi_{R\uparrow}^\dagger \psi_{L\downarrow}^\dagger + \Delta^*(x) \psi_{L\downarrow} \psi_{R\uparrow} \right]$$



# BdG Hamiltonian



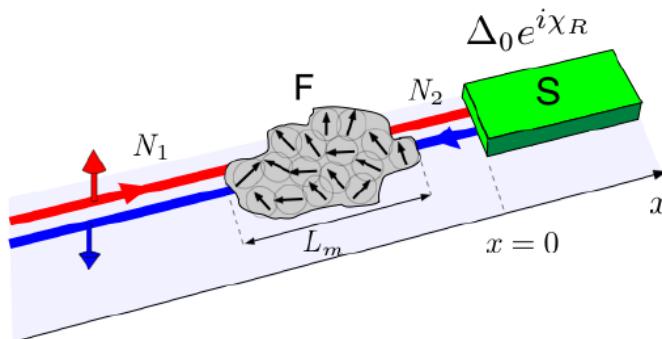
$$H = \frac{1}{2} \int dx \Psi^\dagger \color{red} H_{BdG} \color{black} \Psi$$

$$\Psi^\dagger = (\psi_{R\uparrow}^\dagger, \psi_{L\downarrow}^\dagger, \psi_{L\downarrow}, -\psi_{R\uparrow})$$

$$\color{red} H_{BdG} = \begin{pmatrix} H_{0+FM}^e & \Delta(x)\sigma_0 \\ \Delta^*(x)\sigma_0 & H_{0+FM}^h \end{pmatrix}$$



# BdG Hamiltonian



$$H = \frac{1}{2} \int dx \Psi^\dagger H_{BdG} \Psi$$

$$\Psi^\dagger = (\psi_{R\uparrow}^\dagger, \psi_{L\uparrow}^\dagger, \psi_{L\downarrow}, -\psi_{R\uparrow})$$

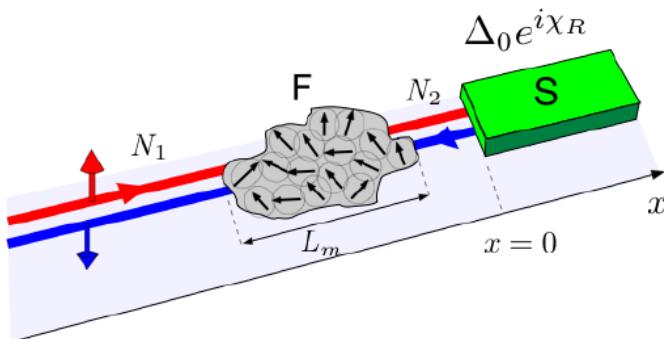
$$H_{0+FM}^e = v_F \sigma_z p_x - \mu \sigma_0 + \vec{m}(x) \cdot \vec{\sigma}$$

$$H_{BdG} = \begin{pmatrix} H_{0+FM}^e & \Delta(x) \sigma_0 \\ \Delta^*(x) \sigma_0 & H_{0+FM}^h \end{pmatrix}$$

$$H_{0+FM}^h = -T H_{0+FM}^e T^{-1} = -v_F \sigma_z p_x + \mu \sigma_0 + \vec{m}(x) \cdot \vec{\sigma}$$



# BdG Hamiltonian



$$H = \frac{1}{2} \int dx \Psi^\dagger H_{BdG} \Psi$$

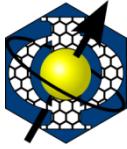
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$$H_{BdG} = \begin{pmatrix} H_{0+FM}^e & \Delta(x) \sigma_0 \\ \Delta^*(x) \sigma_0 & H_{0+FM}^h \end{pmatrix}$$

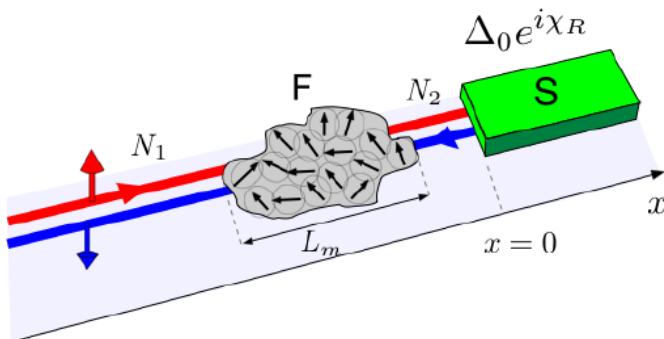
$$H = \sum_{\varepsilon_n \geq 0, j} \varepsilon_n \gamma_{\varepsilon_n, j}^\dagger \gamma_{\varepsilon_n, j}$$

$$\Psi(x) = \sum_{\varepsilon_n \geq 0, j} \left\{ \varphi_{\varepsilon_n, j}(x) \gamma_{\varepsilon_n, j} + [C \varphi_{\varepsilon_n, j}](x) \gamma_{\varepsilon_n, j}^\dagger \right\}$$

charge conjugated wave function



# Symmetries & Majoranas

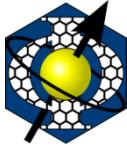


Particle-hole symmetry

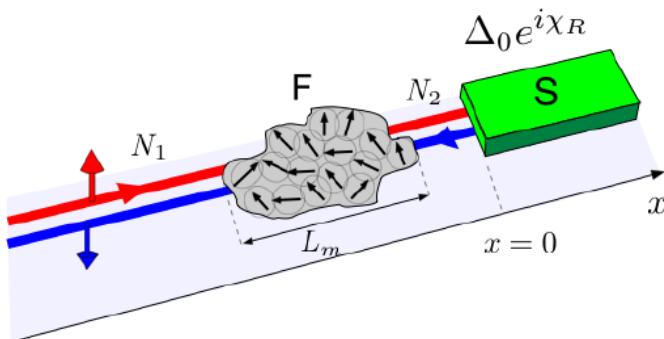
$$CH_{BdG}C^{-1} = -H_{BdG}$$

$$\gamma_{\varepsilon_n, j}^\dagger = \gamma_{-\varepsilon_n, j}$$

$$C = K \tau_y \otimes \sigma_y = K U_C$$



# Symmetries & Majoranas



Particle-hole symmetry

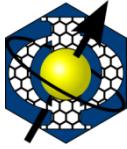
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$$C = K\tau_y \otimes \sigma_y = KU_C$$

$$\Psi(x) = \sum_{\varepsilon_n \geq 0, j} \left\{ \varphi_{\varepsilon_n,j}(x) \gamma_{\varepsilon_n,j} + [C \varphi_{\varepsilon_n,j}](x) \gamma_{\varepsilon_n,j}^\dagger \right\}$$

Majorana fermion



# Majorana fermions vs. anyons



## Majorana fermions

$$\Psi(x) = \sum_{\varepsilon_n \geq 0, j} \left\{ \varphi_{\varepsilon_n, j}(x) \gamma_{\varepsilon_n, j} + [C \varphi_{\varepsilon_n, j}](x) \gamma_{\varepsilon_n, j}^\dagger \right\}$$

fermions

$$\Psi^\dagger(x) = U_C \Psi(x)$$

$$\gamma_{\varepsilon_n, j}^\dagger = \gamma_{-\varepsilon_n, j}$$



# Majorana fermions vs. anyons



Majorana fermions

$$\Psi(x) = \sum_{\varepsilon_n \geq 0, j} \left\{ \varphi_{\varepsilon_n, j}(x) \gamma_{\varepsilon_n, j} + [C \varphi_{\varepsilon_n, j}](x) \gamma_{\varepsilon_n, j}^\dagger \right\}$$

fermions

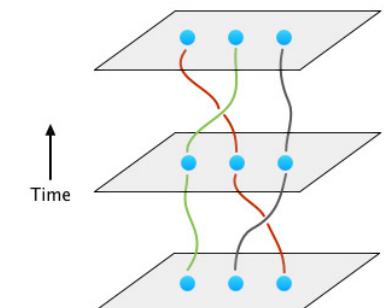
$$\Psi^\dagger(x) = U_C \Psi(x)$$

$$\gamma_{\varepsilon_n, j}^\dagger = \gamma_{-\varepsilon_n, j}$$

Majorana bound states

$$\gamma_{0,j}^\dagger = \gamma_{0,j}$$

anyons





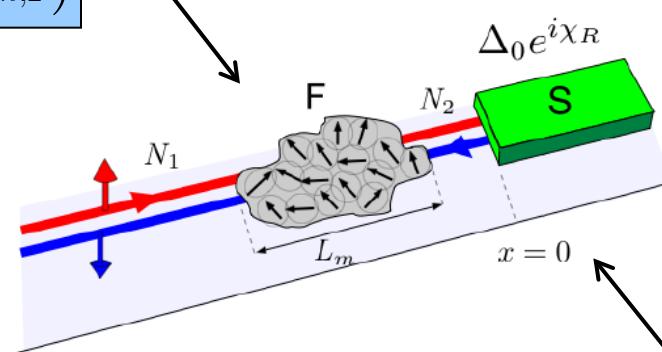
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- Application: Transport signatures in NS junctions



# S-matrix construction

$$\begin{pmatrix} b_{e,1} \\ b_{e,2} \\ b_{h,1} \\ b_{h,2} \end{pmatrix} = \begin{pmatrix} S_0^e(\varepsilon) & 0 \\ 0 & S_0^h(\varepsilon) \end{pmatrix} \begin{pmatrix} a_{e,1} \\ a_{e,2} \\ a_{h,1} \\ a_{h,2} \end{pmatrix}$$



perfect AR

$$\begin{pmatrix} a_{e,2} \\ a_{h,2} \end{pmatrix} = \exp\left(-i \arccos\left(\frac{\varepsilon}{\Delta_0}\right)\right) \begin{pmatrix} 0 & e^{i\chi_R} \\ e^{-i\chi_R} & 0 \end{pmatrix} \begin{pmatrix} b_{e,2} \\ b_{h,2} \end{pmatrix}$$



# S-matrix of FM domain

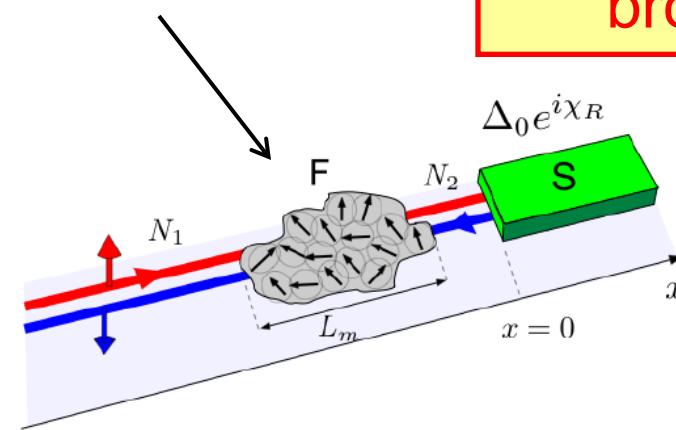
$$\Gamma_m(\varepsilon) \sim k_F L_m$$

$$\Phi_m(\varepsilon) \sim k_F x_0$$

$$S_0^e(\varepsilon) = e^{i\Gamma_m(\varepsilon)} \begin{pmatrix} -ie^{i\Phi_m(\varepsilon)} \sqrt{1-T_\varepsilon} & e^{i\chi_m(\varepsilon)} \sqrt{T_\varepsilon} \\ e^{-i\chi_m(\varepsilon)} \sqrt{T_\varepsilon} & -ie^{-i\Phi_m(\varepsilon)} \sqrt{1-T_\varepsilon} \end{pmatrix}$$

$$\chi_m(\varepsilon) \sim m_z L_m$$

all symmetries  
broken

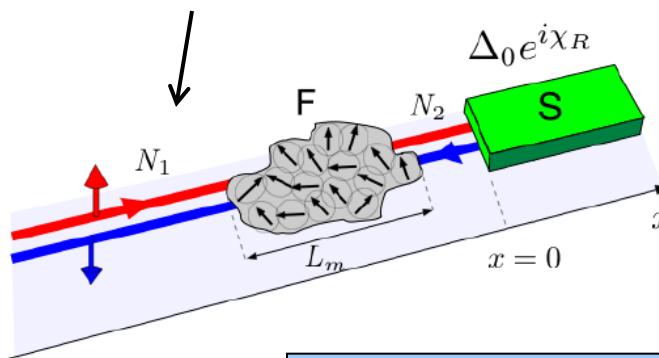




# Andreev reflection

$$\begin{pmatrix} b_{e,1} \\ b_{h,1} \end{pmatrix} = \begin{pmatrix} r_{ee} & \color{red}r_{eh} \\ \color{red}r_{he} & r_{hh} \end{pmatrix} \begin{pmatrix} a_{e,1} \\ a_{h,1} \end{pmatrix}$$

$$R_A = |\color{red}r_{eh}|^2 = |\color{red}r_{he}|^2$$



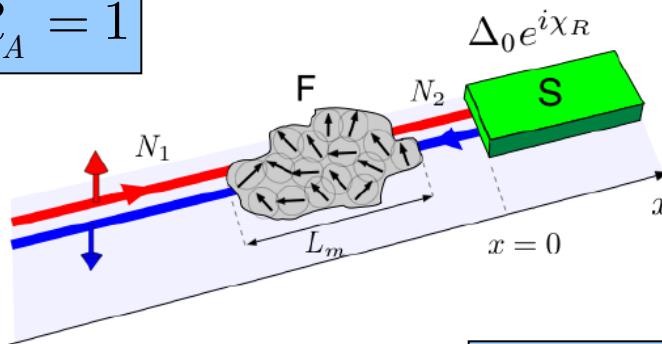
$$\Phi_m^A(\varepsilon) = \frac{1}{2} (\Phi_m(\varepsilon) - \Phi_m(-\varepsilon))$$

$$R_A = \frac{T_\varepsilon T_{-\varepsilon}}{\left(1 - \sqrt{R_\varepsilon R_{-\varepsilon}}\right)^2 + 4 \cos^2 \left[ \arccos \left( \frac{\varepsilon}{\Delta_0} \right) + \Phi_m^A(\varepsilon) \right] \sqrt{R_\varepsilon R_{-\varepsilon}}}$$



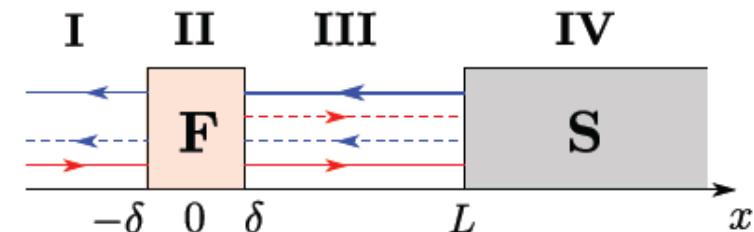
# Resonance condition

$$R_A = 1$$



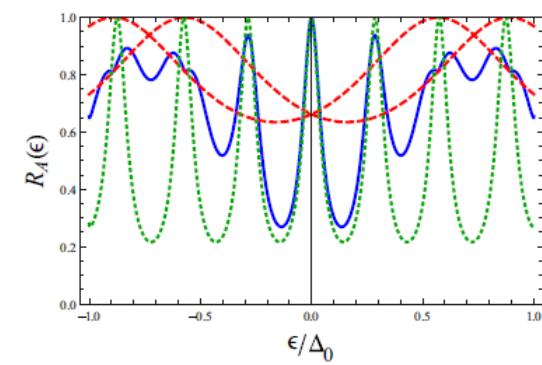
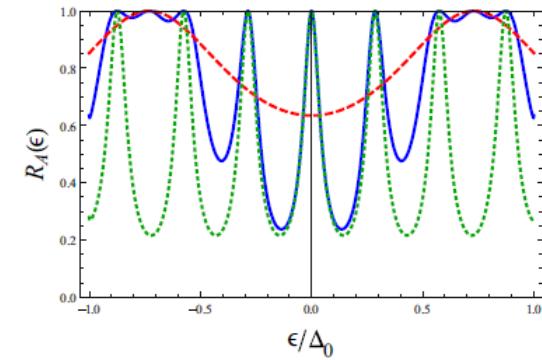
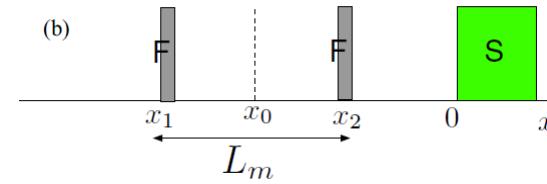
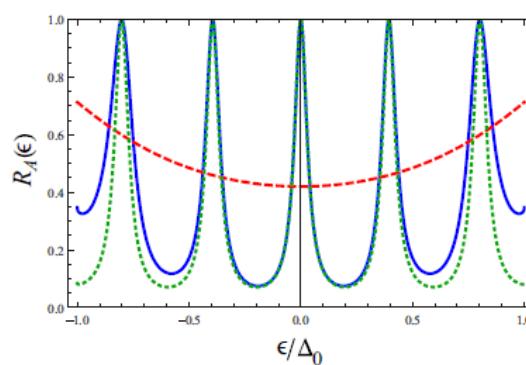
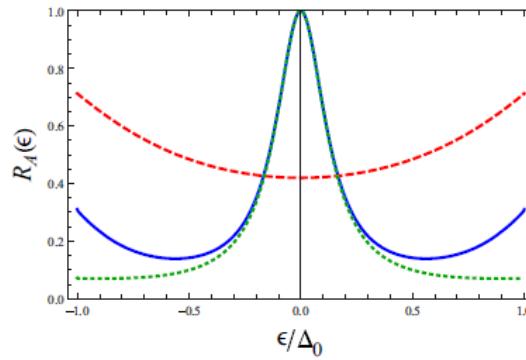
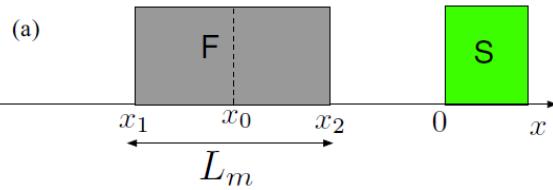
$$\left(\sqrt{R_\varepsilon} - \sqrt{R_{-\varepsilon}}\right)^2 + 4 \cos^2 \left[ \arccos \left( \frac{\varepsilon}{\Delta_0} \right) + \Phi_m^A (\varepsilon) \right] \sqrt{R_\varepsilon R_{-\varepsilon}} = 0$$

Fabry-Perot resonator  
for electrons/holes



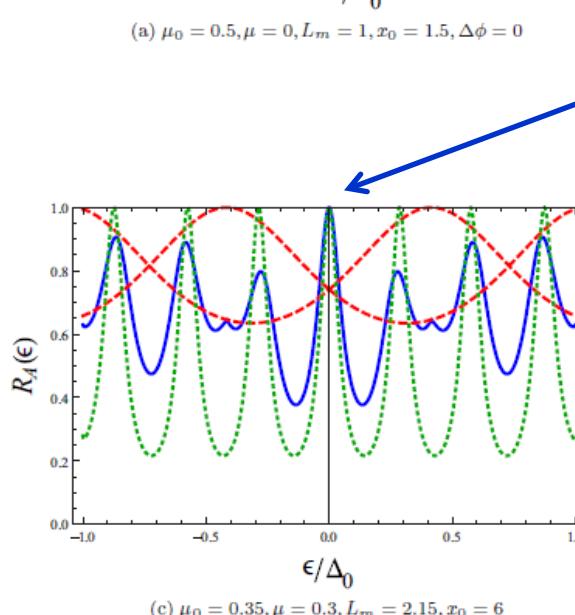
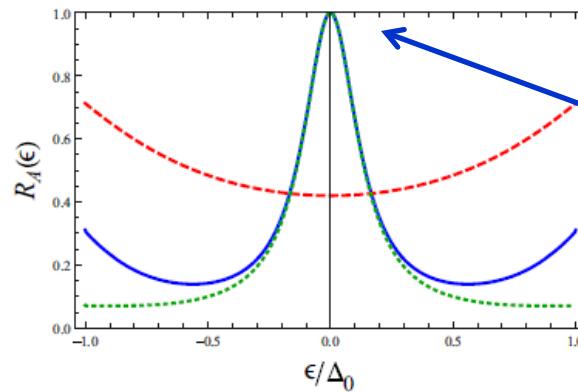


# FM domain vs. double barrier

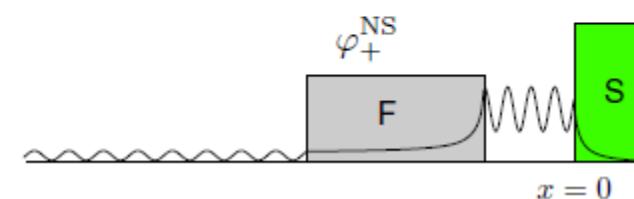
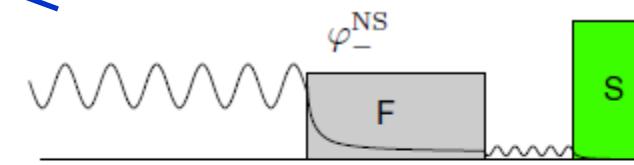




# Robust MBS signature

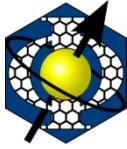


$$C\varphi_{\pm}^{NS} = \varphi_{\pm}^{NS}$$

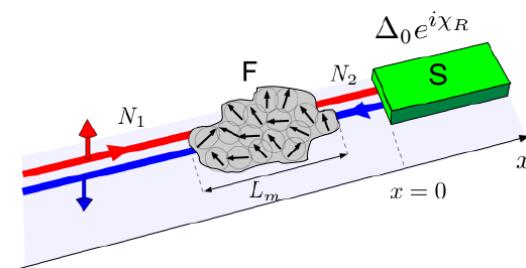
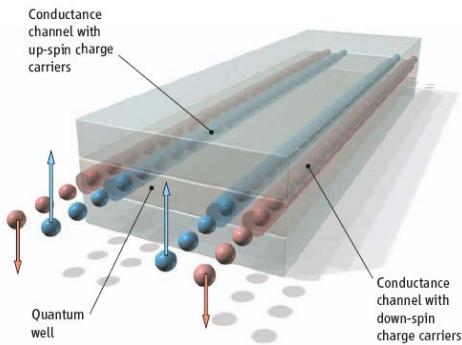


Majorana bound states

$$\varepsilon = 0$$



# Summary Lecture 1



$$H_{BdG} = \begin{pmatrix} H^e_{0+FM} & \Delta(x)\sigma_0 \\ \Delta^*(x)\sigma_0 & H^h_{0+FM} \end{pmatrix}$$

